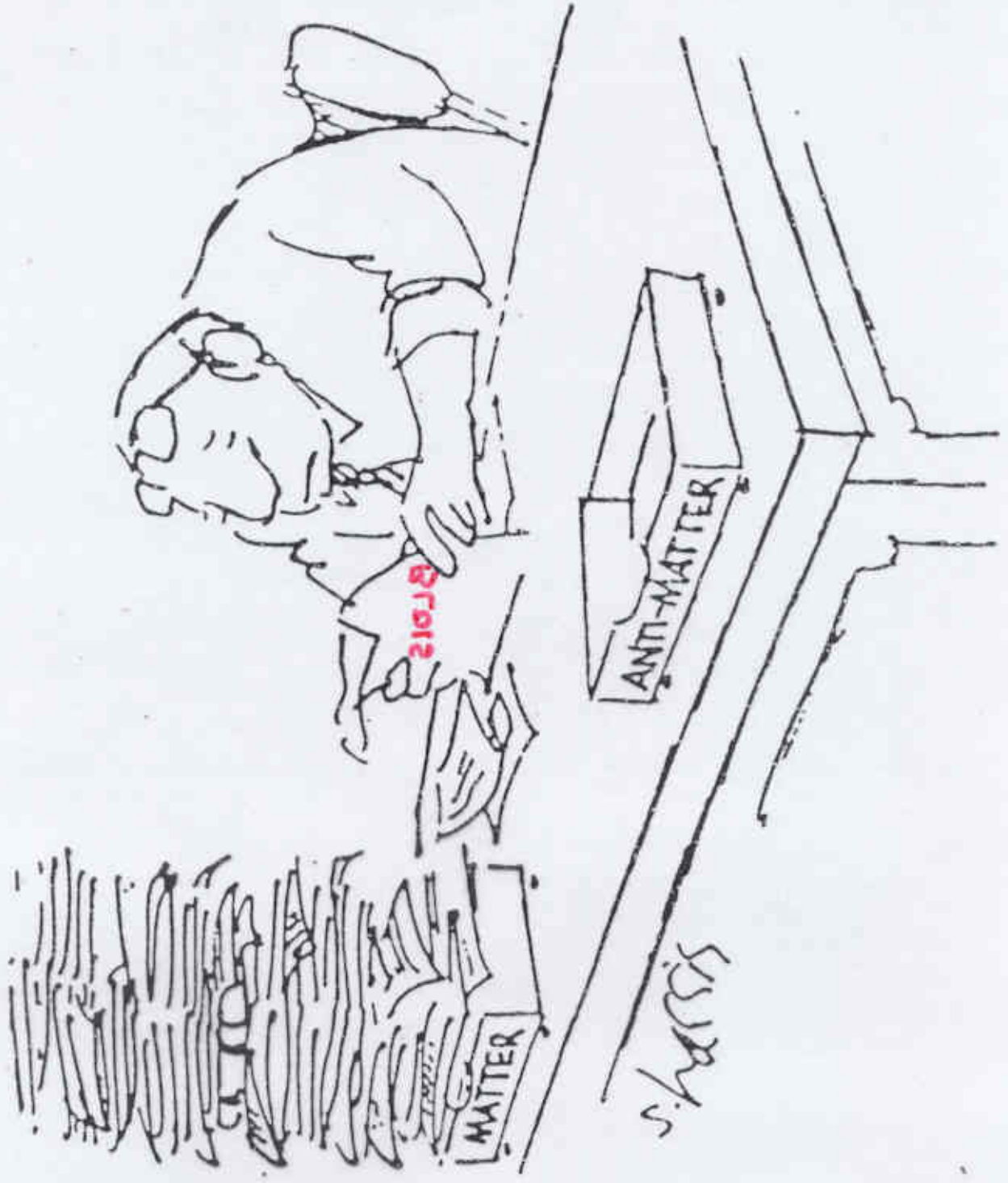


THE BARYON ASYMMETRY OF THE UNIVERSE

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OBSERVATIONS AND THEORETICAL CONSIDERATIONS



THE BARYON ASYMMETRY OF THE UNIVERSE

Big-bang nucleosynthesis calculations compared with cosmic H, He, d and L measurements and cosmic background radiation anisotropy measurements imply:

$$\Omega_B \simeq 0.04,$$

i.e., baryonic matter makes up 4% of the closure density of the universe (*e.g.*, review of Olive 2002; Sievers, *et al.* 2002).

But – if N and \bar{N} were in thermal equilibrium in the early hot big-bang and dropped out of equilibrium owing to the expansion of the universe, *i.e.*, when the interaction rate,

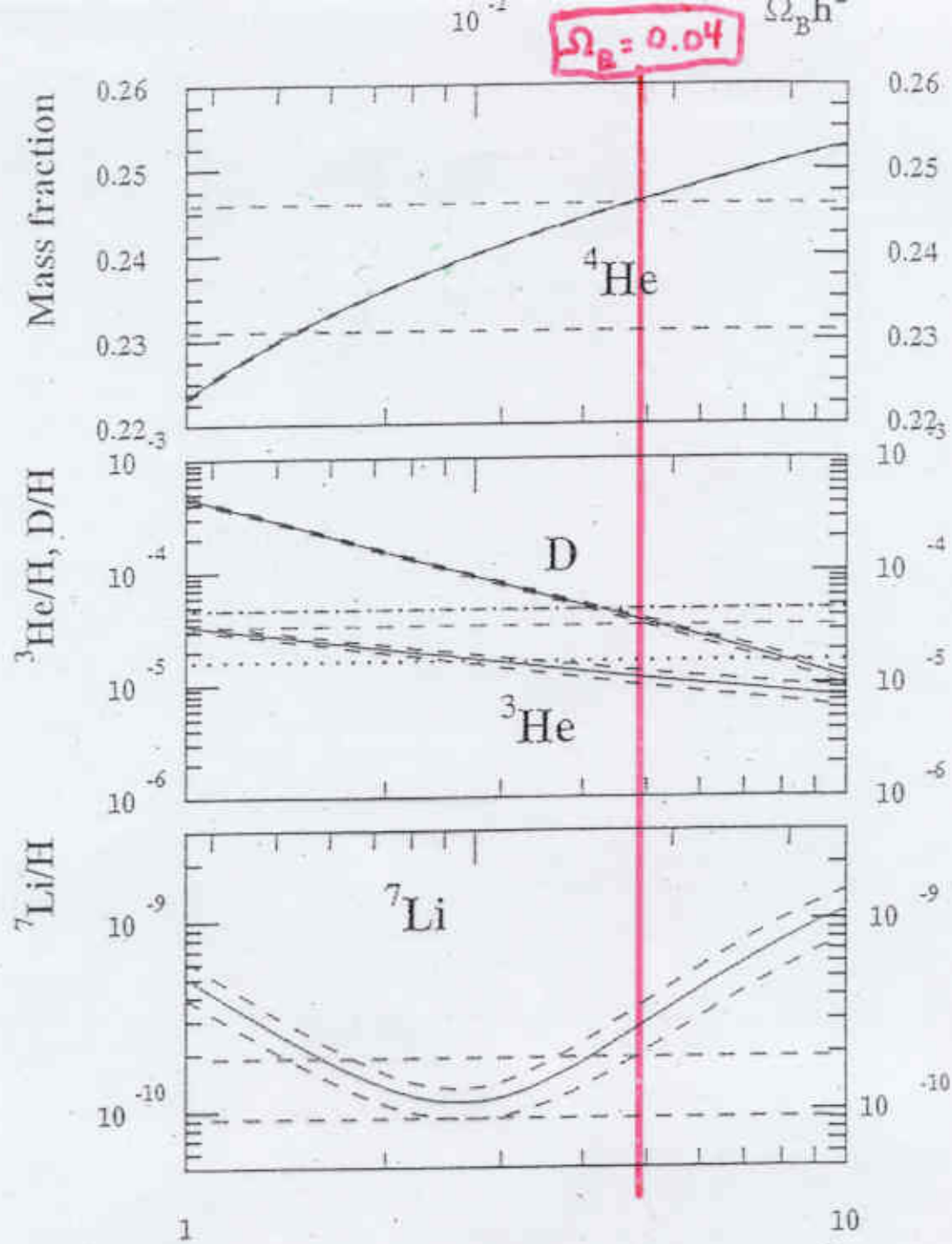
$$\Gamma = n_N v \sigma \sim n_N v m_\pi^{-2}$$

dropped below the expansion rate, H , which would have occurred at $m/T \sim 40$ ($T \simeq 20$ MeV), then we would expect

$$\Omega_B \sim 4 \times 10^{-11},$$

i.e., nine orders of magnitude less than the value implied by cosmological observations (Chiu 1966).

