

Cosmological Matter- Antimatter Asymmetry.

Blois, 16-22 June, 2002

Prediction (Dirac) and discovery (Anderson) of antimatter are among the greatest scientific achievements of XX century.

But:

new discoveries \Rightarrow new problems

I. Why the Universe is strongly
($\sim 100\%$) matter dominated
at least in our neighbourhood?

- answered by Sakharov in 1967

II. Are there astronomically large
domains or objects of antimatter?

Unanswered yet. Challenge
for XXI century (or III Millenium
?)

Dirac, 1934: "... quite possible ...

these stars being built up mainly
of positrons and negative protons.

In fact there may be half
of the stars of each kind.

Observations: $p, e^-,$ nuclei
in our part of the
Universe

a little \bar{p}, e^+ in cosmic rays
(of secondary origin)

Not a single antinuclei yet
observed.

Theory: predicts a lot of
relic antineutrinos, $\sim 55/cm^3$ of
each flavor

B4

Data analysis: γ ray
background, ~ 100 MeV
nearest antigalaxy at

(Stecker, Puget, 1972) $\bar{l} > 10$ Mpc
(Steigman, 1978 - review)

Cohen, De Rujula, Glashow, 1998:

$$\bar{l} > \text{Gpc}$$

if $M : \bar{M} = 50 : 50$ and

adiabatic perturbations

Angular fluctuations of CMBR - bound
depends on scale; no limits
for small scales.

Despite ~~all~~ ^{limits} above:

2007
BS

existence of astronomically large objects of antimatter (gas clouds, antistars, ...) is not excluded and even natural. They may be as close as halo of our galaxy and rich of heavy antinuclei ($\bar{C}, \bar{N}, \bar{O}, \dots, \bar{Fe}?$)

Very interesting to search for them.
Why?

1. Evident (no other motivation is necessary)
2. Understanding of C, CP - breaking in cosmology
3. Proof (or evidence) of inflation.

Local (?) excess of baryons over anti baryons is!

Elbas
BG

$$\left[\beta = \frac{N_B - N_{\bar{B}}}{N_\gamma} \approx 6 \cdot 10^{-10} \quad (N_\gamma \approx 400/\text{cm}^3) \right]$$

Now $N_B \gg N_{\bar{B}}$;

early, at $T \approx 100 \text{ MeV}$, $N_{\bar{B}} \approx N_B (1 - \beta)$

$\beta =$ baryon asymmetry of the universe
(BAU)

5* - questions for XXI century:

1. $\beta = \text{const}$ or may $\beta = \beta(x)$?
2. $\beta > 0$ or somewhere $\beta < 0$?
3. $B_{\text{tot}} = \int \beta d^3x \gtrsim 0$ or $B = 0$?

global charge
symmetry
4. If $\beta = \beta(x)$ what is the characteristic scale l_B ?
5. What kind of anti-objects may exist?

Theory allows anything.

Data possibly exclude $B = 0$
inside horizon, $l_{\text{hor}} \approx 3 \text{ Gpc}$.

Scenarios of Baryogenesis

Pioneers:

A. Sakharov (67) \leftarrow nonconserved B ;

R. Omnès (69, 70) \leftarrow conserved B

\hookrightarrow spatial dynamical separation of M and \bar{M} ; is not known how to achieve at large scale

[crazy idea: \blacksquare higher dimensions, \bar{M} is pushed to another brane?] energy conservation?

modified Omnès, in our world looks as a realization of Sakharov's idea of non-conserved B

3 principles of baryogenesis:

1. Non-conservation of baryons

Weakest point in 1967:

"we exist, hence baryons are stable"

Theory: GUT, electroweak $\Rightarrow \Delta B \neq 0$

Inflationary cosmology is impossible if baryons are conserved.

Our universe is a strong experimental indication in favor of B-nonconservation.

Now: we exist, hence baryons are unstable.

2. Breaking of charge symmetry (C, CP ↓)

Experiment: $K_L \rightarrow 2\pi$, $K^0 \rightarrow \pi e^\pm \nu$
well established.

Theory: many models, do not know which is realized.

3. Deviation from thermal equilibrium

In equilibrium $f = (1 + e^{E/T})^{-1}$

and $f = \bar{f}$ if $m = \bar{m}$ (CPT)

Chemical potentials $\mu = \bar{\mu} = 0$ if B is not conserved

Theory: FRW - cosmology \Rightarrow nonstationary universe

) \hookrightarrow particles with $m \neq 0$ are always out of equilibrium.

