

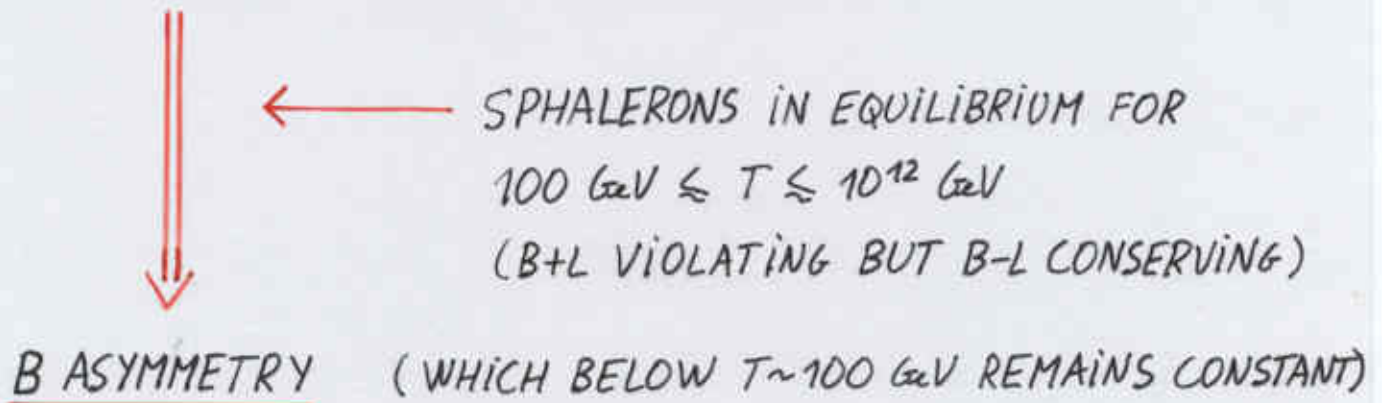
TESTS OF LEPTOGENESIS AT LOW ENERGY

T. HAMBYE
CPT-MARSEILLE

1 | LEPTOGENESIS WITH HEAVY RIGHT-HANDED NEUTRINOS

FUKUGITA-YANAGIDA(86)

- L ASYMMETRY PRODUCED AT $T \gg 100 \text{ GeV}$



- TO GENERATE A L ASYMMETRY AT $T \gg 100 \text{ GeV}$:

HEAVY RIGHT-HANDED NEUTRINOS : $N_{R_i}, i=1,2,3.$

$$\mathcal{L} = \mathcal{L}_{SM} + h_{ij} \bar{N}_{R_i} L_j \Phi^* + \frac{1}{2} \bar{N}_{R_i}^c M_{N_{ij}} N_{R_j} + \text{h.c.}$$

↑
YUKAWA
($\Delta L = 0$)

↑
MAJORANA MASS TERM
($\Delta L = 2$)

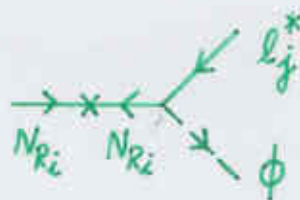
\Rightarrow L VIOLATION IN THE N_R DECAYS

N_R DECAYS:

TREE LEVEL

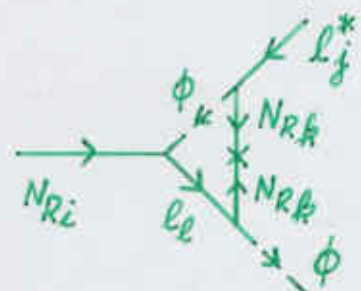
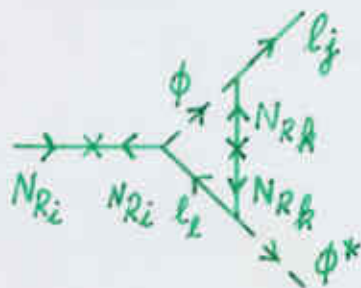


$$L_{FIN} = +1$$



$$L_{FIN} = -1$$

ONE LOOP



↳ DON'T CREATE THE LEPTONS AND ANTILEPTONS IN EQUAL NUMBERS:

$$\epsilon_{N_i} = \sum_j \frac{\Gamma(N_i \rightarrow l_j \phi^*) - \Gamma(N_i \rightarrow l_j^* \phi)}{\Gamma_{N_i}^{TOT}}$$



REQUIRES CP-VIOLATION ← PROVIDED BY THE h_{ij} YUKAWA COUPLINGS

⇒ ϵ_{N_i} AT LOWEST ORDER: TREE LEVEL - ONE LOOP INTERFERENCE.