RELATIVE ABUNDANCES FROM
BIG BANG NUCLEOSYNTHESIS



A. Coc et al. (2002) Phys. Rev. D. $\eta \times 10^{10}$

FIG. 4. Abundances of ^4He (mass fraction), D , ^3He and ^7Li (by number relative to H) as a function of the baryon over photon ratio η . Mean values (solid curves) and 2σ limits (dashed curves) are obtained from Monte Carlo calculations. Horizontal lines represent primordial ^4He , D and ^7Li abundances deduced from observations (see text). For D the dotted lines represents the range of observed values (see Fig. 1) while the dashed lines corresponds to the adopted value of Ref. [20].

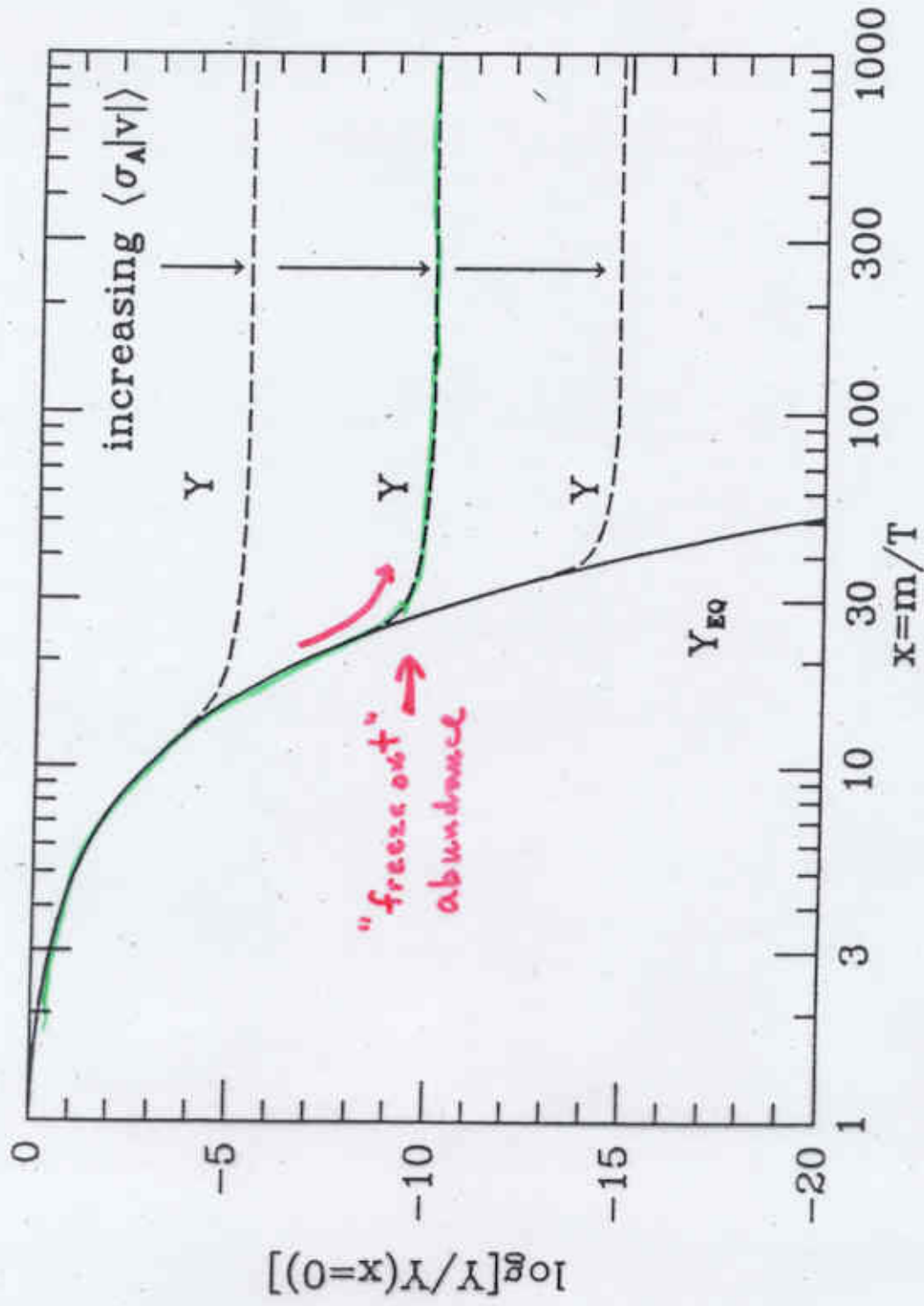


Fig. 5.1: The freeze out of a massive particle species. The dashed line is the actual abundance, and the solid line is the equilibrium abundance. (Kolb & Turner)

You can't make an omelet without breaking
some eggs.

You can't make the Universe without
breaking some symmetries.

Sakharov (1967) Conditions for Matter-
Antimatter Asymmetry

B Violation (Grand Unified Theories)

CP Violation

Non-equilibrium (expansion)

BARYON NONCONSERVATION IN EARLY UNIVERSE

(WEINBERG SCENARIO)

IN BIG-BANG $T = t^{-1/2}$

WHEN T HIGH ENOUGH ($T \gg m_X c^2/k$), LOTS OF X, \bar{X} PAIRS OF GAUGE (OR HIGGS) BOSONS IN THERMAL EQUILIBRIUM

CONSIDER 2 DECAY MODES

PARTICLE	→	MODE	BRANCHING RATIO
X	→	QL	R
X	→	$\bar{Q}\bar{Q}$	$1-R$
\bar{X}	→	$\bar{Q}L$	\bar{R}
\bar{X}	→	QQ	$1-\bar{R}$

IF CP IS VIOLATED $R \neq \bar{R}$

IF X, \bar{X} 'S FALL OUT OF THERMAL EQUILIBRIUM OWING TO EXPANSION OF UNIVERSE AND THEY CAN DECAY FREELY, NET BARYON NO. PRODUCED WILL BE

$$\Delta B = \frac{1}{2} \left[\frac{1}{3}R - \frac{2}{3}(1-R) - \frac{1}{3}\bar{R} + \frac{2}{3}(1-\bar{R}) \right] = \frac{1}{2}(R-\bar{R})$$

IF CP NOT VIOLATED, $R=\bar{R} \rightarrow \Delta B = 0$; ALSO SIGN OF ΔB DETERMINED BY SIGN OF CP VIOLATIONS.

I) EXPANSION RATE OF UNIVERSE $\dot{R}/R \equiv H = T^{-2}$

FOR WEAK AND SUPERWEAK INTERACTIONS WITH FERMI CONST.'S
OF FORM

$$G_X \sim \frac{\alpha_X}{M_X^2}$$

GET COLLISIONAL INTERACTION RATES

$$II) \quad \Gamma_C^* \sim G_X^2 (KT)^6 \sim \frac{\alpha_X^2}{M_X^4} (KT)^6$$

GET DECAY RATES

$$\Gamma_{X \rightarrow 0} \sim \frac{\alpha_X}{M_X^2} M_X^3 N$$

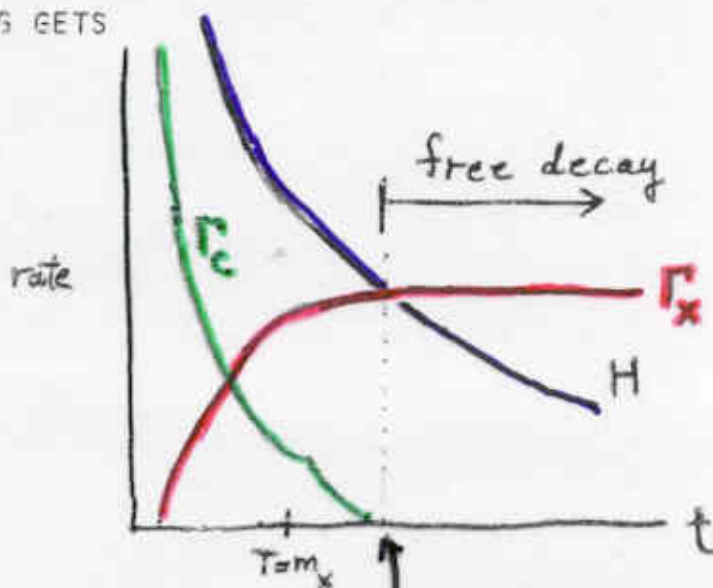
PUTTING IN TIME DILATION

$$III) \quad \Gamma_X \sim \frac{\alpha_X M_X N}{\gamma} \sim \frac{\alpha_X M_X^2 N}{\sqrt{M_X^2 + (KT)^2}}$$

*MORE EXACTLY, TAKING ACCOUNT OF VIRTUAL BOSON PROPAGATION, ETC.,

WEINBERG GETS

$$IV) \quad \Gamma_C \sim \frac{\alpha_X^2 (KT)^6 N}{(M_X^2 + (KT)^2)^2}$$



$T < m_X$ so inverse decay
is blocked by $\exp -m_X/T$

X DECAY SCENARIO:

At $T > M_X$, $n_X \sim n_\gamma$.

