

Maxim Yu. Khlopov

Centre for Cosmoparticle Physics
COSMION

Moscow, Russia

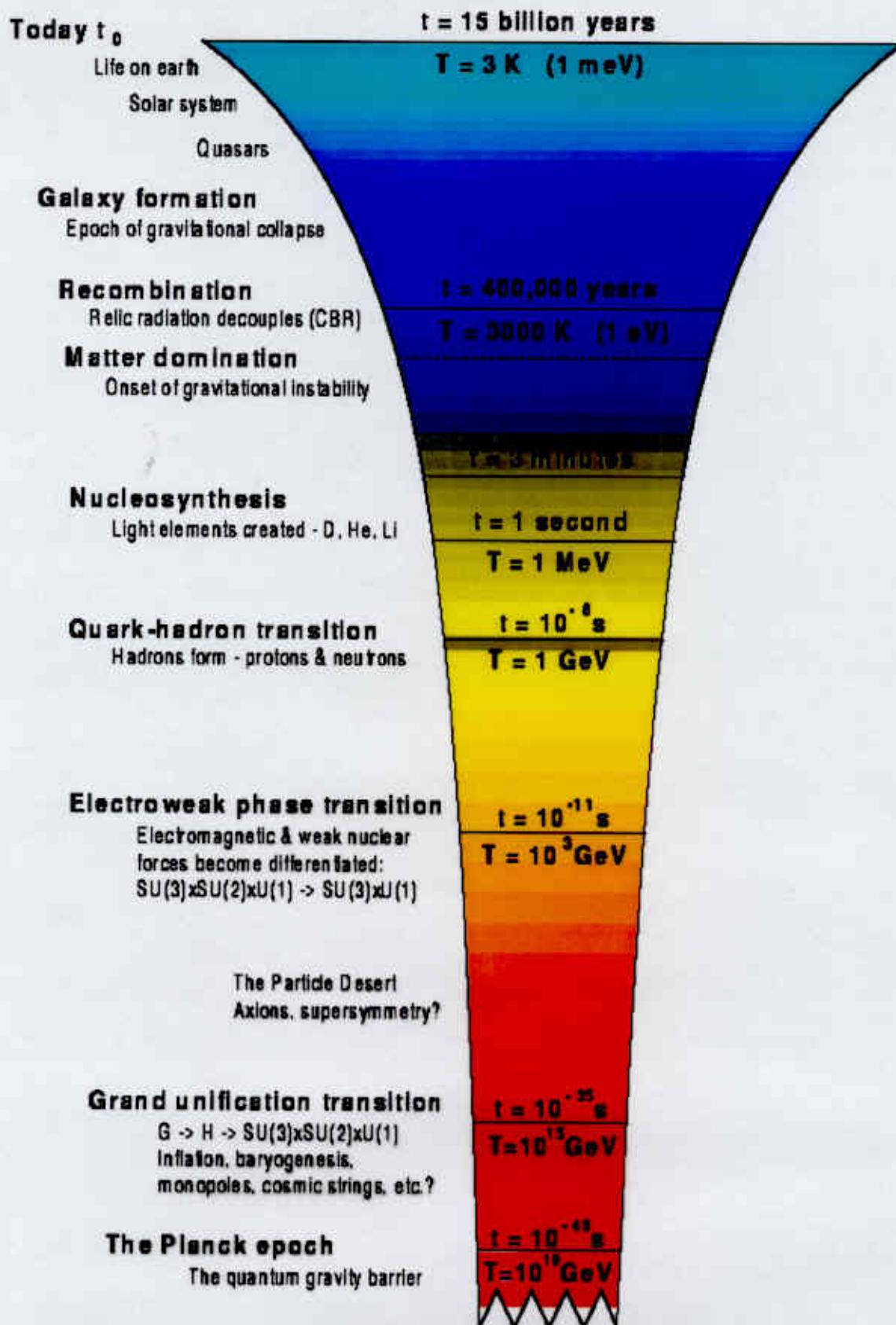
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COSMO PHENOMENOLOGY