⇒ RESULTS:

$$E_{N_i} = \sum_k \frac{1}{8\pi} \frac{M_{Nk}}{M_{N_i}} \left[1 - \left(1 + \frac{M_{Nk}^2}{M_{N_i}^2} \right) \ln \left(1 + \frac{M_{N_i}^2}{M_{Nk}^2} \right) \right]$$

$$\frac{\sum_{j,l} \text{Im} [h_{ij}^* h_{il}^* h_{kj} h_{kl}]}{\sum_j |h_{ij}|^2}$$

⇒ IF THE OUT-OF-EQUILIBRIUM CONDITION IS SATISFIED:

$$\Gamma_{N_i} = \frac{1}{16\pi} \sum_j |h_{ij}|^2 M_{N_i} \lesssim H(T=M_{N_i}) \sim \frac{T^2}{M_{\text{PLANCK}}} \Big|_{T=M_{N_i}}$$

$$\text{THEN: } \frac{m_L}{\Lambda} \sim \sum_i \frac{E_{N_i}}{g^*} \quad (\sim \frac{m_B}{\Lambda})$$

TO HAVE $\Gamma_{N_i} \lesssim H$ AND $\frac{m_L}{\Lambda} \sim 10^{-10}$ WE NEED: $M_{N_i} \gtrsim 10^{10} \text{ GeV}$

FOR $M_{N_i} \gtrsim 10^{10} \text{ GeV}$ FROM THE SAME INTERACTIONS ARE GENERATED IN ADDITION THE ν MASSES FROM THE SEESAW MECHANISM: $m_\nu \sim \frac{h^2 v^2}{M_{N_i}}$

MOREOVER THIS MODEL IS QUITE NATURAL IN THE CONTEXT OF GUT SUCH AS SO(10)

⇒ VERY ATTRACTIVE MODEL!

BUT: WHAT ABOUT ITS TESTABILITY?

↳ INDIRECT TESTS:

- IN GENERAL: WITH $3 \nu_L + 3 N_R$:

6 PHASES

3 COMBINATIONS
IN THE LIGHT ν
MASS MATRIX

3 OTHER COMBINATIONS DECOUPLE
↳ e.g. NECESSARY FOR TESTING
LEPTOGENESIS

- SUPERSYMMETRIC SEESAW MODEL: NON DECOUPLING EFFECTS FROM THE h_{ij} AND M_{N_i} ON THE SOFT SLEPTON MASS TERMS ← RGE'S

⇒ INFLUENCE ON:

- e^- AND μ^- EDM
- T-ASYMMETRY IN $\mu \rightarrow e e e$
- SLEPTON OSCILLATIONS AT COLLIDERS
- $\mu \rightarrow e \gamma$, $\tau \rightarrow \mu \gamma$
- SNEUTRINO-ANTISNEUTRINO OSCILLAT° ($\Delta L=2$).
- ...



DAVIDSON-
IBARRA (01),
ELLIS ET AL (01)
BLAZEK-KING (00),
CASAS-
IBARRA (01)
LAVIGNAC ET AL
(01)

...

⇒ COMBINED WITH ν OSCILLAT° (AND $0\nu 2\beta$) EXPTS, CHANCE FOR A LOW ENERGY PROGRAM OF LEPTOGENESIS TESTS

- GUT WITH TEXTURES: MORE RELATIONS BETWEEN LEPTOGENESIS PARAMETERS AND LIGHT ν MASS PARAMETERS

BÜCHMULLER, PLUMACHER, (98...)

BRANCO, GONZALEZ-FELIPE,

JOACHIM, REBELO (01)

JOSHIPURA, PASCHOS,

RODEJOHANN (01)

...



NEED FOR MORE PRECISE INFORMATION ON THE ν MASS MATRIX (ν HIERARCHY...)