If there is a CP asymmetry δ , then

$$n_B \sim \delta n_X \sim \delta n_\gamma$$

But $s \sim g_* n_\gamma$, so that $n_B/n_\gamma \sim \delta/g_*$.

At GUT temperatures $g_* \sim 10^2 - 10^3$.

Therefore the CP asymmetry at the GUT temperature required to explain the present value of n_B/n_γ is

$$\delta \sim 10^{-8}$$

$$\sim 5 \times 10^{-10}$$

At low temperatures, the CP asymmetry observed for the decay modes

$$K_L^0 \rightarrow \pi^\mp e^\pm \nu(\bar{\nu})$$

is $\sim 3 \times 10^{-3}$.

Cosmic Antimatter

To be or not to be; that is the
question!

Baryon Symmetric Domain Cosmology (BSDC)

- R.W. Brown & F.W. Stecker, Phys. Rev. Letters 43, 315 (1979)
- K. Sato, Phys. Lett. 99B, 66 (1981)
- F.W. Stecker, Nucl. Phys. B252, 25 (1985)

4th Condition needed for Local,
Instead of Global, Asymmetry:

- Spontaneous CP breaking at GUT level followed by moderate inflation.

SPONTANEOUS SYMMETRY BREAKING



HIDDEN SYMMETRY



DOMAINS

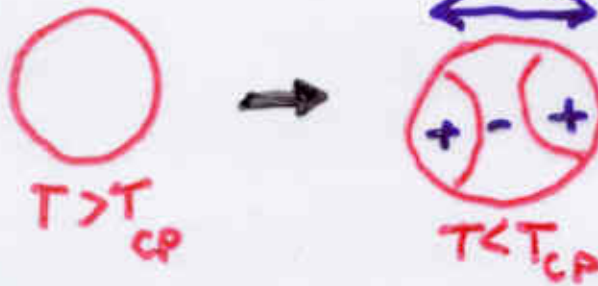
cf. FERROMAGNETISM

NO
MAGNETIZATION

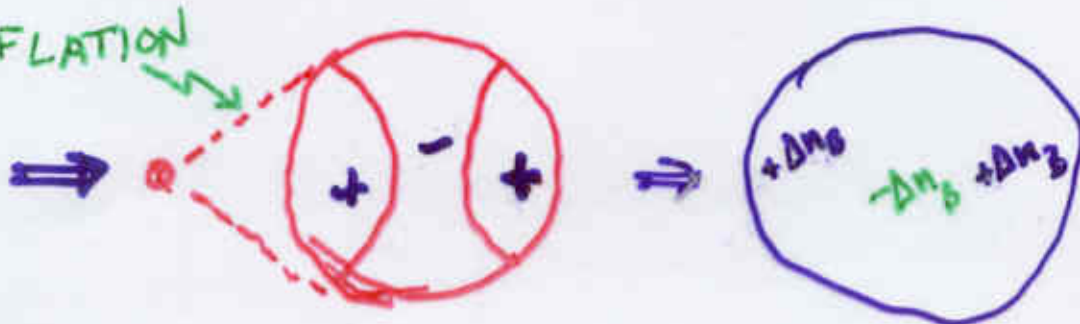


CP CASE

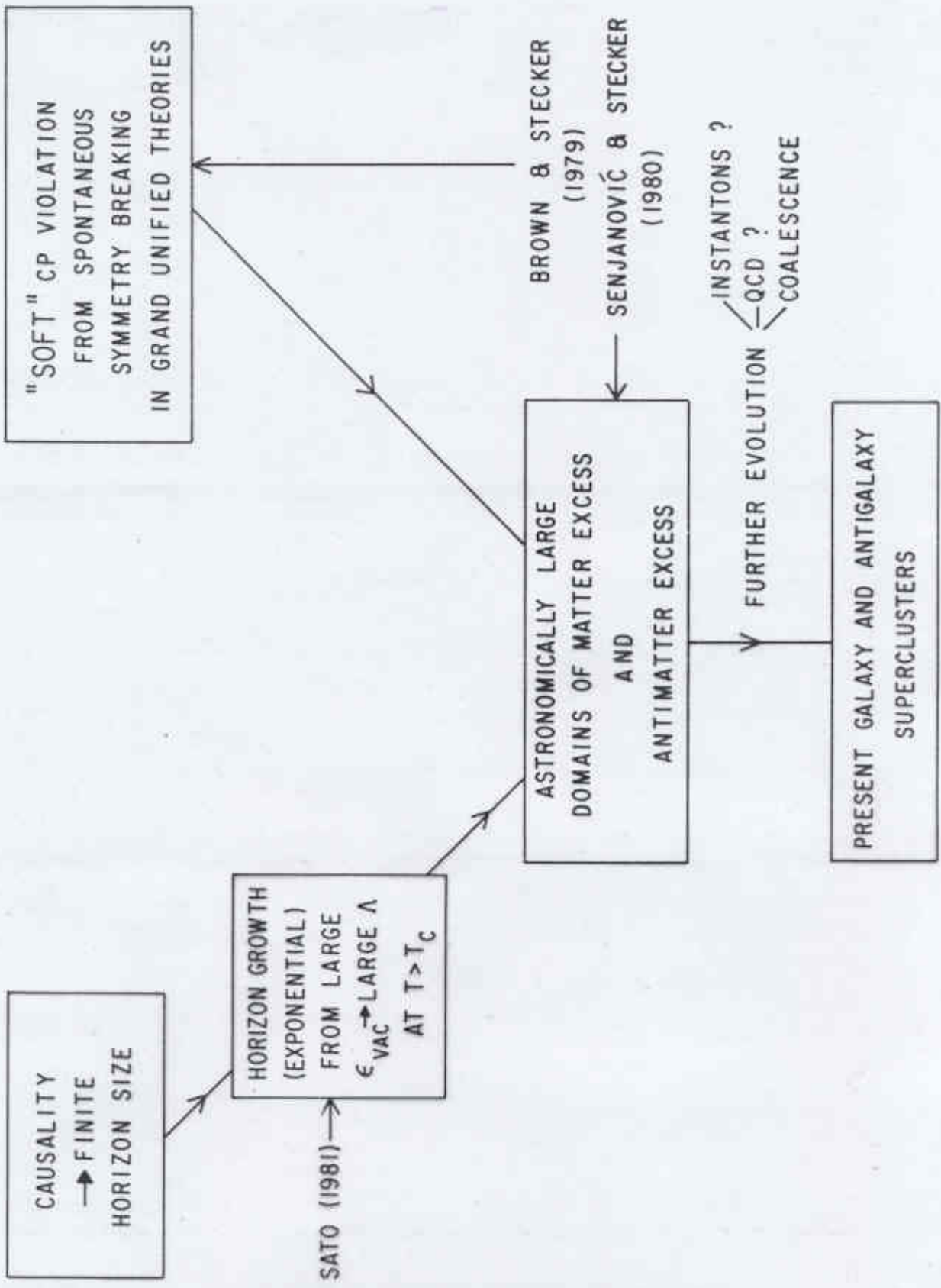
CP SYMMETRIC



INFLATION



SIMPLEST BARYON SYMMETRIC BIG-BANG SCENARIO



Technical Problems with BSDC

- Fine tuning of inflation
- Domain Walls

POSSIBLE SOLUTION TO THE CP DOMAIN WALL PROBLEM

(Mohanty & Stecker, Phys. Lett. 143B (1984) 351)

1. Assume that there is a heavy fermion sector, ψ_i
2. In analogy with QCD, form $\langle \bar{\psi}_i \psi_i \rangle$ condensates in a supercooled, strongly interacting phase. Get terms like

$$\text{Tr} \Phi^3 m(T) e^{3i\beta} + h.c. \propto \cos(3\theta + \beta)$$

where β is a nontrivial phase which breaks the CP degeneracy and drives the universe to a "true vacuum" after the baryon (and antibaryon) asymmetries are created.