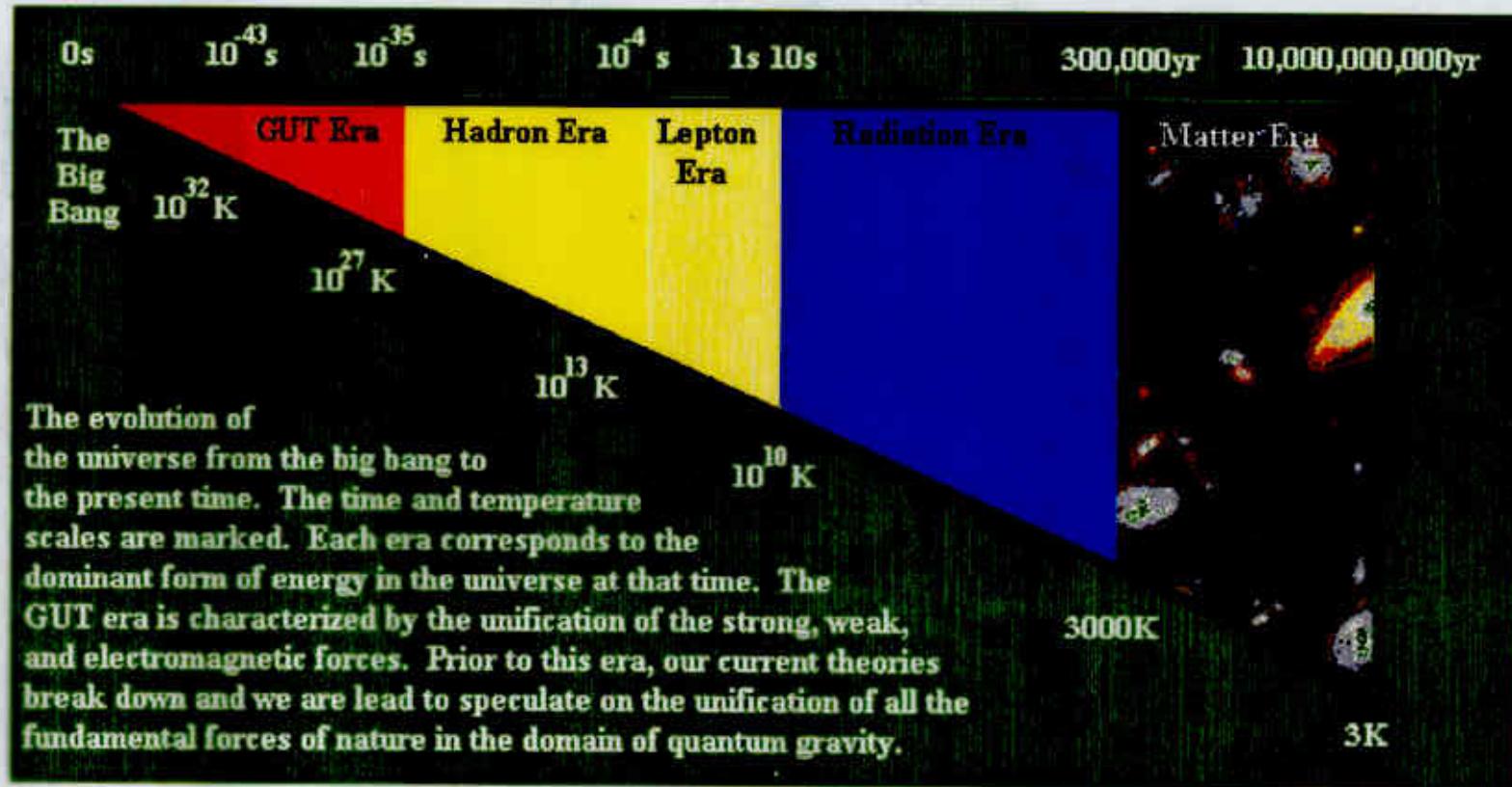
of

NEW PHYSICS

Big Bang history



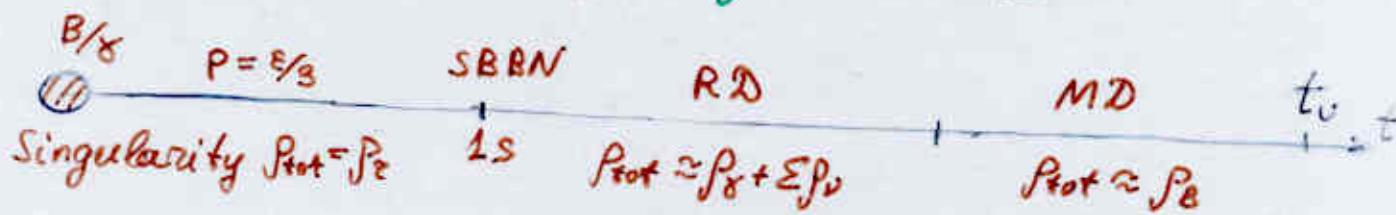
The thermal history of the Universe



New standards in cosmology

New paradigm of inflationary Universe with baryosynthesis and nonbaryonic dark matter

Old Big Bang scenario



New big Bang scenario



Trivial comment - The true history of the Universe should account for all the three phenomena (all the three are based on the extensions of the SM-cosmology), necessary for the modern cosmology \Rightarrow Nontrivial consequence - No practical realization of New big Bang scenario can be reduced to these three phenomena only. The true history of the Universe should be much more complicated.

The true mystery of the crisis of the old big Bang scenario is that it happens (if does, c.f. SBBN - ?) so late. The true cosmological picture is effectively masked by the old scenario. We should refine the astrophysical probes for new cosmological phenomena to the extent, sufficient for their discrimination.

Inflation

To solve the problems of Big Bang Universe (horizon, flatness, initial δT_p , primordial magnetic monopoles, ...)

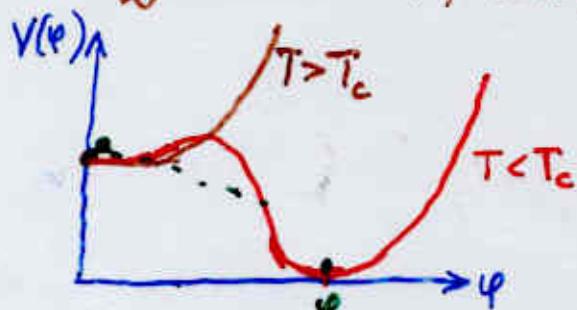
$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\epsilon + 3p) \quad \text{if } p < -\epsilon/3 \quad \frac{\ddot{a}}{a} > 0$$

deceleration *acceleration*

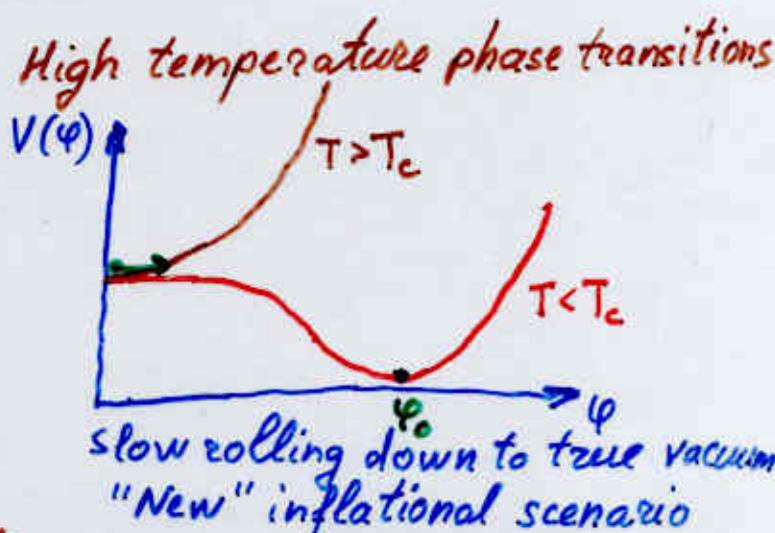
Inflaton? Its physical nature?

New elementary scalar field.

GUT Higgs scalar: φ , $V(\varphi, T)$



Strong I order phase transition.
"Old" inflational scenario



Chaotic inflation



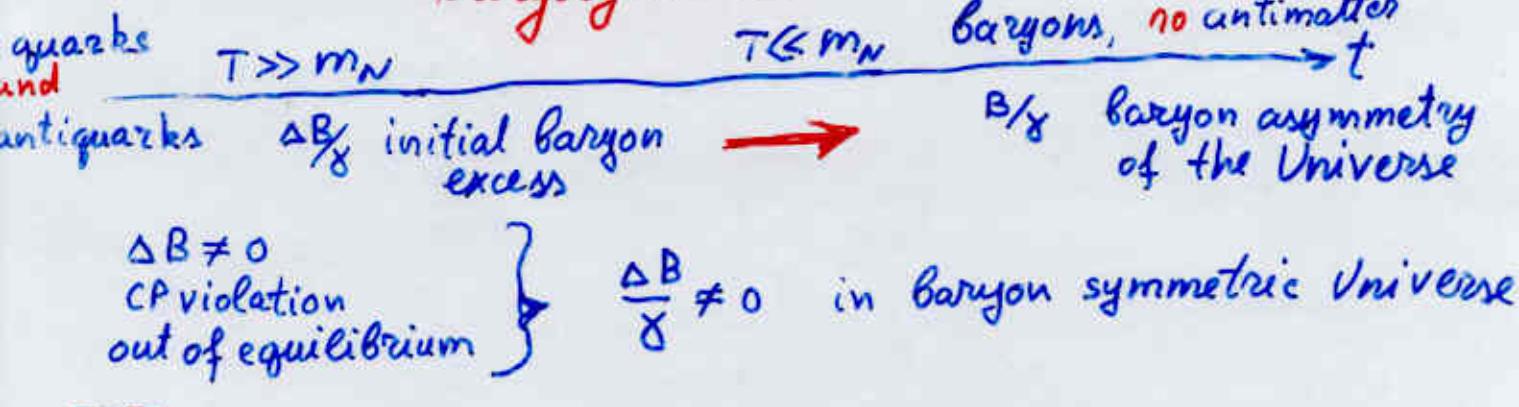
Slow rolling down to true vacuum for any scalar field φ
 c.f. $V(\varphi) = \frac{1}{2}m^2\varphi^2$