2) E_6 LEPTOGENESIS

T.H., E. MA, M. RAIDAL,
U. SARKAR, PLB512 (01)

* IN E_6 THE N_R ARE PREDICTED AS IN $SO(10)$

↳ IN THE FUNDAMENTAL REPRESENTATION: 27

$$\begin{array}{l} = 15 : \text{A STANDARD MODEL GENERATOR} \\ + 1 : N_R \\ + 4 : 2 \text{ } SU(2)_L \text{ DOUBLETS: } \tilde{H} \text{ AND } \tilde{H}^c \\ + 6 : 2 \text{ } SU(3)_C \text{ TRIPLETS: } \tilde{R} \text{ AND } \tilde{R}^c \\ + 1 : 1 \text{ } SU(2)_L \text{-} SU(3)_C \text{ SINGLET} \end{array} \left. \begin{array}{l} \} 16 \text{ OF } SO(10) \\ \} 10 \text{ OF } SO(10) \\ \} 1 \text{ OF } SO(10) \end{array} \right\}$$

* LOW ENERGY TESTS: MORE IN E_6

↳ FROM THE OBSERVATION OF A Z' OR A W' AT THE
~ FEW 100 GeV - FEW TeV SCALE \leftarrow FROM AN EXTRA
 $U(1)$ OR $SU(2)$.

• FOR $SO(10)$: A Z' OR W' WOULD EXCLUDE:

- THE STANDARD LEPTOGENESIS MECHANISM
- THE SEESAW MECHANISM \leftarrow IN $SO(10)$



IMPOSSIBLE TO HAVE A $SO(10)$
SUBGROUP AT \sim TeV WITHOUT
HAVING N_R LIGHT: $M_{N_R} \sim$ TeV.

- IN E_6 : AMONG THE MANY POSSIBLE $U(1)$ OR $SU(2)$ AT $\sim 1\text{TeV}$, THERE EXIST (ONLY) TWO CONSISTENT WITH LEPTOGENESIS AND SEESAW



WHICH ALLOW HEAVY N_R
 (i.e. FOR WHICH N_R ARE
 SINGLET $\Rightarrow M_{N_R}$ NOT PROTECTED)

$$\Rightarrow \left. \begin{array}{l} SU(3)_C \times SU(2)_L \times U(1)_Y \times \underline{U(1)_N} \\ \text{OR} \\ SU(3)_C \times SU(2)_L \times \underline{SU(2)_R^A} \times U(1)_{Y_L+Y_R'} \end{array} \right\} \begin{array}{l} \text{POSSIBLE GAUGE} \\ \text{SYMMETRIES AT} \\ \sim 1\text{TeV} \end{array}$$

$$(Q_N = \cos \alpha Q_Y + \sin \alpha Q_X \leftarrow \tan \alpha = \sqrt{\frac{7}{15}})$$

$E_6 \rightarrow SO(10) \times U(1)_\psi$ $SO(10) \rightarrow SU(5) \times U(1)_\chi$

$$SU(2)_R^A : (\ell^c, \nu^c) : \text{DOUBLET}$$

IN THESE 2 CASES THE MODEL WORKS FINE FOR
 GENERATING LEPTOGENESIS AND ν MASSES

WITH IN ADDITION A RICH PHENOMENOLOGY AT $\sim 1\text{TeV}$:

- ONE Z' FOR $U(1)_N$ AND $1Z' + 2W_R$ FOR $SU(2)_R^A$

$$\text{EXAMPLE: } \Gamma(Z' \rightarrow \nu \bar{\nu} + S \bar{S}) = \frac{62}{15} \Gamma(Z' \rightarrow \ell^+ \ell^-) \leftarrow U(1)_N$$

$$m_{W_R} = \frac{\cos 2\theta_W}{\cos \theta_W} m_{Z'} = 0.84 m_{Z'} \leftarrow SU(2)_R^A$$

- 11 EXTRA PARTICLES FROM THE 27 AT $\sim 1\text{TeV}$

\Rightarrow WOULD PROVIDE STRONG INDICATIONS FOR THE ORIGIN
 OF LEPTOGENESIS AND ν MASSES THROUGH THE N_R

BUT:

IS IT REALLY IMPOSSIBLE
TO CONSTRUCT A LEPTOGENESIS
MODEL AT A DIRECTLY TESTABLE
SCALE ($\sim 1-10$ TeV) ???