...and each domain grows into
either a matter or an antimatter
supercluster...

That's ridiculous!!!

...and each domain grows into
either a matter or an antimatter
supercluster...

That's ridiculous!!

CP VIOLATION

We finally understand weak interactions



J. Fabergé, CERN Courier, 6, No. 10, 193 (October 1966). (Courtesy of Madame Fabergé.)

PRODUCTION OF BARYON AND ANTI-BARYON EXCESSES

NECESSARY CONDITIONS:

- I. NON-EQUILIBRIUM (SUPPLIED BY EXPANSION OF THE UNIVERSE)
- II. PROCESSES WHICH DO NOT CONSERVE BARYON NUMBER (GUTS)
- III. BREAKING OF CP SYMMETRY

CRUCIAL QUESTION: HOW IS CP SYMMETRY BROKEN?

POSSIBILITIES:

- A) CP BROKEN AT ALL TIMES AND TEMPERATURES BY EXPLICITLY PUTTING "HARD" CP VIOLATION INTO THE LAGRANGIAN OF THEORY IN THE FORM OF COMPLEX YUKAWA COUPLINGS F

$$L_Y = F \psi \bar{\psi} \psi$$

OR COMPLEX HIGGS SELF-COUPLINGS λ

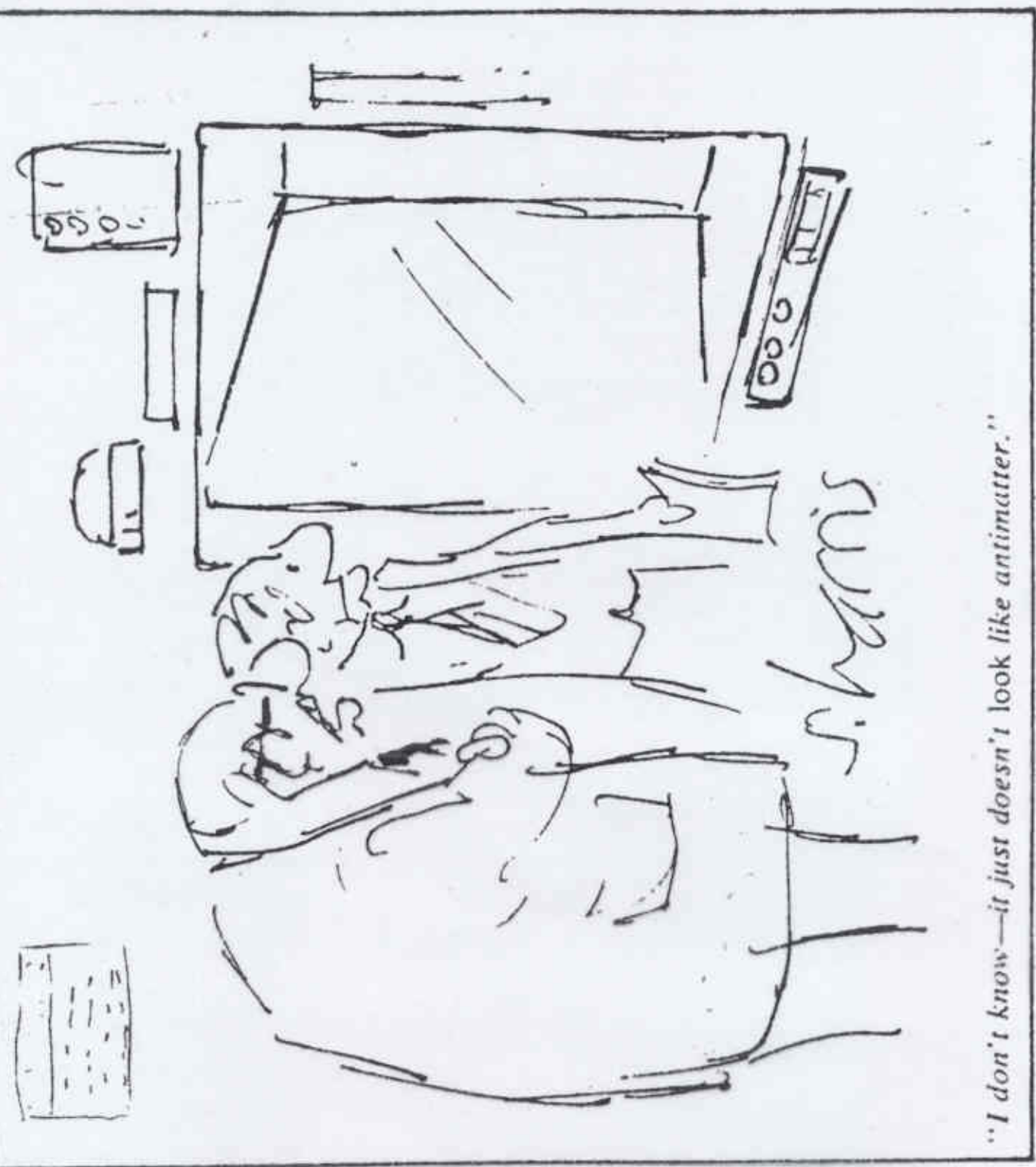
$$L_H = \lambda \phi^4$$

B) CP VIOLATION "SOFT" :

1. START WITH REAL COUPLINGS IN THE LAGRANGIAN SO THAT IT IS CP - INVARIANT
2. BREAK CP SPONTANEOUSLY THROUGH COMPLEX VEV'S IN THE HIGGS FIELDS

ADVANTAGES:

1. CAN USE PERTURBATION THEORY TO CALCULATE CP VIOLATION
2. SOLVES THE "STRONG CP PROBLEM" IN QCD
3. IS HARD TO AVOID IN ANY CASE IN A "REALISTIC" GUT



"I don't know—it just doesn't look like *animatter*."