Extended inflation

Self supported inflation

Phenomenology of inflation (A.D.Linde)
should find physical grounds and cosmoparticle
relevance

Baryosynthesis



GUT

$\Delta B \neq 0$
Baroleptonic gauge fields, X $\leftrightarrow P \rightarrow e^+ \pi^0$ decay

$$n_X = n_{\bar{X}} \quad X \rightarrow q\bar{q} \quad z \\ \rightarrow \bar{q}l \quad z - z \\ \bar{X} \rightarrow \bar{q}\bar{q}_c \quad \bar{z} \\ \rightarrow q_c l \quad 1 - \bar{z} \quad \left. \right\}$$

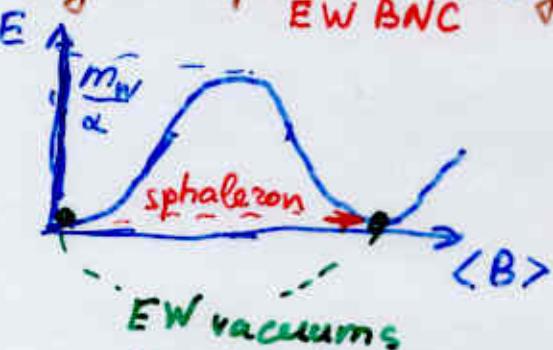
Owing to CP violation $z \neq \bar{z}$

$$\frac{\Delta B}{8} = (z - \bar{z}) \frac{n_X}{8}$$

SUSY

squarks, baryons with $s=0$, $\langle n_{sq} \rangle \neq 0$ squark condensate

High Temperature Baryon Nonconservation in Electroweak Theory
EW BNC



At $T \gg m_W$ sphalerons are in the equilibrium leading to $B + L \rightarrow 0$.

$\frac{\Delta B}{8}$ with $(B-L)=0$ can be washed down

No direct relation to proton decay, but maybe (?) related to strong BNC at high energies. High temperature BNC takes place even in the standard model, but respective Baryosynthesis seem to need smth beyond it.

$(\Delta L=2)$ and EW BNC

Physics of Majorana neutrino mass

$$\frac{v_L m_\nu \bar{\nu}_R}{\Delta L=2}$$

Out of equilibrium $\Delta L=2$ processes together with $(B+L) \rightarrow 0$ due to EW BNC can induce the observed baryon asymmetry of the Universe

Different physical models for inflation and baryosynthesis are not, in fact, alternative. They follow from different physical grounds and should be treated (in general) independently

Inflation

Strong 1st order phase transition



GUT Higgs potentials
Hierarchy of GUT sym. Breaking

New or chaotic inflation scenario



refined GUT Higgs potentials
SUSY, R^2 gravity

extended inflation



Kaluza-Klein models or
superstrings

Baryosynthesis

(CP violation in
non equilibrium

$$x \xrightarrow{q\bar{q}} \frac{q\bar{q}}{q} e \xrightarrow{\Delta B} \frac{\Delta B}{q}$$



GUT B-nonconservation

(sq) condensate with
 $\langle B \rangle \neq 0$



SUSY GUT models

$$(B+L) \neq 0$$



extended SM
high T B nonconservation in
EW model

$$\Delta L \neq 0$$



Majorana neutrino mass
physics

They should be selfconsistently considered together with the physics of nonbaryonic dark matter.

Hidden parameters of the modern Cosmology

The parameters ρ_{tot} , ρ_e , ρ_{dm}
finding physical basis in
inflation

$$\rho_{\text{tot}} = \rho_{\text{crit}} \quad (\text{simplest case})$$

baryosynthesis

$$\rho_e = f \left(\frac{\Delta_B}{\epsilon} \right)$$

and nonbaryonic dark matter

$$\rho_{\text{d.m.}} = m_{\text{d.m.}} n_{\text{d.m.}}$$

are accompanied by
additional (HIDDEN)
parameters ρ_{PBH} , $\rho_{\bar{e}}$, ...

In particle physics ($\hbar = c = 1$)

metastable particles : $\tau \gg \frac{1}{m}$

In Cosmology, to be of cosmological significance

$$\tau \gg t(T \sim m) = \frac{m_{pe}}{m} \cdot \frac{L}{m}$$

Following Noether's theorem :

Strict symmetry



Conserved charge

The lightest particle, possessing strictly conserved quantum number should be absolutely stable.

So, conservation of electron charge \Rightarrow stability of electron

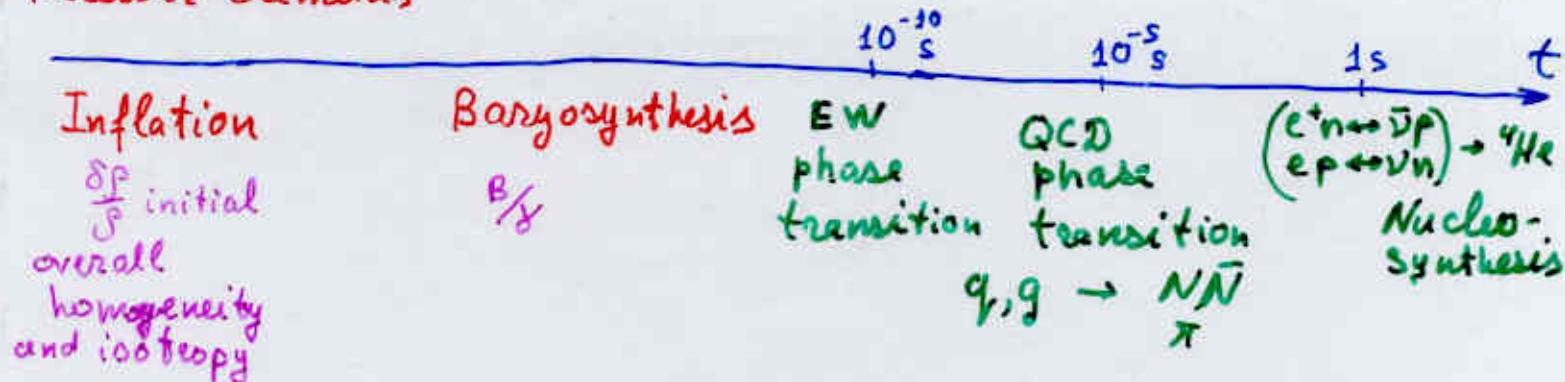
baryon charge (approximate) conservation
(meta) stability of proton

New strict (approximate) particle symmetries lead to new (meta)stable particle.

In this way cosmology probes the most fundamental structure of particle theory.

