Investigating Lorentz and CPT symmetry with antihydrogen

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Outline

Lorentz & CPT violation

 the standard-model extension
 signal properties

 Hydrogen and Antihydrogen

 1S to 2S
 Hyperfine Zeeman
 Results in other systems
 Neutral mesons, Atomic systems, Electrodynamics



Approaches to testing CPT

	Phenomenological	Theoretical Framework
PROS	 Practical approach 	 Consistency with standard mode Experimental breadth Predicts signal types
CONS	 Unclear theoretical basis Experiment dependent Limited predictive ability 	•Challenging to find

Standard-model extension: a consistent, microscopic, general theoretical framework allowing Lorentz and CPT violation

Kostelecký, Potting, PRD **51** 3923 (1995) Colladay, Kostelecký, PRD **55** 6760 (1997); PRD **58** 116002 (1998)



Kostelecký, Potting, PRD **51** 3923 (1995) Colladay, Kostelecký, PRD **55** 6760 (1997); PRD **58** 116002 (1998)



Lorentz Transformations

Standard-model extension properties

Conventional	Unconventional
Gauge structure Power-counting renormalizabilit Energy and momentum conserva Quantization Microcausality Spin-statistics Observer Lorentz covariance	y tion Particle Lorentz non-covariance CPT violation

Kostelecký, Potting, PRD 51 3923 (1995) Colladay, Kostelecký, PRD 55 6760 (1997); PRD 58 116002 (1998)

Special cases of standard-model extension include:

Realistic noncommutative field theories: $[x^{\mu}, x^{\nu}] = i\theta^{\mu\nu}$

-- q plays role of prescribed background tensor in QFT

Carrol et al, PRL 87 141601 (2001)

I sotropic high-energy limit:

-- preferred frame,

-- one parameter for each species: Coleman, Glashow PRD 59 116008 (1999) Others...

Sidereal variations

Kostelecký, PRL 80 1818 (1998)

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cold antihydrogen:

- -- theoretical analysis possible for simple system
- -- clean bounds should be possible
- -- atomic clocks: no antiatoms, analysis complex, need models

Antiprotonic atoms

ASACUSA - third talk this session, Widmann

Antiprotonic He, for example:

-- potential for tests involving antiproton/proton comparison

1S-2S transition in trapped antihydrogen and hydrogen

Allowed 1S-2S transitions

Only the c and d states are trapped

$$\begin{aligned} |d\rangle_n &= \left|\frac{1}{2}, \frac{1}{2}\right\rangle \quad, \\ |c\rangle_n &= \sin\theta_n \left|-\frac{1}{2}, \frac{1}{2}\right\rangle + \cos\theta_n \left|\frac{1}{2}, -\frac{1}{2}\right\rangle \\ |b\rangle_n &= \left|-\frac{1}{2}, -\frac{1}{2}\right\rangle \quad, \\ |a\rangle_n &= \cos\theta_n \left|-\frac{1}{2}, \frac{1}{2}\right\rangle - \sin\theta_n \left|\frac{1}{2}, -\frac{1}{2}\right\rangle \end{aligned}$$

$$\tan 2\theta_n \approx \frac{(51 \text{ mT})}{n^3 B}$$

What are the effects on the c \rightarrow c and d \rightarrow d transitions?

Shifts in energy levels

Hydrogen with electron and proton angular momenta J and I: $\Delta E^{H}(m_{J}, m_{I}) \approx (a_{0}^{e} + a_{0}^{p} - c_{00}^{e}m_{e} - c_{00}^{p}m_{p}) + (-b_{3}^{e} + d_{30}^{e}m_{e} + H_{12}^{e})m_{J}/|m_{J}| + (-b_{3}^{p} + d_{30}^{p}m_{p} + H_{12}^{p})m_{I}/|m_{I}|$

Antihydrogen: reverse signs of a, d and H parameters

No zero-order effect in $d \rightarrow d$ transition, since n=1 and n=2 states have identical spin:

$$|m_{J} = +1/2, m_{I} = +1/2\rangle$$

Bluhm, Kostelecký, Russell, PRL **82** 2254 (1999)

Signal in $|c > \rightarrow |c >$ transition

Spin mixing is different in 1S and 2S \rightarrow unsuppressed signal

For $c \rightarrow c$ in hydrogen:

Hyperfine Zeeman transitions

Shifts are in opposite directions

$$\delta\nu_{c\to d}^{H} \approx (-b_{3}^{p} + d_{30}^{p}m_{p} + H_{12}^{p})/\pi$$
$$\delta\nu_{c\to d}^{\overline{H}} \approx (b_{3}^{p} + d_{30}^{p}m_{p} + H_{12}^{p})/\pi$$

Instantaneous comparison:

$$\nu^{H}_{c \to d} - \nu^{\overline{H}}_{c \to d} \approx -2b_{3}^{p}/\pi$$

gen With 1mHz resolution,

$$\left|b_3^p\right| \le 10^{-27} \,\mathrm{GeV}$$

Bluhm, Kostelecký, Russell, PRL **82** 2254 (1999) Cleaner than any competing experiment

Hydrogen maser experiment

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Neutral Mesons

Relevant mesons: K, D, $B_{d'}$, B_s and antiparticles

Time evolution via 2×2 effective hamiltioni an Λ : $i\partial_t \Psi = \Lambda \Psi$