) First order phase transitions during hot stage \Rightarrow strong deviation from equilibrium.

VB: none of the three above is obligatory.

Necessary conditions for creation of matter and antimatter domains

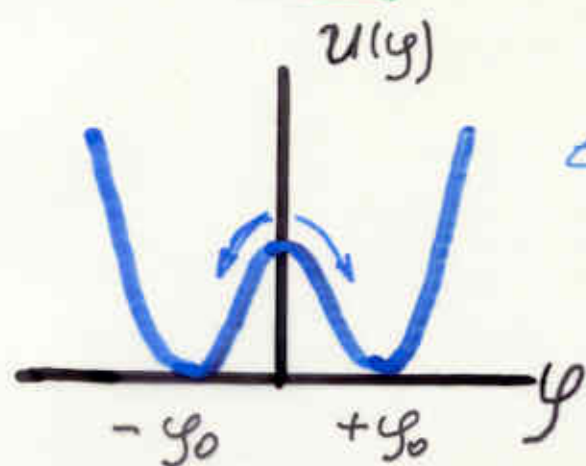
1. Different signs of C (CP) -
breaking in different regions
(Brown, Stecker, spont. CP breaking)

2. Exponential (inflationary) but
moderate blow-up of regions
with a definite sign of C, CP -
breaking.

(K. Sato, before inflation was
proposed)

Mechanisms of C, CP - breaking ³¹³ in cosmology.

1. Explicit — complex constants in Lagrangian
↳ $\beta = \text{const}$ (space point independent)
2. Spontaneous (T.D. Lee)



← complex field $\psi(x)$

50% of $y = +y_0$
50% of $y = -y_0$

Globally charge symmetric universe

Domain wall problem (Kobzarev, Okun, Zeldovich)

Necessary $l_B > l_{\text{horizon}}$ to solve domain wall problem or to destroy the walls

Even if walls were eliminated,

B and \bar{B} were in close contact and $l_B > l_{\text{CMB}}$
(CDRG)

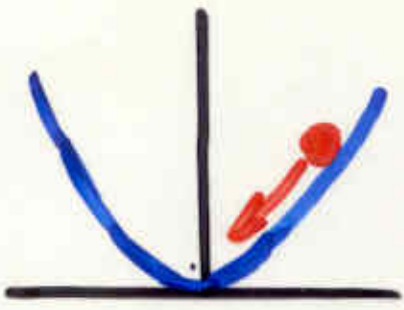
3. stochastic c, CP-breaking (by chaotic initial conditions); does not manifest in particle physics

E.g. $\psi(t, x) \neq 0$ initially

↳ complex field

[inflation $\langle \psi^2 \rangle \sim H^3 t$ chaotically]

If $\psi(t \rightarrow \infty) \rightarrow 0 \Rightarrow$ no c, CP breaking observable today,



no domain walls,
but CP-odd
amplitudes $\neq 0$
during baryogenesis.

Possibly both explicit and stochastic were operative.