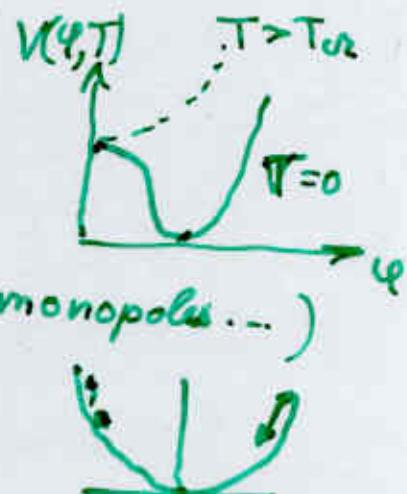
Very early Universe

Necessary elements



Possible processes

Spontaneous breaking of symmetry at $T < T_{cr}$



Topological defects (strings, walls, monopoles...)

Early dust-like stages

Post-inflationary stage $\langle p \rangle = 0$
 $\varphi = \varphi_0 \cos m t$ $t \gg \frac{1}{m}$

Freezing of metastable particles

$T \gg m$ - equilibrium

$T < m$ $n_m/n_2 \propto \exp(-M/T) \xrightarrow{\text{freezing}} \exp(-m/T_f)$

freezing $\tau(m \bar{m} \rightarrow \dots) > t$

$$n_m / (\bar{n}_{ann} v) t < 1$$

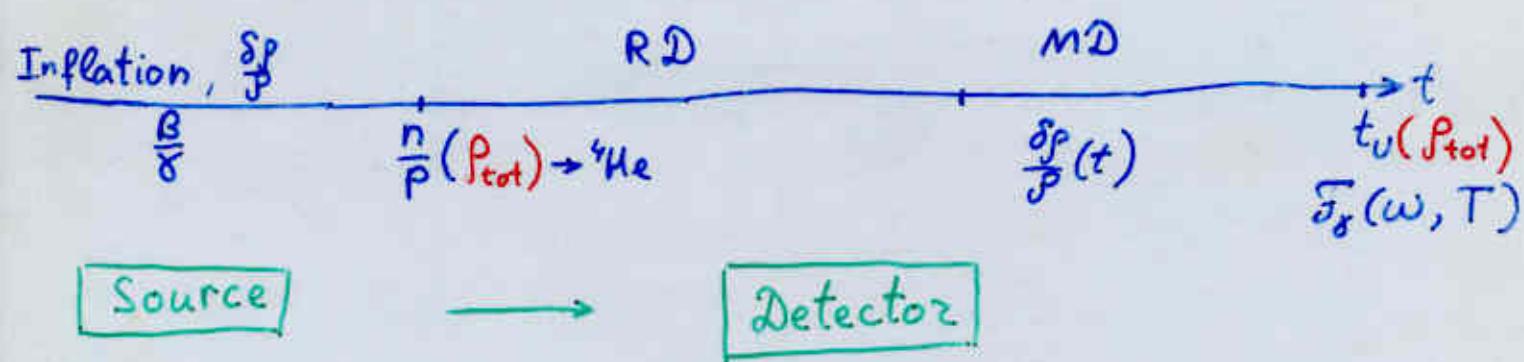
$$v = \frac{n_m}{n_2} \Big|_{\text{frozen}} \sim \frac{1}{m \cdot \bar{n}_{ann} \cdot m_{pe}}$$

At $v_m > T$

$$\rho_m = m \cdot v \cdot n_g > \rho_2 \sim T \cdot n_g$$

m -domination in the Universe until $t \approx$

Gedanken experiment



Astrophysical data as the sample
of experimental data from this gedanken
experiment.

Matter - structure of spatial distribution
- chemical composition

Thermal background - spectrum
- isotropy

Nonthermal background - X-ray, γ -ray, UV, IR, optical, radio

Cosmic rays - e^+e^- , N , nuclei

galactic magnetic fields antiprotons, antinuclei?

Dark matter fluxes

Magnetic monopole fluxes

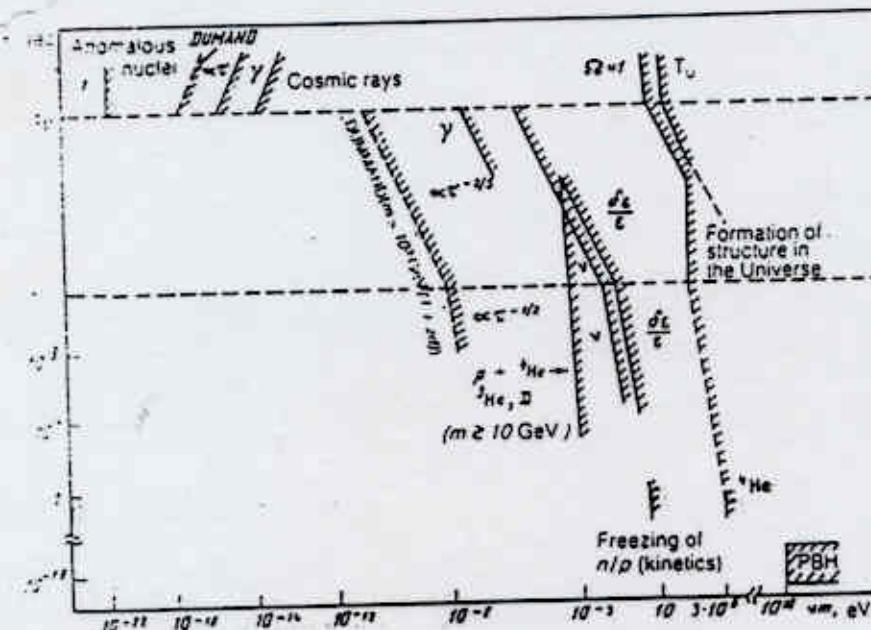
ν -background

Cosmic axions

Different properties
of hypothetical particles
and fields cause
different effects in
these data

Cosmoarcheology

Astrophysical data as experimental sample
from the Gedanken experiment with
cosmological consequences of new physics



Integral methods

t_U - age of the Universe, constraining the total density of the modern Universe

γ - primordial light element abundances constrain the total (^4He , ^2H) and baryon (D) densities in the period of SBBN

$(\frac{\Delta T}{T})_e$ - total and baryon densities in the period of recombination

MD stage in the period of LSS formation

Differential methods

$(\frac{\Delta T}{T})_\nu \rightarrow \frac{\delta \epsilon_\nu}{\epsilon_\nu}$ electromagnetic calorimeter $\delta \epsilon_\nu$

Inequilibrium cosmological nucleosynthesis $F(\rho, \langle \epsilon \rangle)$
 $\eta_P; f_i(P)$ $\langle \epsilon \rangle \gg T$

Cosmological nonthermal cosmic ray, γ , $\nu \bar{\nu}$ backgrounds

COSMOARCHEOLOGY

We can put constraints on hypothetical effects at different stages of cosmological evolution from UNIVERSAL TOOLS