COSMIC RAY $\bar{\text{He}}$ LIMITS

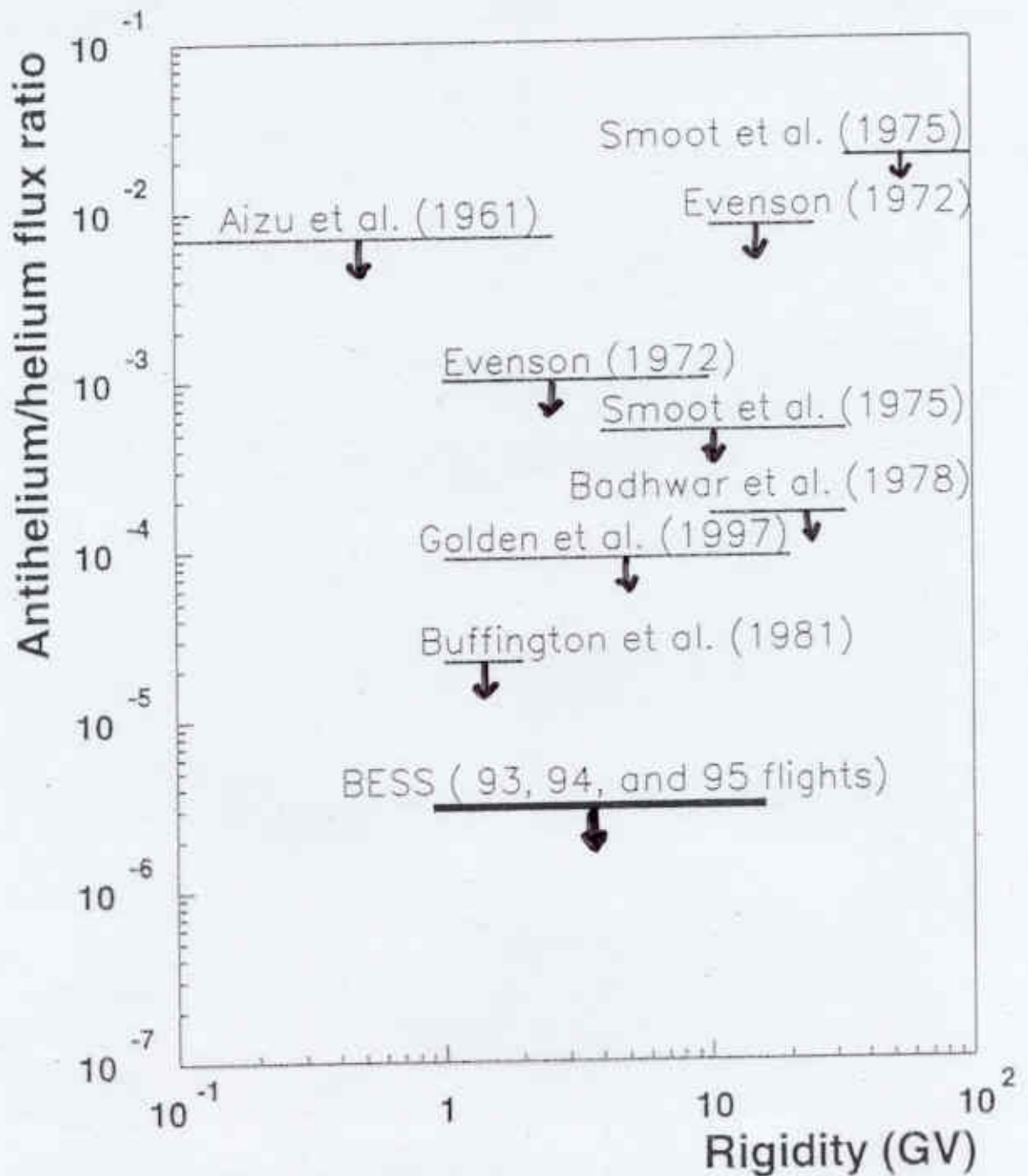


FIG. 5. The resultant upper limit of $\bar{\text{He}}/\text{He}$ flux ratio of the present paper together with previous limits.

Saeki, et al.
astro-ph/9710228

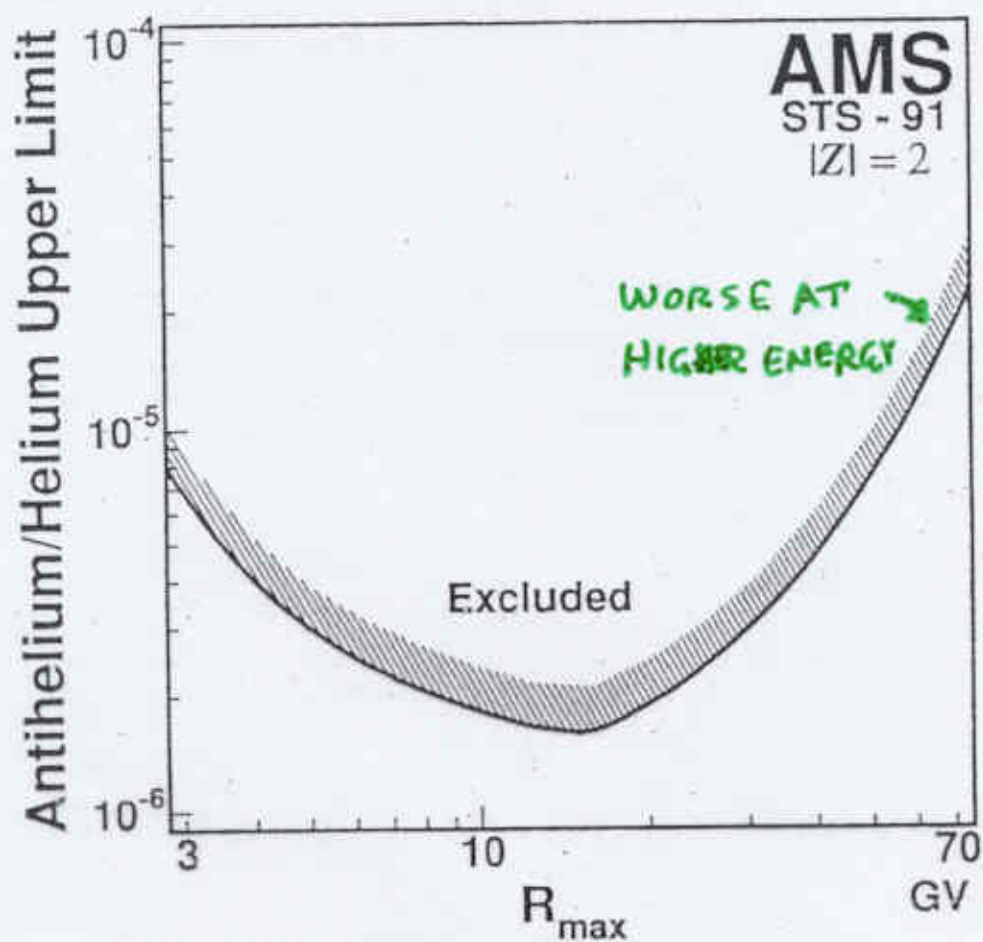


Fig. 9. Upper limits on the relative flux of antihelium to helium, at the 95% confidence level, as a function of the rigidity interval $R = 1.6$ GV to R_{\max} . These results are independent of the incident antihelium spectra.