3) PROBLEMS TO INDUCE LEPTOGENESIS AT $\sim 1-10$ TeV

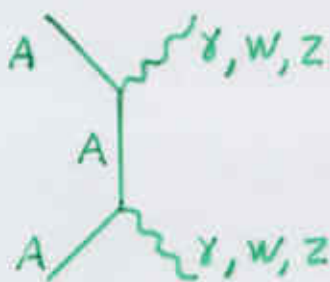
- A) NON EQUILIBRIUM CONDITION: $\frac{\Gamma_A}{T} < H(T=M_A)$
 $\frac{?}{T} < \frac{?}{T^2}$

↳ MUCH MORE DIFFICULT TO SATISFY
AT $T \sim 1-10$ TeV THAN AT $T \sim 10^{10}-10^{15}$ GeV

↳ FOR EXAMPLE: AT $T \sim 10$ TeV, $\Gamma_A < H$ REQUIRES $h \lesssim 10^{-6}$
 $\Rightarrow \frac{\Gamma_A}{\Lambda} \sim \frac{\Sigma N}{g_*} \sim \frac{1}{g_*} \frac{1}{8\pi} h^2 \sim \underline{\underline{10^{-15}!}}$

- B) SCATTERING DAMPING EFFECTS:

↳ IF THE DECAYING PARTICLE A IS NOT NEUTRAL AND $SU(2)_L$ SINGLET, THE GAUGE SCATTERINGS ARE VERY FAST



\Rightarrow NO DEPARTURE FROM THERMAL
EQUILIBRIUM FOR A

\Rightarrow ASYMMETRY DAMPED BY ~ 5
ORDERS OF MAGNITUDE AT ~ 10 TeV

↳ WE NEED A NEUTRAL AND $SU(2)_L$ SINGLET DECAYING
PARTICLE

• C NEUTRINO MASS CONSTRAINTS:

↳ VALUES OF THE L VIOLATING COUPLINGS REQUIRED TO INDUCE THE m_ν ARE TOO LARGE TO SATISFY THE CONDITION $\Gamma < H$ AT 1-10 TeV

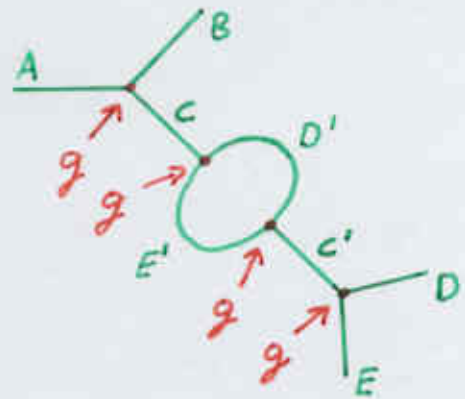
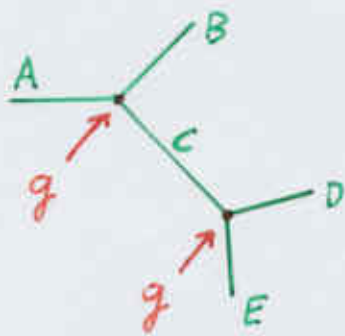
↳ GIVES A SUPPRESSION OF $m_{L1\Delta}$

↓
OF ORDER ~ 100
IN THE SEESAW
MODEL WITH N_{Ri}

↓
OF ORDER $\sim 10^6$ IN THE
RADIATIVE ν MASS MODELS
(WITH R-PARITY VIOLAT^o, ZEE MODEL)

4) A SIMPLE SOLUTION WITH 3-BODY DECAYS

T.H., NPB633 (2002)



$$\Rightarrow \Gamma_A \sim g^4 \text{ AND } \epsilon_A \sim \frac{g^6}{g^4} \sim g^2$$

2 BODY DECAYS

$$\Gamma \sim h^2 < H \Rightarrow h^2 \lesssim 10^{-12}$$

↑
AT $T \sim 10 \text{ TeV}$

$$\Rightarrow \epsilon \sim h^2 \lesssim \underline{\underline{10^{-12}}}$$

\Rightarrow NEED FOR A MUCH HIGHER SCALE

3 BODY DECAYS

$$\Gamma \sim g^4 < H \Rightarrow g^4 \lesssim 10^{-12}$$

↑
AT $T \sim 10 \text{ TeV}$

\Downarrow
 $g \lesssim 10^{-3}$

$$\Rightarrow \epsilon \sim g^2 \sim \underline{\underline{10^{-6}}}$$

\Rightarrow VERY NATURAL AT 1-10 TeV

MOREOVER: A VALUE OF A L-VIOLATING COUPLING g OF ORDER 10^{-3} IS TYPICALLY THE VALUE WE NEED TO INDUCE RADIATIVELY ν MASSES OF ORDER $\sim 0.1-1 \text{ eV}$.