↳ $\bar{M} : M < 1$
~~~~~



# Models of Baryogenesis

LEP/1  
B15

## 1. GUT - baryogenesis.

gauge bosons,  $m_x \sim 10^{16}$  GeV

$X \rightarrow qq, \bar{q}l$ ; B nonconserved

$$\frac{B(X \rightarrow qq)}{B(\bar{X} \rightarrow \bar{q}\bar{q})} \neq 1 \iff CP \text{ nonconservation}$$

Deviation from thermal equilibrium  
at  $T = m_x$ :

$H =$  Hubble  
parameter

$$\frac{H(T = m_x)}{\Gamma_x} \approx \frac{m_x}{10^{17} \text{ GeV}} \leftarrow \text{quite large}$$

Problem: probably universe was  
never at  $T \approx m_x$  and  
 $X$ -bosons never were  
abundant.

Naturally gives  $\beta = \text{const}$ ,  
no antimatter

## 2. Electro-weak baryogenesis

CERN  
B16

Fuzmin  
Rubakov  
Shaposhnikov

Characteristic temperatures:

$$T_{EW} = (\text{a few}) \times 100 \text{ GeV} - \text{TeV}$$

Normal, known, "low-energy" physics

$$\underline{SU(3) \times SU(2) \times U(1)}$$

Baryons do not conserve due to quantum chiral anomaly.

Thermal equilibrium is broken if EW phase transition is first order.

Problem: if  $m_{\text{Higgs}} > 100 \text{ GeV}$  phase transition is second order.

EW not good for creation of  $\bar{M}$   
(as well as GUT)



### 3. Baryo-through-lepto

317  
Fukugita  
Yanagita

(combination of 1 and 2).

Heavy Majorana neutrino  $\nu_M$   
(invented for explanation of small masses of ordinary  $\nu$ 's)

$$m_{\nu_M} \sim 10^{10} \text{ GeV}$$

Decays of  $\nu_M$  naturally do not conserve  $L$ ,  
produce lepton asymmetry

EW processes convert  $L$  into  $B$   
in thermally equilibrium state.

B-thru-L is also bad for  $M$ .

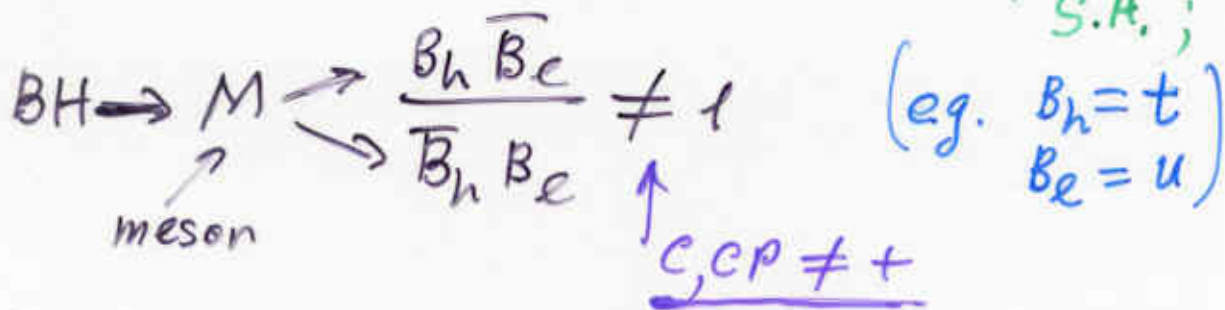
Problems: deviation from equilibrium of  
 $\nu_M$  at  $T \sim m_{\nu_M}$  is small;  
CP-breaking effects are probably  