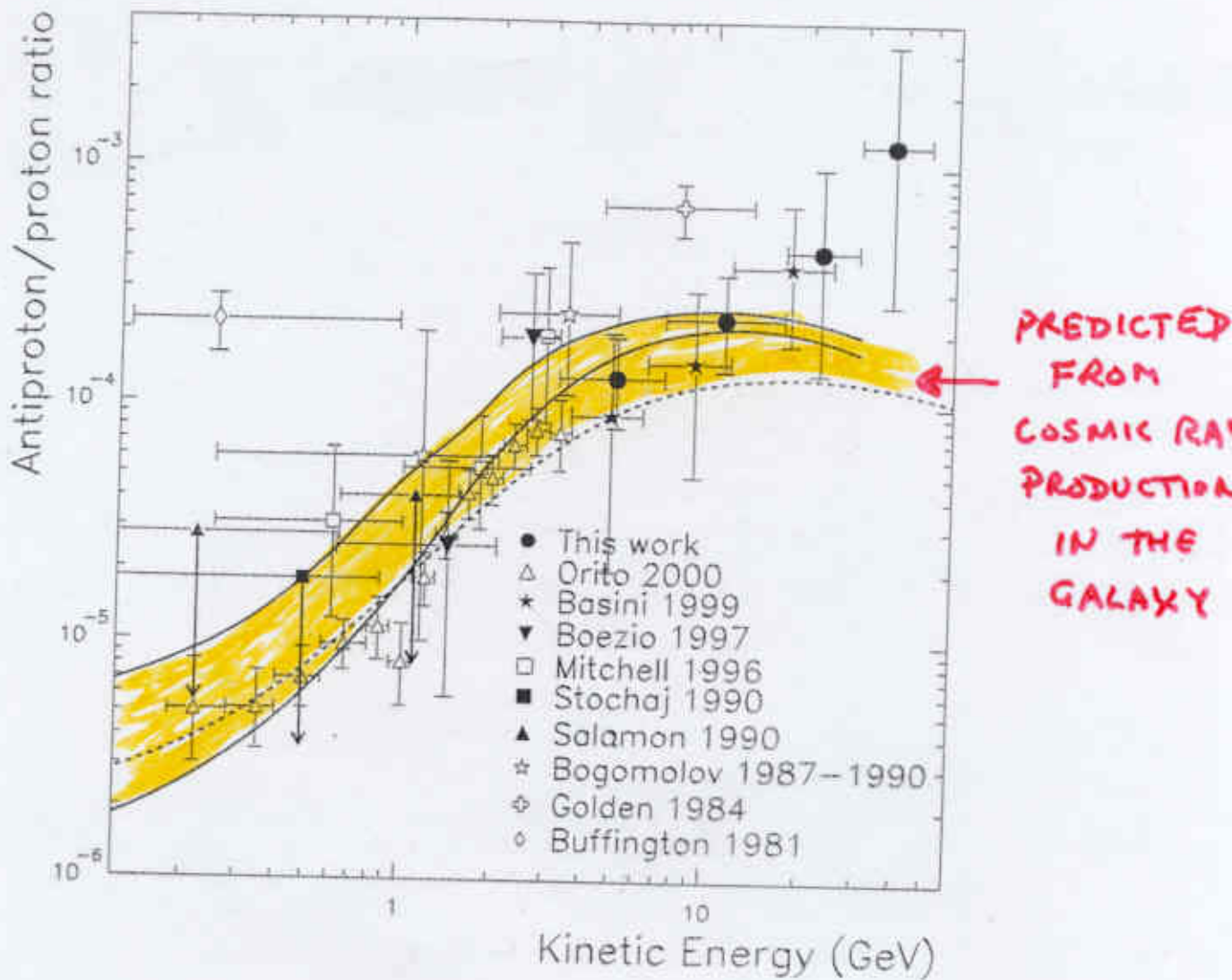
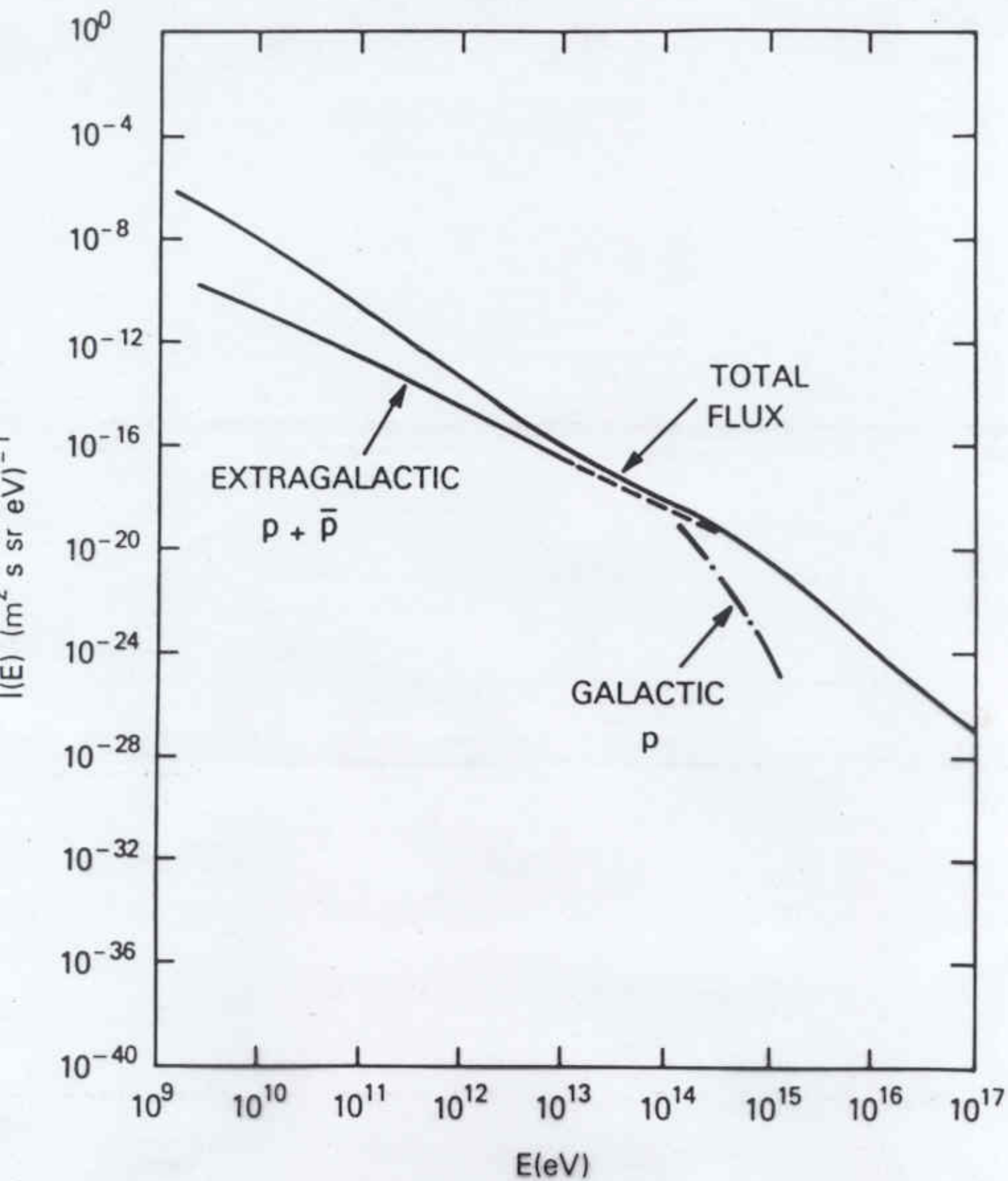


Fig. 10.— The \bar{p}/p ratio at the top of the atmosphere obtained in this work compared with previous measurements (Buffington, Schindler, & Pennypacker 1981; Golden et al. 1984; Bogomolov et al. 1987, 1990; Salamon et al. 1990; Stochaj et al. 1990; Mitchell et al. 1996; Boezio et al. 1997; Basini et al. 1999; Orito et al. 2000). The lines are the calculations of interstellar antiprotons assuming a pure secondary production during the propagation of cosmic rays in the Galaxy by Simon, Molnar, & Roesler (1998) (solid lines, upper and lower limits of the calculation) and by L. Bergström and P. Ullio 1999 (private communication) (dashed line).

Boezio et al.
 astro-ph/0103513



Constraints on Large-Scale Baryon- Antibaryon Domains

CBR Distortion

Gamma-Ray Background Radiation

Cosmic Ray Antihelium*

*If it can reach us.

And today on our stage we are reuniting, after billions of years of separation, Matt & Anti-Matt, twins who were seperated at birth in the Big Bang.

THE
MUMASON
SHOW

© 1995 GABE MARTIN



Sub-Atomic Talk-Show Disasters

Distortion of CBR by Annihilation Energy Input in High Energy Particles

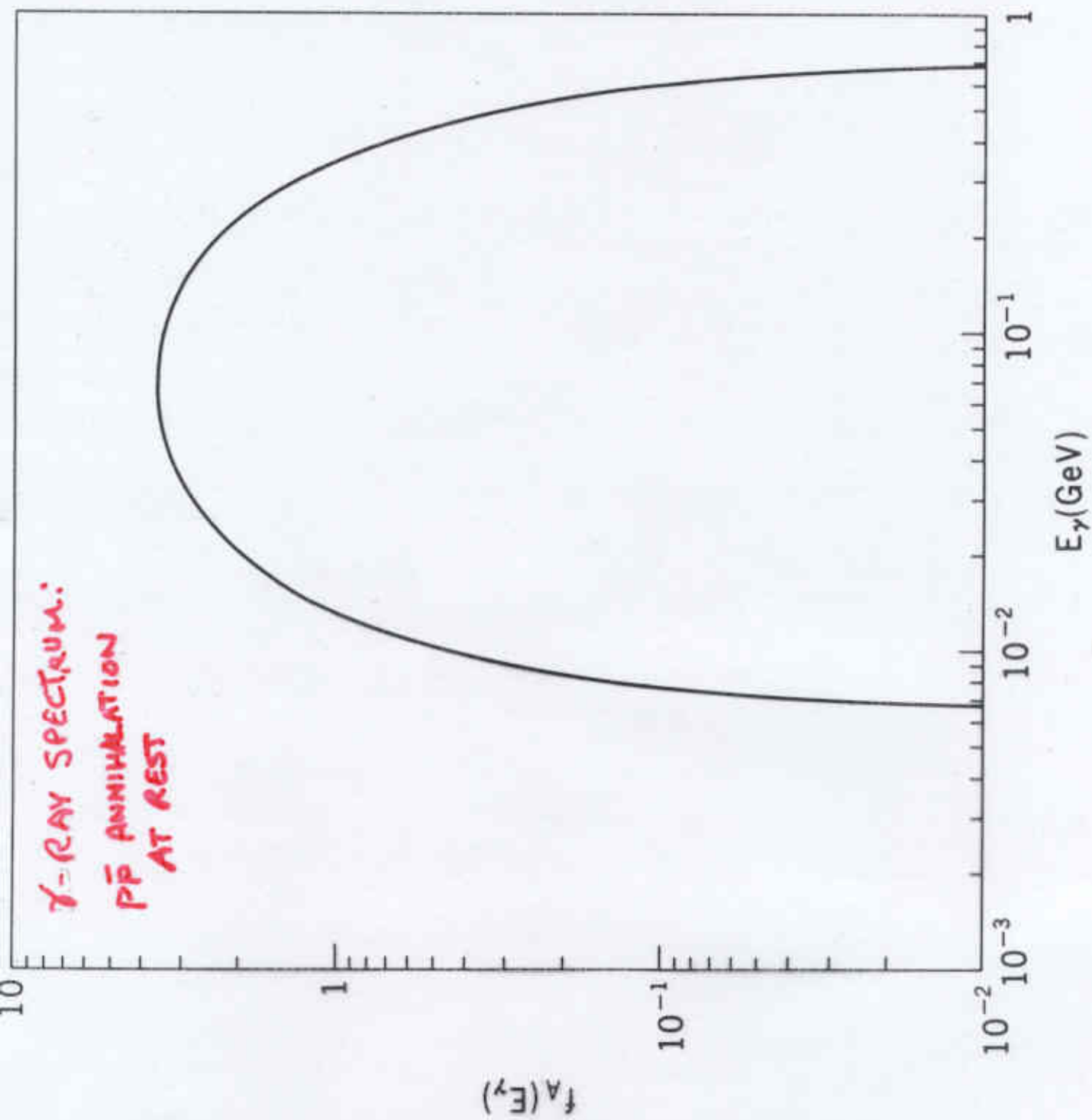
Parameter $y = (\text{fractional energy input})/4$:

From COBE measurements, y is less than
 2.5×10^{-5} (Wright et al. 1994)

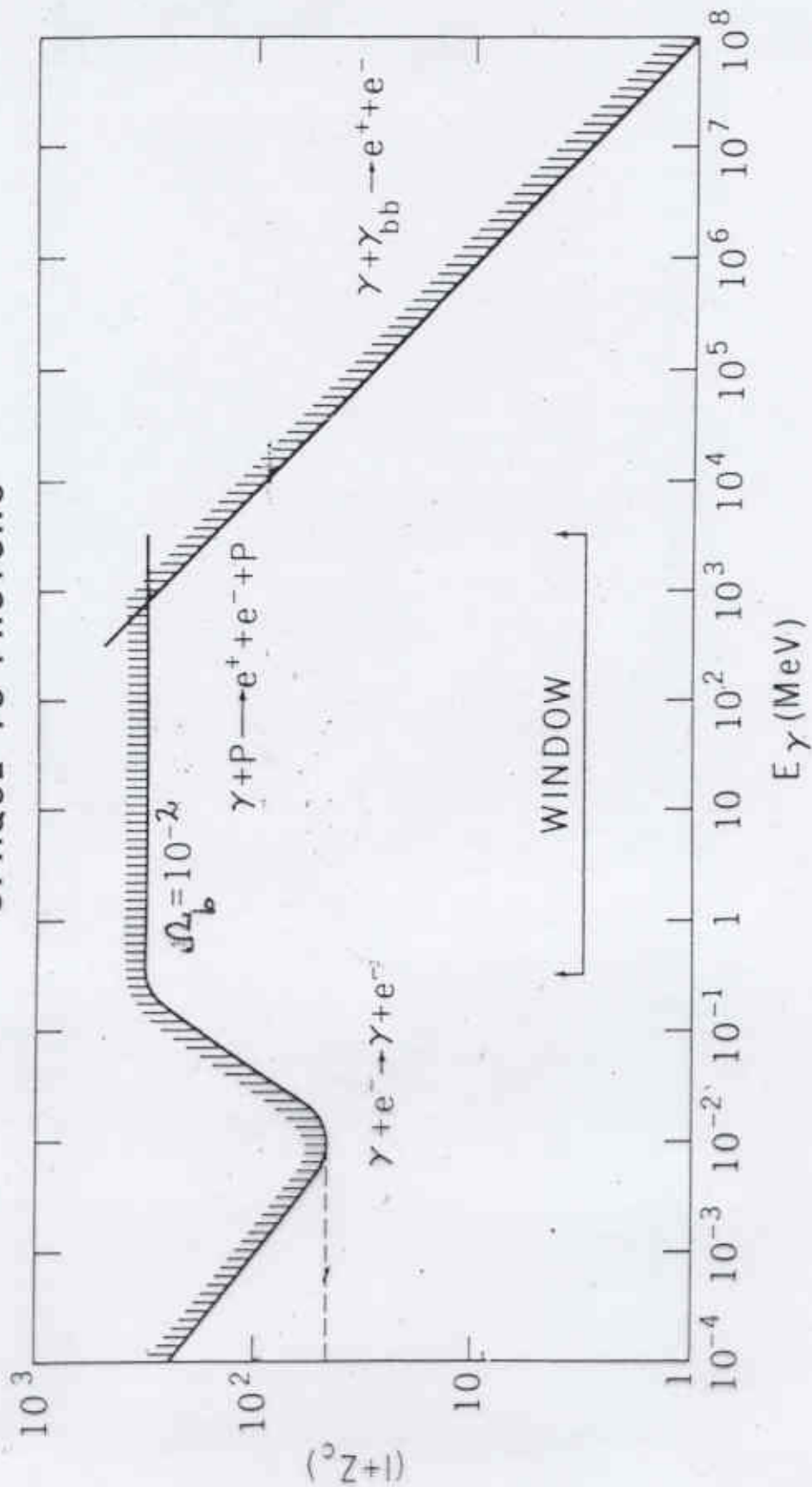
Predicted y less than 9×10^{-9} (Cohen et al. 1998)

No Conflict!

**γ -RAY SPECTRUM:
P \bar{P} ANNIHILATION
AT REST**



REDSHIFT AT WHICH THE UNIVERSE BECOMES OPAQUE TO PHOTONS



This integro-differential equation takes account of γ -ray production, absorption, scattering, and redshifting and is of the form

$$\frac{\partial \mathcal{I}}{\partial t} + \frac{\partial}{\partial E} [-EH(z)\mathcal{I}] = \mathcal{Q}(E, z) - \kappa_{AB}(E, z) \mathcal{I} + \int_E^\infty \epsilon(E') \kappa_{sc}(E, z) \mathcal{I}(E'; E') dE' \quad (1)$$

where

$$\mathcal{I}(E, z) \equiv (1+z)^{-3} I(E, z)$$

$$\mathcal{Q}(E, z) \equiv (1+z)^{-3} Q(E, z)$$

and

$$\frac{\partial \mathcal{I}}{\partial t} = - (1+z) H(z) \frac{\partial \mathcal{I}}{\partial z},$$