CPT symmetry $\Leftrightarrow \mathbf{x} = 0$ where $\mathbf{x} \approx \Lambda_{11} - \Lambda_{22}$

Standard - model extension:

$$\boldsymbol{x} \approx \boldsymbol{b}^{m} \Delta a_{m}$$

Decay probabilities have

velocity â dependence

sidereal dependence

Assumption of constant \hat{i} is incompatible with QFT

Kostelecký, PRL **80** 1818 (1998); PRD **61** 016002 (1999); PRD **64** 076001 (2001)

Neutral Mesons: status of bounds $on u_{\mu}$

		tilde coefficient	Prestage <i>et al</i> .	Lamoreaux <i>et al.</i>	Chupp et al.	Berglund <i>et al.</i>	Bear <i>et al.</i>
Clock-comparison	р	b	*	*		-27	*
	p	d	*	*	-	-25	*
bounds	p	9D.J	*	*	2	-25	*
	р	CQ.J	*	_	200	-	-
	р	9Q,J	*	_		-	<u> </u>
	р	C_	*	*	*	-	-
⁹ Be ⁺ and H maser	р	g_	*	*	*	-	-
	р	CXY	*	*	*	_	-
²⁰¹ Hg and ¹⁹⁹ Hg	р	9xy	*	*	*	-	-
	n	bı	-27	-29		-30	-31
² Ne and ³ He	n	d	-25	-26	-	-28	-29
1994 and 1330 c	n	gn.i	-25	-27	(<u></u>	-28	-29
	n	CO.J	-25	—	-		_
$3 \sqcup_{0}$ and $129 \lor_{0}$	n	90.1	*	3 <u></u> 3	X 		-
	n	C	-25	-27	-27		-
	n	g	*	*	*	_	-
	n	CXY	-25	-27	-27	100	
Sidoroal CoV	n	g xy	*	*	*	-	-
Sidereal, Gev	е	b,	-	-	-	-27	-
	е	dJ	-	<u> </u>	_	-22	_
	е	9D.J	-	-	_	-22	
	е	CQJ	_			_	_
	е	9Q,J	-		-	-	-
	е	C_	-	_	-	-	-
Kostelecký, Lane	е	g_	-	-	_	-	-
	е	CXY	-	· _ · ·	—	_	_
PRD 60 116010 (1999)	е	gхү		8 — 3	-	1000	-

¹²⁹Xe/³He Maser

Best test for neutron parameters in standard-model extension

$$\tilde{b}_{\perp}^{n} \equiv \sqrt{(\tilde{b}_{X}^{n})^{2} + (\tilde{b}_{Y}^{n})^{2}} = (4.0 \pm 3.3) \times 10^{-31} \text{ GeV}$$

Bear et al, PRL 85 5038 (2000)

Space-based clock-comparison tests

ACES, PARCS, RACE, SUMO:

Cs, Rb, and H masers planned for International Space Station; Also Superconducting microwave oscillator experiment

Advantages of satellite platform

Orbit inclination and precession gives access to many more coefficient About 50 to 60 coefficients for each experiment Greater speeds and rotation rates accessible in space Boost factor $v = 10^{-4}$ c Earth period 24 hours; Space station 92 minutes Possible technological advantages

Bluhm, Kostelecký, Lane, Russell, PRL 88 090801 (2002)

Penning-trap bounds

Results in context of standard-model extension

Туре	Particles	Result	Reference
lnstantaneou anomaly frequency	^S Electron, positron	10-21	Dehmelt et al, PRL 83 4694 (1999)
Instantaneou cyclotron frequency	^S Hydrogen ion, antiproton	10-26	Gabrielse et al, PRL 82 3198 (1999)
Sidereal anomaly frequency	Electron	10-21	Mittleman et al, PRL 83 2116 (1999)

Bluhm, Kostelecký, Russell, PRL **79** 1432 (1997); PRD **57** 3932 (1998)

Muons

Muonium spectroscopy of hyperfine Zeeman levels

$$\sqrt{(\tilde{b}_X^{\,\mu})^2 + (\tilde{b}_Y^{\,\mu})^2} \le 2 \times 10^{-23} \text{ GeV}$$

CERN and BNL g-2 Experiments

Anomaly-frequency comparisons b_z^m< 10⁻²³ GeV

Sidereal anomaly-frequency variations (estimated) b_m<10⁻²⁴ GeV

> Bluhm, Kostelecký, Lane, PRL **84** 1098 (2000); Hughes et al, PRL **87** 111804 (2001); Deile et al, CPT '01 Proceedings

m-

Electrodvnamics

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}(k_F)_{\kappa\lambda\mu\nu}F^{\kappa\lambda}F^{\mu\nu}$$

Standard-model extension predicts birefringence
 Sensitive tests using astrophysical sources

Bounds from pulse spreading: $|k^a| < 3x10^{-16}$

Bounds from differential polarization ro $|k^a| < 2x10^{-32}$

Kostelecký, Mewes, PRL **87** 251304 (2001); hep-ph/0205211

"Best test of Special Relativity"

Summary

Standard-model extension: a viable theoretical framework that allows Lorentz and CPT violation

Trapped Antihydrogen

Unsuppressed signal occurs in the 1S-2S transition

Another occurs in the hyperfine Zeeman transition

Comparisons of hydrogen and antihydrogen would produce clean bound on an untested parameter combination

Bounds exist in other systems neutral mesons, atomic systems, electrodynamics