also small and resulting  
baryon asymmetry may be too low (?)

People working on B-thru-L are more  
optimistic



#### 4. Black hole evaporation.

(Ya.B.Z.; BB  
S.K.; A8)



$B_h$  and  $\bar{B}_h$  are more efficiently back captured.

Baryons outside, anti baryons inside BH  
BH may possibly completely disappear.

Could give  $\beta \sim \beta_{obs}$  if  $\rho_{BH}(t=t_{decay}) \sim \rho_{tot}$

[similar model in particle physics  
 $p +$  (heavy stable antibaryons)]

Bad for antimatter.

Problems: unnaturalness  
mechanism of early BH formation?

CERN 11  
B13

## Intermediate conclusion.

### BAD NEWS

All simple models discussed above are bad for antimatter

simultaneously <sup>but</sup> they are not very good for creating the baryon asymmetry of the observed magnitude.

An introduction an additional complex scalar field  $y(t, x) \rightarrow$  stochastic CP-violation, may give rise to cosmic antimatter.

### GOOD NEWS

There are (at least) two known models which can explain the observed baryon asymmetry and naturally produce plenty of antimatter.



## 5. Spontaneous baryogenesis

(A. Cohen,  
J. Kaplan)SSB of  $U(1)_{B \text{ or } Y}$ :

$$\mathcal{L} = \lambda (|y|^2 - f^2)^2 + \bar{\Psi} (i \not{\partial} + m) \Psi + g \bar{\Psi} \Psi \Psi - |\partial \Psi|^2 \dots$$



$$y = f e^{i\theta}, \quad \Psi_B \rightarrow \Psi_B e^{i\theta}$$

baryonic current

$$\mathcal{L} \rightarrow f^2 (\partial \theta)^2 + \underbrace{i \partial_\mu \theta \cdot \bar{\Psi}_B \gamma^\mu \Psi_B}_{\text{baryonic current}} + \dots$$

Homogeneous  $\theta = \theta(t) \Rightarrow \dot{\theta} \sim \mu_B$ , i.e.

$$\dot{\theta} = \mu_B \text{ (chemical potential)}$$

$B \sim \dot{\theta}$  in thermal equil  
not exactly true