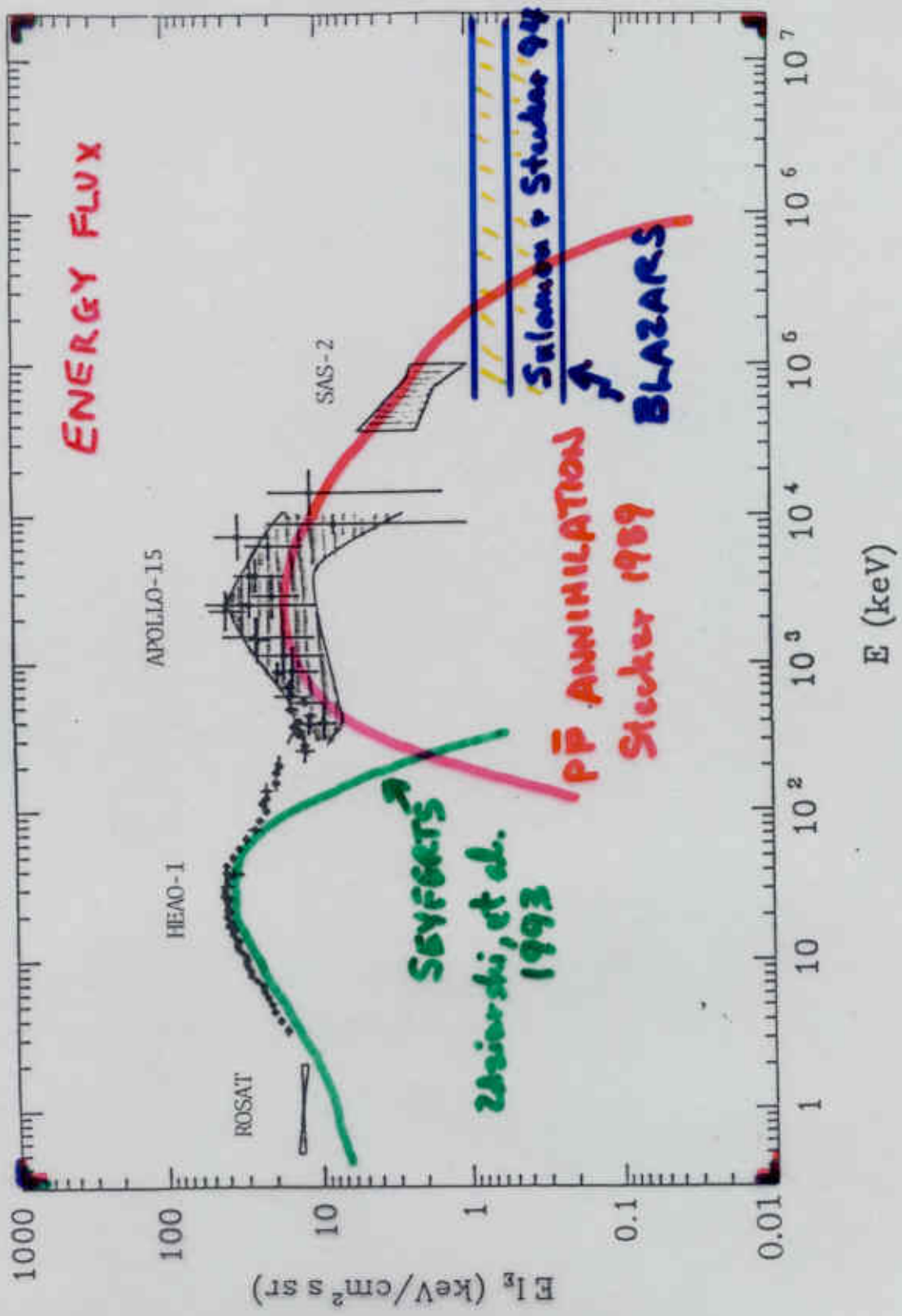
$$H(z) = H_0 (1+z)^{-1/2}$$

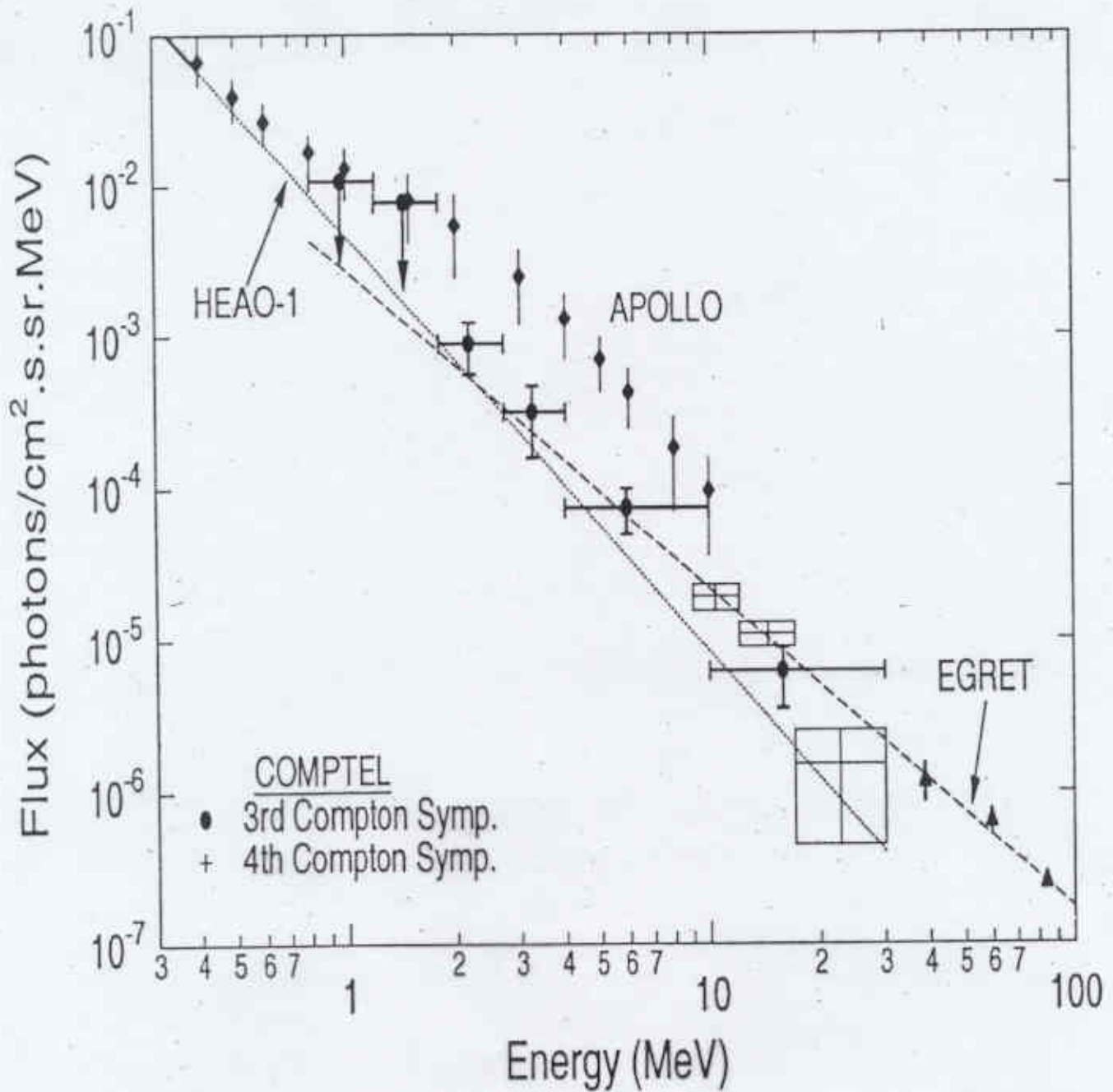
Cosmological Photon Transport (CPT) Equation

- Stecker, Morgan & Brodeur, *Phys. Rev Letters* 27, 1469 (1971)

The second term in eq. (1) expresses energy loss from the redshift effect. The third term is the γ -ray source term from $p\bar{p}$ annihilation primarily into π^0 's. The absorption term is from pion production and Compton interactions with electrons at high z and the scattering integral puts back Compton scattered γ -rays at lower energies $E < E'$.

EXTRAGALACTIC X- AND γ -RAY
BACKGROUND





PHOTON FLUX *

* LARGE INTRINSIC BACKGROUND SUBTRACTIONS!

Minimum Domain Size

- From Gamma-Ray Background* Constraints:

A) 10 Mpc (Stecker & Puget 1972 – based on data of Trombka et al. (1972))

B) 1 Gpc (Cohen et al. 1998 – based on data of Kappadath et al. 1995)

*all data in the critical MeV range are questionable – need dedicated experiment!

Summary of Astrophysical Information:

Completely mixed antimatter ruled out by thermodynamics (Chiu 1966).

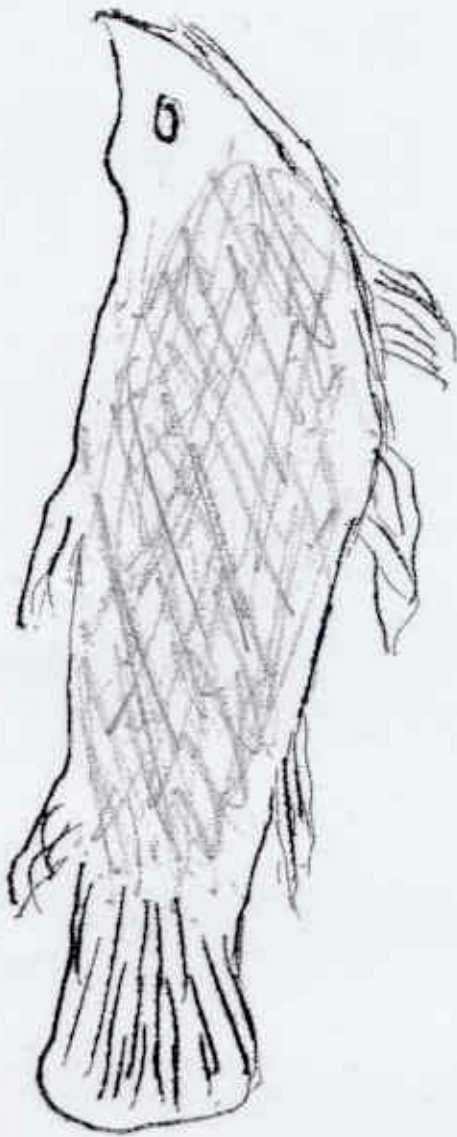
Sub-galaxy size antimatter domains ruled out by big-bang nucleosynthesis (Kurki-Suonio & Sihvola 2000; Rehm & Jedamzik 2002).

Large-scale antimatter in our galaxy ruled out by low-energy cosmic ray measurements.

Gamma-ray background data indicate that matter-antimatter domains must be at least galaxy supercluster size.

"HELICANTH"

12



Conclusions:

There is no evidence for cosmologically significant amounts of antimatter in the universe.

With inflation and spontaneous CP breaking one can get a either a locally or globally baryon asymmetric universe.

Astrophysical data are not conclusive. It is still possible that there are domains of antimatter galaxies in the universe; absence of evidence is not evidence of absence. We need reliable gamma-ray background measurements as a conclusive test.