1. Age of the Universe $\rho_{tot} \leq \rho_{max}$ for any form of matter
2. Primordial ^4He abundance $\rho_{tot} \leq \rho_{max}$ in the modern Universe
3. Formation of Large Scale Structure $\rho_{tot} \leq \rho_{max}$ for any form of matter
Excludes radiation dominancy in the period of large scale structure formation. $\rho_{tot} \leq \rho_{max}$ at $10^4 \leq t \leq 10^{16}$ s
4. Absence of primordial Black holes and effects of their evaporation
Probes early dust like stages and inhomogeneity of early Universe, starting from $t \sim 10^{-28}$ s

5. Spectrum of Black Body Radiation Background
Absence of $\Delta T_f(v)$ distortions constrain electromagnetic energy release, starting from $z > 10^8$
6. Nonthermal electromagnetic cosmic backgrounds
Probes sources of inequilibrium particles at $z > 10^2$
7. $D, ^3\text{He}, ^6\text{Li}, ^7\text{Be}, \dots$ abundances probing inequilibrium cosmological nucleosynthesis at $t > 10^5$ s
8. Cosmic rays, cosmic $\bar{\rho}$ and γ fluxes, sensitive to hypothetical sources of inequilibrium particles at $z < 10^3$
[M.Yu. Kh. se Evgenyi Sedelnikov, ~~GTRC~~ Seminar]
9. Cosmic High Energy Neutrino Background, probing hypothetical sources of superhigh energy particles at $t > 10^5$ s

Inhomogeneous Baryosynthesis

CP violation in the standard model of EW interaction is ascribed to KM mass matrix.

But

$$m_f = h_f \langle \varphi \rangle = 0 \text{ at } T > T_{EW}$$

when $\langle \varphi \rangle = 0$

Spontaneous CP violation may lead to spatial variation of CP violating phase $\chi(x)$

$$\text{GUT: } \frac{\Delta B}{\delta} \propto \text{Im} \chi(x) \Rightarrow \frac{\Delta B(x)}{\delta}$$

$$\text{SUSY: } V(B) \Rightarrow n_{sq} \rightarrow \Delta n_B$$

$$n_{sq}(x) \Rightarrow \Delta n_B(x)$$

Spatial variation of baryonic density of scalar quark condensate

$$\text{EW BNC} + (\Delta L=2):$$

axions are CP violating phase in the period

$$T_{PQ} > T > \Lambda \sim 1 \text{ GeV}$$

$$\Delta L(x) \Rightarrow \Delta B(x)$$

Coexistence of $\varphi_0(x) = \text{const}$ and $\varphi_1(x)$ can lead to local antibaryosynthesis in the Universe with global dominancy of baryons.

Antibaryon domains are in this case probes for the mechanisms of baryosynthesis and CP violation at high temperatures

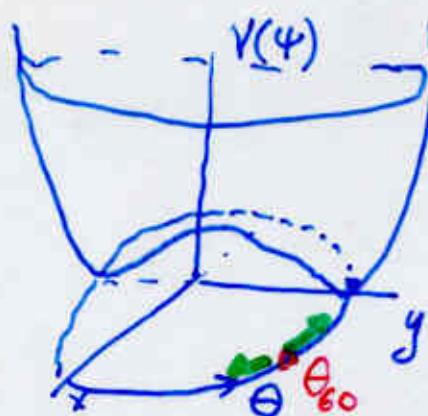
Inflationary Model with Spontaneous Baryogenesis

(KHLOPOV, RUBIN, SAKHAROV, Phys. Rev D, 2000)

$$\mathcal{L} = h \theta \bar{\psi} Q$$

$$\Delta L \neq 0$$

$$\Delta B \neq 0$$

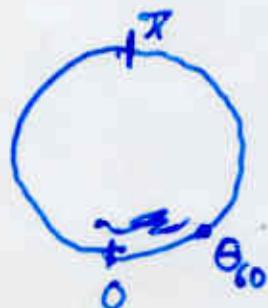


$$\Phi = f e^{i\theta}$$

$N = 60 \rightarrow e\text{-folding}$, corresponding to the modern horizon.

$$\theta_{60} \rightarrow \ell_h(t_0)$$

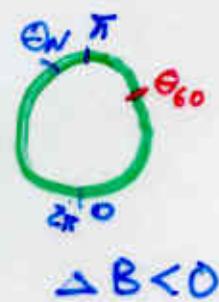
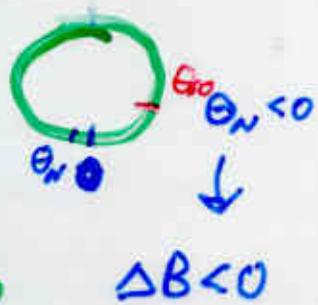
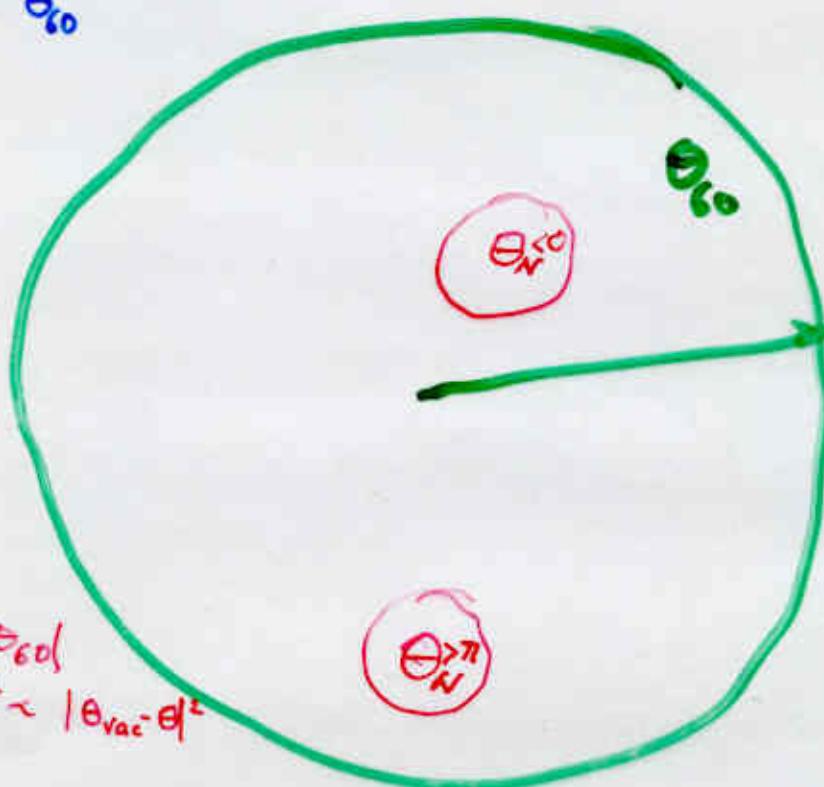
$$\Delta \theta \sim \frac{M_{\text{infl}}}{2\pi f}$$



$$\theta_N < 0, \text{ or } \theta_N > \pi$$

$$\begin{aligned}\theta_N &< 0 \\ |\theta_N| &< |\theta_{60}| \\ \Delta B &\propto \Theta^3 \\ p_{\bar{e}} &\ll p_e\end{aligned}$$

$$\begin{aligned}\theta_N &> \pi \\ |2\pi - \theta_N| &> |\theta_{60}| \\ \Delta B &\propto \Theta^2 \sim |\theta_{\text{vac}} - \theta|^2 \\ p_{\bar{e}} &> p_e\end{aligned}$$