Still  $B \sim \dot{\theta}$  in thermal equilibrium and without C, CP-violation

determined by initial conditions, possibly by stochastic quantum fluctuations during inflation.

Sign of  $\dot{\theta}$  is arbitrary  $\rightarrow$  matter  
 $\rightarrow$  antimatter

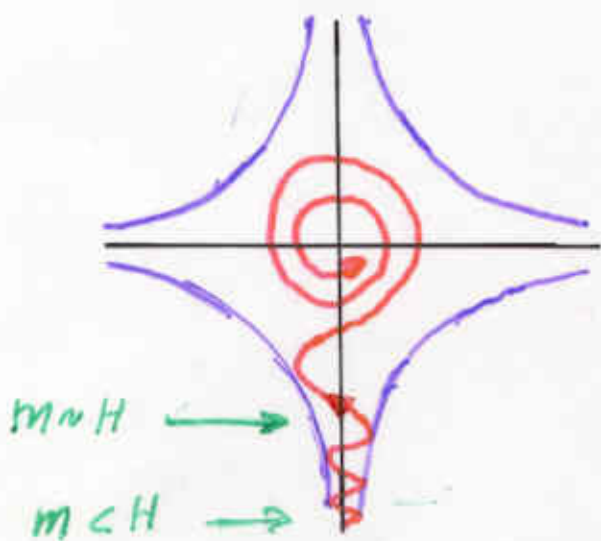


Necessary ingredients:

- scalar field,  $\chi$ , with  $B \neq 0$ ;
- flat directions in  $U(\chi)$ ;
- non-conservation of baryonic charge of  $\chi$ , i.e.  $u = u(\chi)$  but not  $u(|\chi|)$ .

Natural in susy models.

$$u(\chi) = \lambda |\chi|^4 (1 - \cos 4\theta), \quad \chi = |\chi| e^{i\theta}$$



$$g_{\text{valley}} = \frac{\pi}{2} \cdot n$$

$$\langle \chi^2 \rangle \sim \frac{H^3 t}{2\pi} \text{ at inflation}$$

$$\ddot{\chi} + 3H\dot{\chi} + u'(\chi) = 0 \leftarrow \text{Newton law}$$

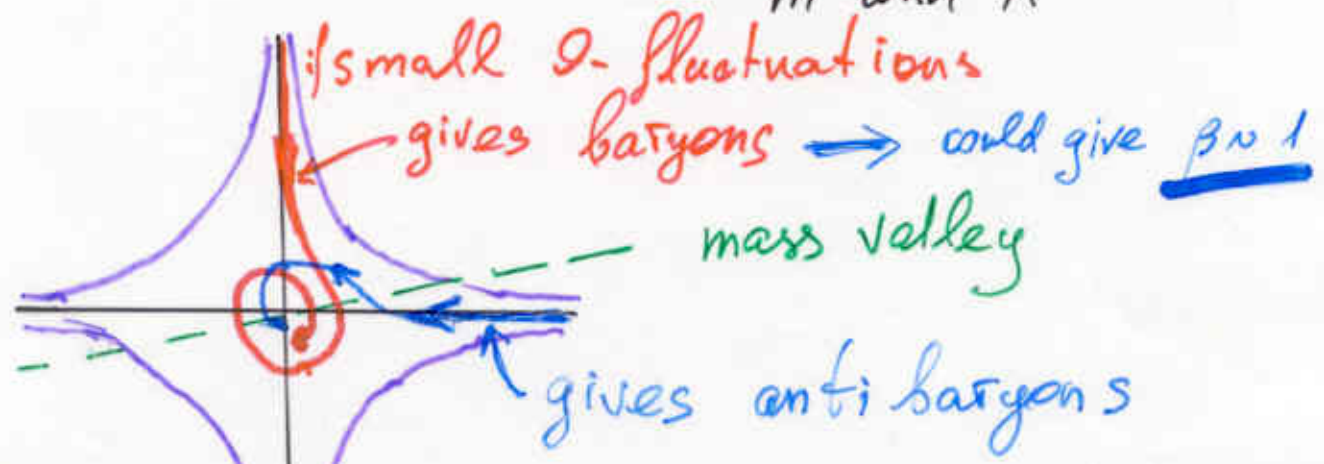
$$J_{\mu}^B(\chi) = \chi^* \overleftrightarrow{\partial}_{\mu} \chi, \quad \underline{B = \dot{\theta} |\chi|^2} \leftrightarrow \text{angular momentum}$$

$\beta$  is determined by chaotic initial conditions (at inflation); no CP is necessary; however the effect may be small

# Explicit CP-violation:

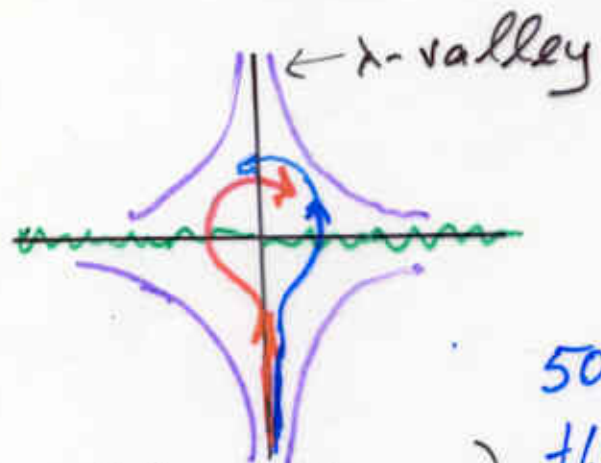
$$u(x) = \lambda |x|^4 (1 - \cos 4\theta) + m^2 |x|^2 [1 - \cos(2\theta + 2\alpha)]$$

↑  
relative phase of  $m$  and  $\lambda$



$B_{tot} = 0$  but  $L_B > L_{horizon}$   
(by inflation)

# CP-conserved model:



$\alpha = 0$

$m$  and  $\lambda$  - valleys

50% and 50% ← as in the case of spontaneous CP-violation but no domain wall problem

$\alpha \neq 0$ , but small  
90% B and 10%  $\bar{B}$   
or anything one wants



# Evolution of high density anti baryonic bubbles

BZ

1. Form Black Holes if  $\left(\frac{\delta\rho}{\rho}\right)_{\text{horizon}} \sim 1$ .