⇒ IF THE LAST L-VIOLATING DECAY IS A 3-BODY DECAY
AT $T \sim 1-10 \text{ TeV}$, LEPTOGENESIS AS WELL AS ν MASSES
COULD BE GENERATED NATURALLY AT THIS SCALE

⇒ ALTERNATIVE AND TESTABLE GENERAL FRAMEWORK
OF LEPTOGENESIS AND ν MASSES AT THE $\sim 1-10 \text{ TeV}$ SCALE
% USUAL FRAMEWORK AT THE GUT SCALE.

5] AN EXPLICIT MODEL

T.H., NPB 633 (2002)

• WE NEED DECAYS OF A NEUTRAL $SU(2)_L$ SINGLET: WE TAKE 3 N_{R_i}

• N_{R_i} 3-BODY DECAYS REQUIRE HEAVIER VIRTUAL PARTICLES:

WE TAKE CHARGED SCALAR $SU(2)_L$ SINGLETS: $S_{1,2}^\pm$

+ WE TAKE 2 HIGGS DOUBLETS: $H_{1,2}$

• PRELIMINARY CONDITION: N_{R_i} WITH $M_{N_i} \sim 1-10 \text{ TeV}$

REQUIRE TINY $h_{ij} \bar{L}_i N_{R_j} H_{1,2}$

YUKAWA COUPLINGS TO AVOID

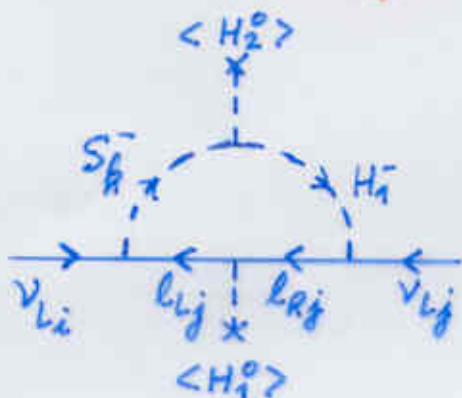
$m_\nu \gtrsim \text{eV}$.

\Rightarrow WE ASSUME THEM NEGLIGIBLE

• YUKAWA INTERACTIONS:

$$\mathcal{L} \ni \underbrace{f_{ij}^k}_{\uparrow} \bar{L}_i^T C^{-1} \gamma_2 L_j S_k^- + \underbrace{\lambda_S^k}_{\uparrow} H_1^T \gamma_2 H_2 S_k^- + \underbrace{h_{Rij}^k}_{\uparrow} \bar{L}_{Rj}^T C^{-1} N_{Ri} S_k^-$$

COUPLINGS OF THE ZEE MODEL \leftarrow VIOLATE L

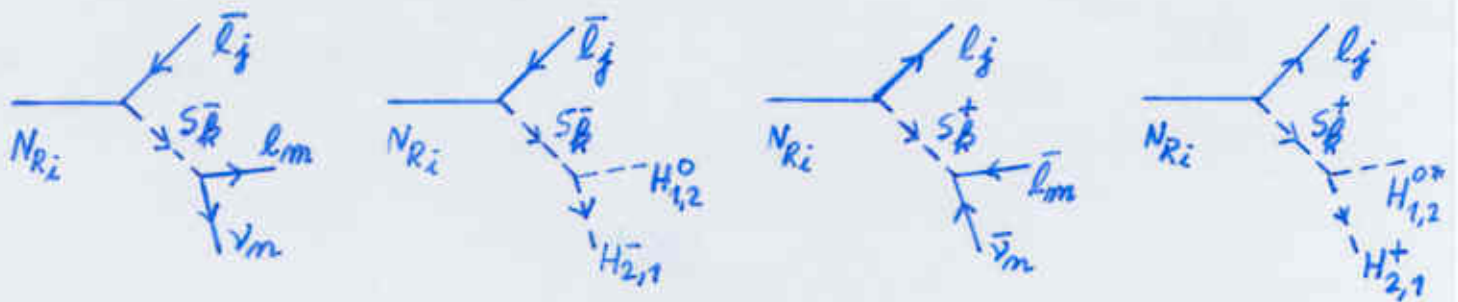


$$\Rightarrow (m_\nu)_{ij} = \frac{1}{(4\pi)^2 k} \sum \frac{\lambda_S^k}{m_{S_k}^2 - m_{H^\pm}^2} \frac{\sqrt{2}}{v_1} f_{ij}^k$$