Antimatter in the Baryon asymmetrical Universe

Diffused Antiworld

Local anti-Baryosynthesis in baryon asymmetrical Universe may lead to antimatter domains with $\Sigma_{\bar{B}} \ll \Sigma_B$

SBB Nucleosynthesis results in this case in "unnatural" abundances of antim nuclei

$$\bar{H} : \bar{D} : {}^3\bar{\text{He}} : {}^4\bar{\text{He}} \neq H : D : {}^3\text{He} : {}^4\text{He}$$

In diffused antiworld \bar{D} and ${}^3\bar{\text{He}}$ may be most abundant (after \bar{H}) and serve as the signature of antimatter.

At $\Sigma_{\bar{B}} \ll 1$, $\Sigma_{\bar{B}} \ll \Sigma_B$ antigalaxy formation may not take place (At $\Sigma_{\bar{B}} < 10^{-5}$ recombination may be strongly suppressed and the years scale of ionized antimatter is of the order of cosmological horizon)

Statistical analysis of antimatter domain distribution is needed.

Preliminary estimates show, that domains as small as 1 kpc can survive and even be present in the Galaxy.

Antimatter globular cluster in our Galaxy – the probe for the matter origin

M.Yu. Khlopov *

Center for CosmoParticle Physics "Cosmion", Miusskaya Pl.4, 125047 Moscow, Russia

Institute of Applied Mathematics, Miusskaya Pl.4, 125047 Moscow, Russia

*Moscow Engineering Physics Institute (Technical University), Kashirskoe Sh.31, 115409 Moscow,
Russia*

Abstract

The existence of a globular cluster of antimatter stars in our Galaxy is shown to be a probable signature of the mechanism of inhomogeneous baryosynthesis. The observed gamma ray flux puts the constraint on the total mass of such anti-cluster. The expected signatures in cosmic ray experiments are discussed.

1 Introduction

Baryon asymmetry of the Universe is considered as the one of the most important features of the modern cosmology. It is commonly based on the statement that no macroscopic amount of antimatter is present around us at least on the scale of the local supercluster of galaxies. The statement is generally supported by the negative results of the direct searches for antimatter in the vicinity of the Solar system and by the severe constraint on matter-antimatter annihilation, following both from the observed gamma background and from the analysis of the influence of annihilation in the early Universe on the spectrum of relic radiation and on the light element abundance. However, the both types of evidences, definitely excluding the equal amounts of matter and antimatter around us, do not exclude the principal possibility of the existence of macroscopic amount of antimatter, putting only upper limit on its average density and possible distribution.

The idea of antimatter probing the origin of the matter in baryon asymmetrical Universe was first put forward in Ref. [1] (see also Ref. [2]). It was shown (see for review Ref. [3]) that practically all the existing mechanisms of baryosynthesis Ref. [4] can under some condition lead to antibaryon excess. In the other words the idea of inhomogeneous baryosynthesis was put forward in Ref. [1] and in Ref. [2]. From this viewpoint antimatter represents the high amplitude

*email: mkhlopov@orc.ru

Forms of Antimatter Domains

$$\Omega_{\bar{B}} = \frac{\langle \rho_{\bar{B}} \rangle_{\text{inside domain}}}{\rho_{\text{cr}}}$$

$\Omega_{\bar{B}} < 10^{-4}$ Diffused Antiworld

BBN: $10^{-5} < \Omega_{\bar{B}} < 10^{-4}$ ${}^3\bar{\text{He}} < \bar{D}$; \bar{H}, \bar{D}

$$\Omega_{\bar{B}} < 10^{-5} \quad \bar{H}$$

No recombination

$\rho_{\bar{B}} < \rho_{\text{c}}$ inside domain at $\Omega_{\bar{B}} < 10^{-4}$

No gravitational instability

Such domains can not be in the superclusters of galaxies and should be in voids

$$10 \text{ Mpc} < R < 1 \pm 10 \text{ Mpc}$$

γ -sources from voids

$\Omega_{\bar{B}} > 10^{-4}$ after recombination and gravitational instability

can lead to antistar formation

Globular cluster of antistars
in the Galaxy

(1)

$$M \sim 10^4 \text{ to } 10^6 M_\odot$$

(2)

low $\rho_{\bar{B}}$ gas

(3)