$$\frac{\delta\rho}{\rho} \sim \frac{m_B}{3T} \quad \frac{n_B}{n_\gamma} = \frac{1}{3}\beta \quad \frac{m}{T}$$

$$M_{\text{horizon}} = \frac{8\pi}{3} \cdot \gamma t^3 \cdot \rho = 5 \cdot 10^4 M_\odot \left(\frac{t}{\text{sec}}\right)$$

$\overline{BH}$  would possess atmosphere of antimatter

2. Before hydrogen recombination diffusion is slow and atmosphere might survive

3. Recombination <sup>in</sup> high  $\beta$  regions occurs earlier  $T = 1-10 \text{ eV}$

Hence atmosphere would not be stripped by CMBR at the onset of structure formation.

4. Annihilations could heat the plasma and create secondary ionization  
How strong?



If  $\chi$  is coupled to inflaton

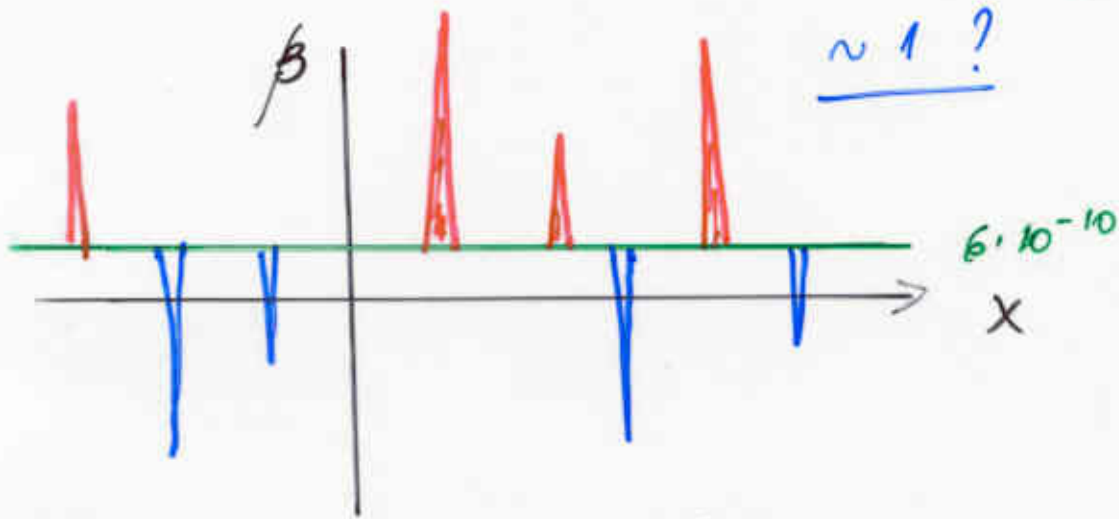
$$\lambda(\chi^2 + \chi^{\dagger 2})(\phi - \phi_1)^2$$

(most general renormalizable coupling)



then:

AD + J. Silk



$$\frac{dN}{dM} \sim \exp\left\{-c \ln^2 \frac{M}{M_0}\right\}$$

↑ primordial BH

with  $M = 1 - 10^9 M_\odot$

plus uncollapsed clouds of antimatter, antistars, etc ... Quasars?

For large  $\beta$  BBN efficiently produces heavier elements (does not stop on  $Li^7$ ) and  $\bar{0}$ ,  $\bar{c}$ ,  $\bar{N}$  may be observed! Fe?

1. Complex scalar field  $\varphi$
2. Inflation created condensate of  $\varphi$ :  $\langle \varphi \rangle = \varphi(x)$   
or bubble from  $\mathbb{Z}$  order p.t.

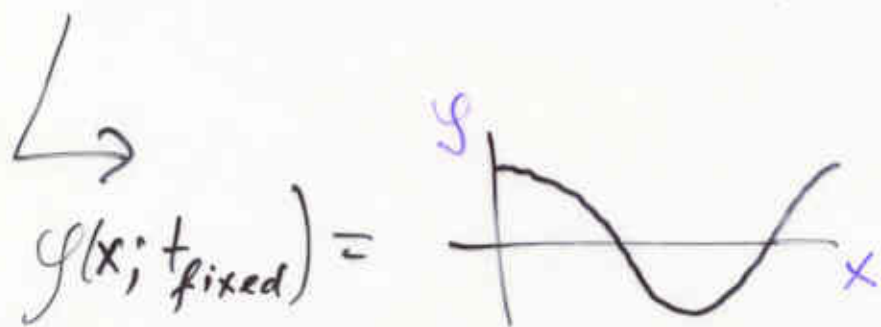


$$3. u(\varphi) = \lambda |\varphi|^4$$

↳ spatial fluctuations of baryon asymmetry at astronomically large scale

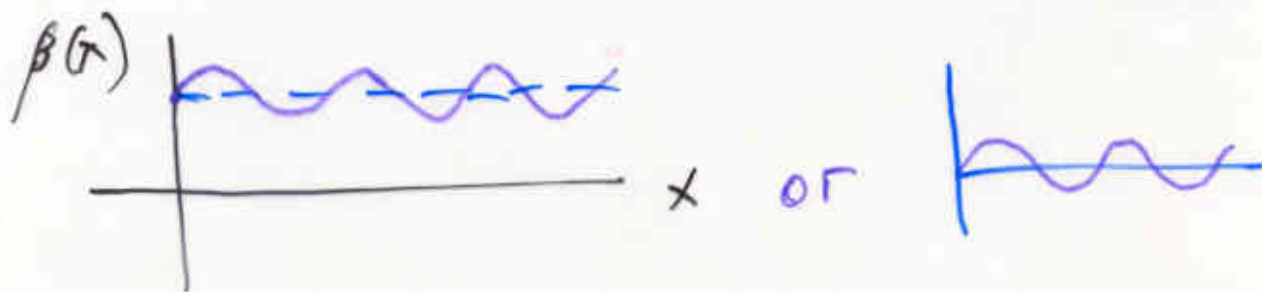
End of inflation:  $\varphi = \varphi(t; x)$

$$\omega \sim \sqrt{\lambda} \varphi_i(x)$$



$$\lambda \sim 100 \text{ Mpc} \text{ (could be)}$$