low $\rho_{\bar{B}}$ of gas

annihilation

Can survive

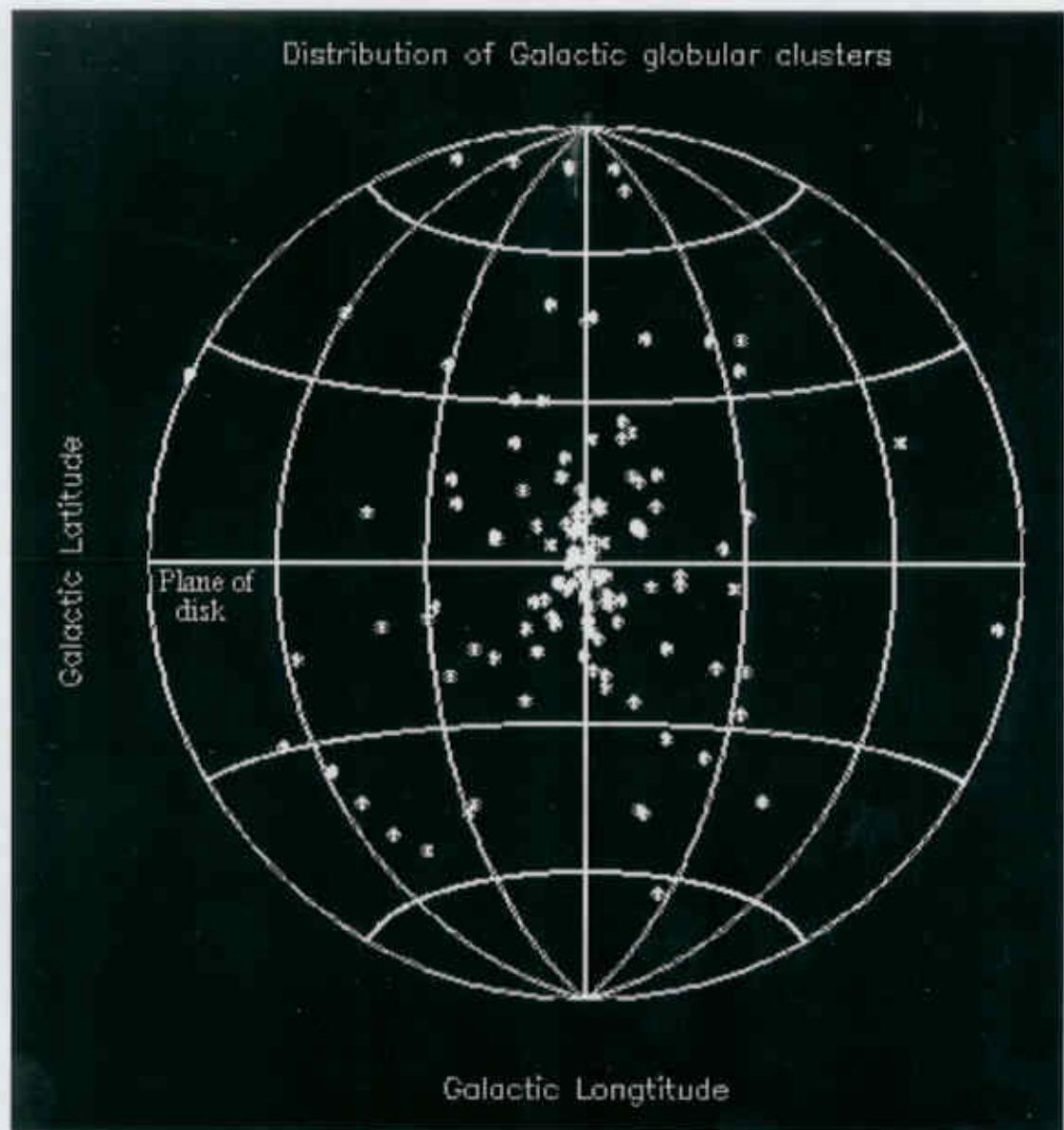
spherical component

No antistars in disc

Acceleration in multiple annihilation

of antinuclei: $\bar{A} + p \rightarrow (\bar{A}-1)$

Globular clusters



Globular cluster population in our Galaxy consists of 147 confirmed globular clusters. In the spherical component of Galaxy the globular clusters move with velocity 150km/s – 170km/s.

If the high – density antimatter region was present in our Galaxy, it could survive in the form of globular cluster at large galactocentric distance.

Antimatter globular cluster

$$M_{\text{surv}} \leq M < M_{\text{max}}$$

$$M_{\text{surv}} \sim \rho_{\bar{e}} l_s^3 \sim 10^3 M_\odot$$

for $\rho_{\bar{e}} \sim \rho_e$

$$M_{\text{max}} \leq 10^5 M_\odot$$

from $\mathcal{F}_\delta^{\text{obs}} \approx \mathcal{F}_\delta(\bar{\rho}_{\text{ann}})$


$$\ast \bar{\rho} \quad \text{if } r > 5 \text{ kpc}$$
$$n_{\text{gas}} = 5 \cdot 10^{-4} \text{ cm}^{-3}$$

$$\mathcal{F}_\delta(\bar{\rho}_{\text{ann}}) \propto \dot{\bar{\rho}}$$

Elliptical galaxies

$$\dot{m} = \frac{0.1 M_\odot}{10'' M_\odot}$$

Antimatter pollution

- The integral effect of antimatter globular clusters may be estimated by the analysis of antimatter pollution of the galaxy by that globular clusters (K.M.Belotzky, Yu.A.Golubkov, M.Yu.Khlopov, R.V.Konoplich, A.S.Sakharov 1998)
- There are two main mechanisms of antimatter loss by the globular clusters

I. The stationary mass loss by antimatter stars in the form of stellar wind

$$\dot{M} \approx 10^{-12} M_{SUN} - \text{per solar mass per year}$$

II. The antimatter supernova explosion $\dot{M} \approx 10^{-13} M_{SUN} - \text{per solar mass per year.}$

The model of galactic annihilation

- The $\bar{p}p$ annihilation cross section

$$\sigma_{ann}(P < 300 \text{ MeV}/c) = \frac{2\pi c \alpha}{v_c^2 (1 - \exp(-2\pi c \alpha / v_c))} \cdot (160 \text{ mb})$$

- The model of halo

$$n_H(z) = n_H^{halo} + \Delta_H(z); \quad \Delta_H(z) = \frac{n_H^{disk}}{1 + (z/D)^2}; \quad n_H^{halo} = 5 \cdot 10^{-4} \text{ cm}^{-3}; \\ n_H^{disk} = 1 \text{ cm}^{-3}; \quad D = 100 \text{ pc};$$

90% of the halo mass is a non - baryonic dark matter

- For the \bar{p} with velocities $\approx 10^3 \text{ km/s}$ (stellar wind) the confinement time in the halo, starting from the distances $z \approx 2 \text{ kpc}$ is less than annihilation time ("two - zone" leaky box model)
In the gaseous disk the situation is opposite
- The \bar{p} are collecting in the halo during the confinement time $\approx 5 \cdot 10^8 \text{ yrs}$ increasing the gamma flux
- During the large confinement time the \bar{p} are being spread over the halo with constant number density not dependent on the position of antistar cluster and under the usual acceleration mechanisms in the halo their spectrum comes to the stationary form