If  $\omega < \Gamma_{BS} \Rightarrow \beta \sim \varphi(x; t_{BS})$



or explicit CP breaking

Large scale excluded by CMB only very large  $\lambda \gg \sim \text{Gpc}$  is allowed

GR - bound:  $L_B \gtrsim c \tau$

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B25

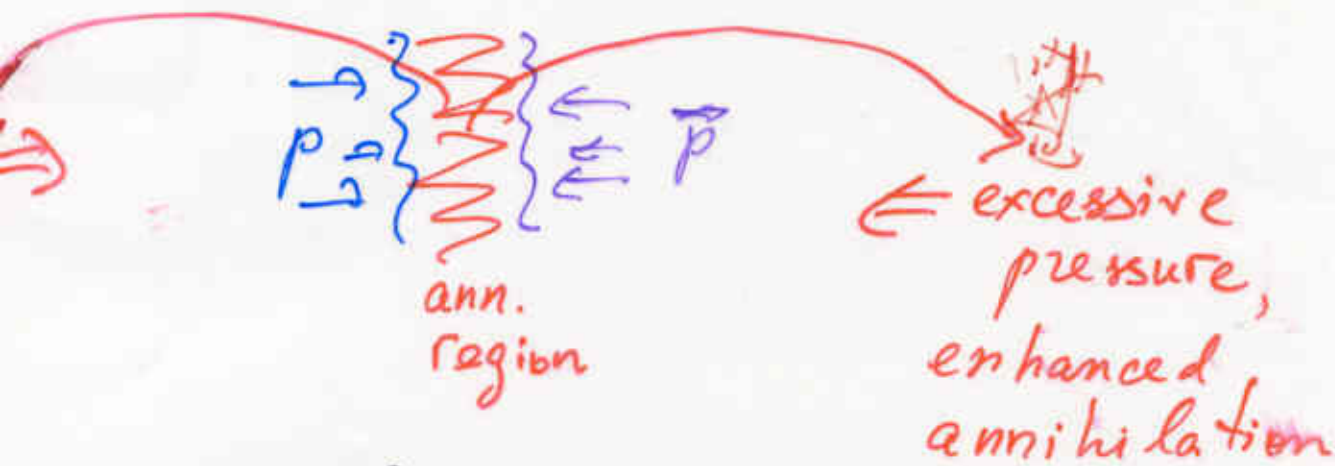
1. If  $B$  and  $\bar{B}$  are separated by

$\Delta L > 20 \text{ Mpc} \Rightarrow$  too large  $\delta T / T$

$\uparrow$  this small scale is already accessible to measurements.

2.  $\Delta L < 20 \text{ Mpc} \Rightarrow p, \bar{p}$  diffusion

brings them into contact



Valid for  $M \approx \bar{M}$

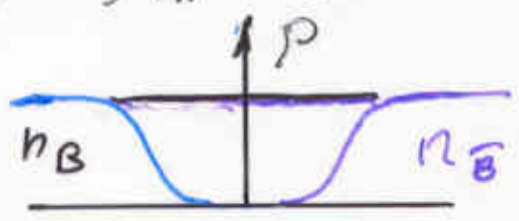
and adiabatic perturbations.

Small (how small?) amount of  $\bar{M}$  nearby is not excluded.

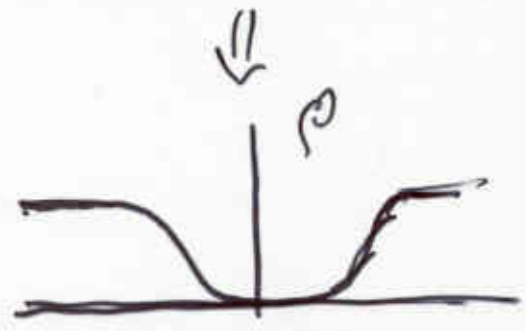


# Isocurvature perturbations:

1)  $\delta\rho_{in} = 0$ , but different chemical content



$$\Delta\rho_{in} = 0$$



$$-\Delta\rho \neq 0$$



$$\Delta T \neq 0$$

↑  
excessive  
pressure and  
M/M separation

Non-relativistic  
matter cools faster

Role of gravity of large scale  
structure in DM-??

# Conclusion.

1. Simplest models of baryogenesis predict no antimatter but they are also not very good for explanation of the observed asymmetry
2. Spontaneous BG and Affleck-Dine one can explain observed asymmetry and naturally predict large fraction of ~~a~~ antimatter.
3. In a version of AD-model with coupling to inflaton baryon rich compact objects may be in galactic halo and even make all DM from BH and anti-BH surrounded correspondingly by  $M$  or  $\bar{M}$ .
4. Heavy  $\bar{A}$  may be even more abundant than  $\overline{He^4}$ .
5. In model 3 detailed predictions are possible (if  $DM = BH + \bar{BH}$ ) but more (a lot) work is necessary.