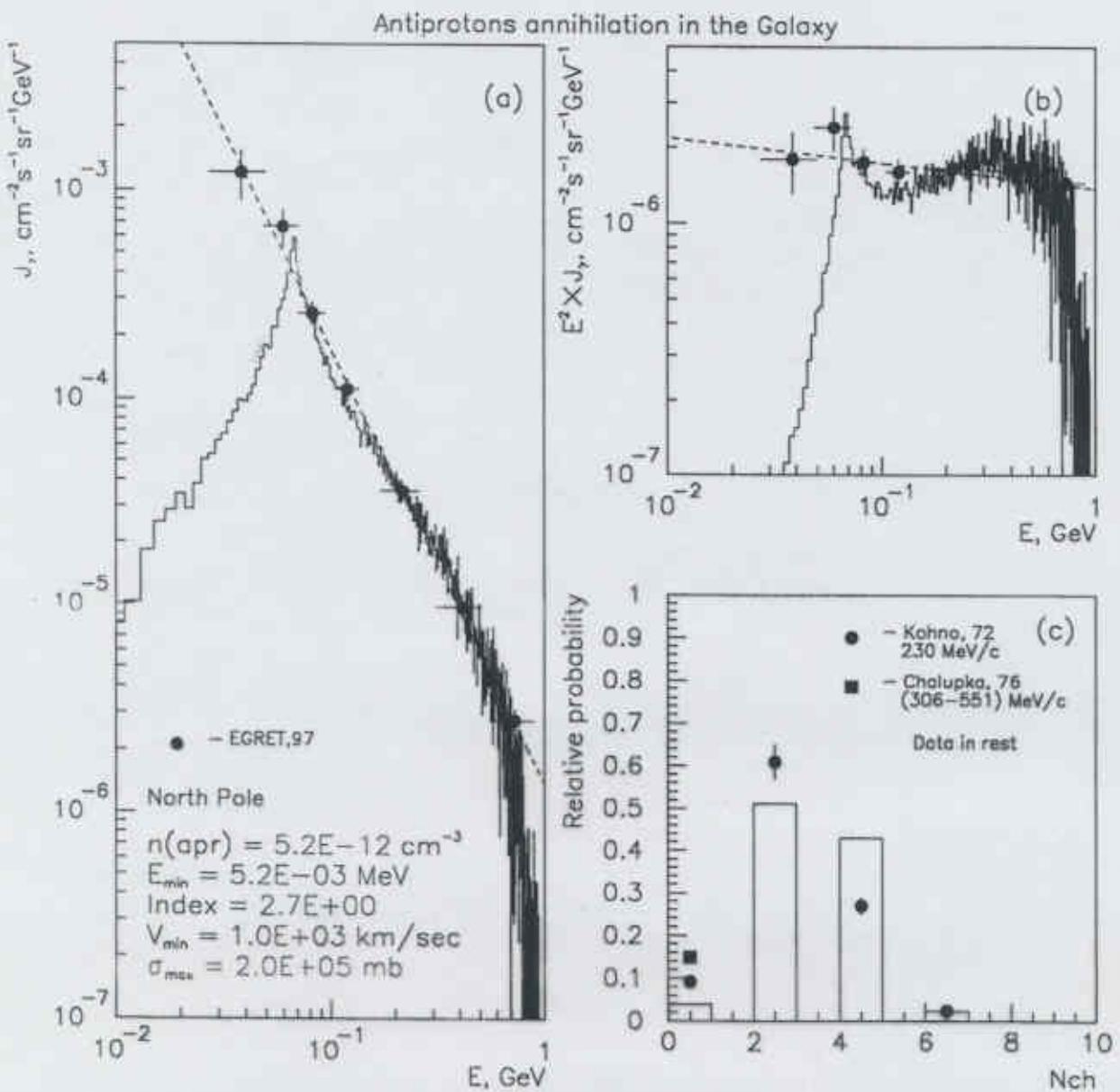


Figure 3: Comparison of the calculated differential fluxes of γ quanta from $\bar{p}p$ annihilation with experimental data *EGRET* [20] on diffuse gamma background (a,b). There is also shown the comparison of the charged multiplicity distribution in the annihilation model described in the text with the existent experimental data (c).

The uncertainty

The actual distribution of magnetic field in our Galaxy.

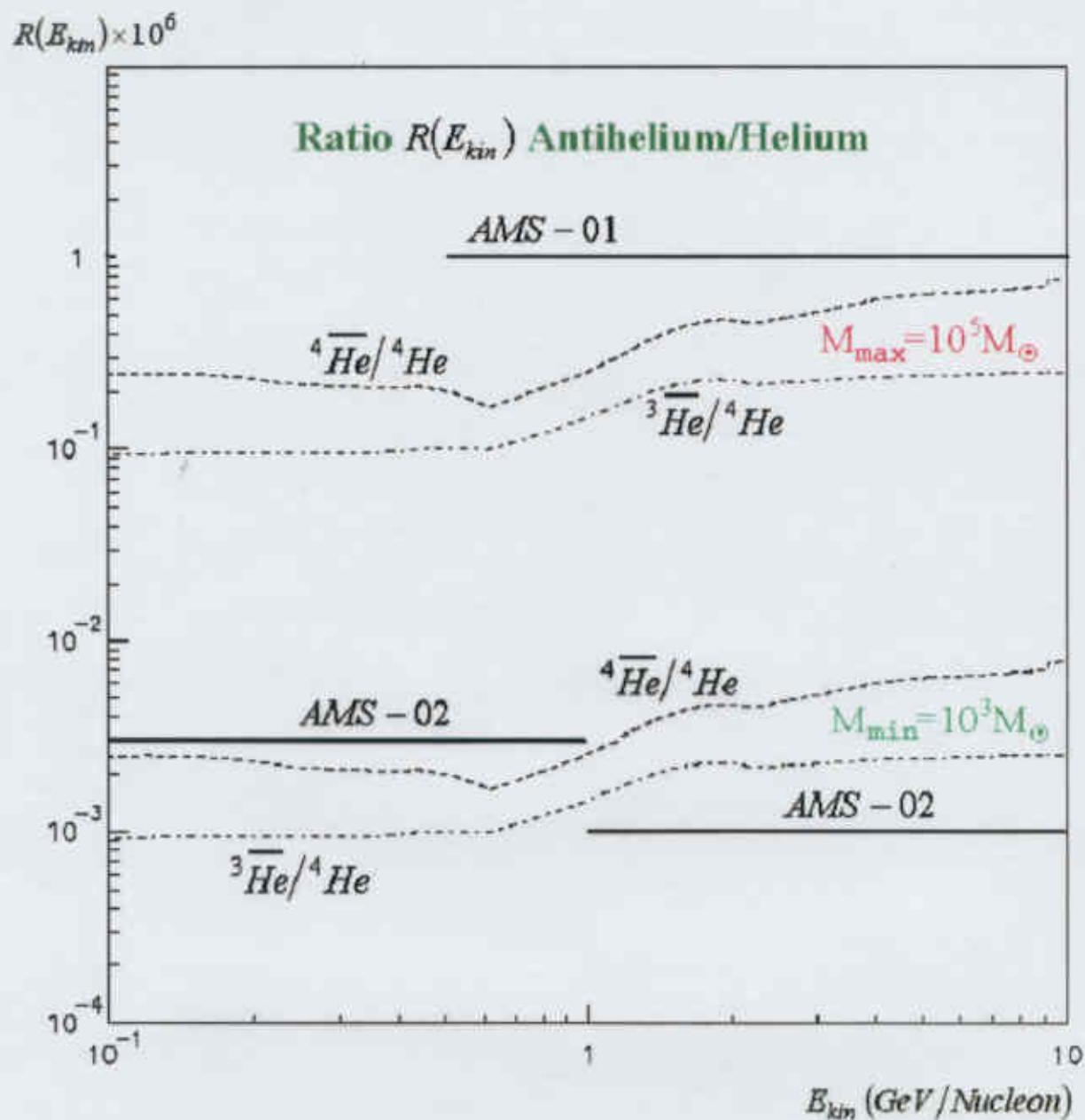
The mechanism of cosmic-ray acceleration

The relative contribution of disc and halo particles into the cosmic-ray spectrum

The acceleration of matter and antimatter cosmic rays are similar

The contribution of antinuclei into the cosmic-ray fluxes is proportional to the mass ratio of globular cluster and Galaxy.

Anti-helium flux



Khlopov

Cosmoparticle Physics



Maxim Yu Khlopov

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6 МЕЖДУНАРОДНАЯ КОНФЕРЕНЦИЯ
ПО КОСМОМИКРОФИЗИКЕ

(КОСМИОН-2003)

Москва - С.Петербург – Париж-Медон, 1-22 июня, 2003

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ON COSMOPARTICLE PHYSICS
(COSMION-2003)

Moscow - St.Petersburg – Paris-Meudon, 1-22 June, 2003