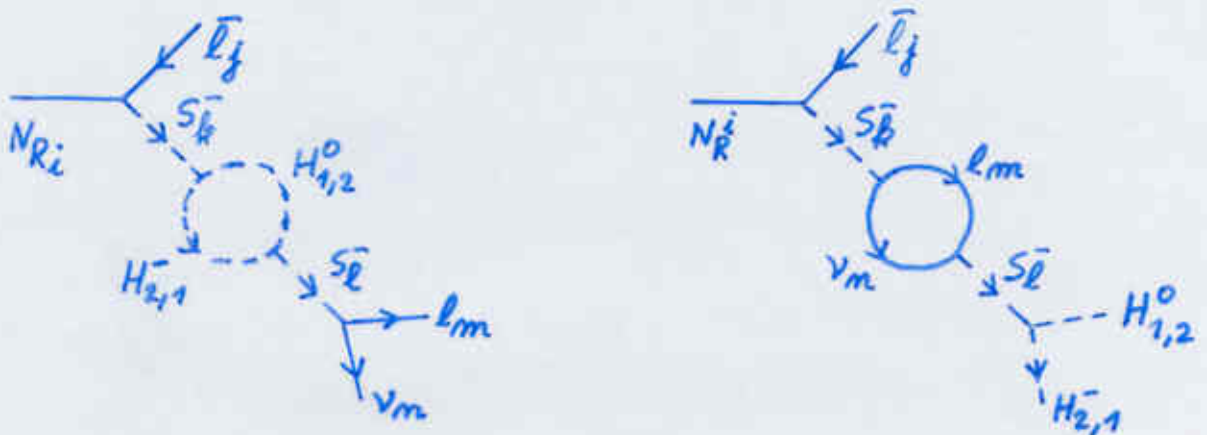
$$\cdot (m_{L_j}^2 - m_{L_i}^2) \ln \frac{m_{S_k}^2}{m_{H^\pm}^2}$$

\Rightarrow A MINIMAL WAY TO INDUCE THE ν MASSES.

• LEPTOGENESIS:



⇒ THE LEPTONIC ASYMMETRY IS OBTAINED FROM THE INTERFERENCE WITH:



$$\Rightarrow \epsilon_{N_i} = \frac{1}{\Gamma_{N_i}} \frac{1}{(2\pi)^3} \frac{1}{12} \frac{\pi}{(4\pi)^2} \frac{M_{N_i}^5}{m_{S_1}^2 m_{S_2}^2} \cdot \sum_{j,m,n}$$

$$\begin{aligned} & \cdot \left[\text{Im} \left[h_{Rij}^{2*} h_{Rij}^1 \lambda_s^{1*} \lambda_s^2 \right] \left(\frac{|f_{mn}^1|^2}{m_{S_1}^2} - \frac{|f_{mn}^2|^2}{m_{S_2}^2} \right) \right. \\ & + \text{Im} \left[h_{Rij}^2 h_{Rij}^{1*} f_{mn}^1 f_{mn}^{2*} \right] \left(\frac{|\lambda_s^1|^2}{m_{S_1}^2} - \frac{|\lambda_s^2|^2}{m_{S_2}^2} \right) \\ & \left. + \text{Im} \left[f_{mn}^2 f_{mn}^{1*} \lambda_s^1 \lambda_s^{2*} \right] \left(\frac{|h_{Rij}^1|^2}{m_{S_1}^2} - \frac{|h_{Rij}^2|^2}{m_{S_2}^2} \right) \right] \end{aligned}$$

$$\Rightarrow E_{N_i} \sim \frac{1}{16\pi} \frac{\lambda_S^2 \beta^2}{\lambda_S^2 + \beta^2 M_{N_i}^2} \frac{M_{N_i}^2}{m_S^2}$$

WITH:

$$\Gamma_{N_i} \sim \frac{1}{(2\pi)^3} \frac{1}{48} h_R^2 (\lambda_S^2 + \beta^2 M_{N_i}^2) \frac{M_{N_i}^3}{m_{S_1}^2 m_{S_2}^2}$$

\Rightarrow IT IS EASY TO GET $\Gamma < H$ AND $\frac{m_L}{\Lambda} \sim 10^{-10}$:

EXAMPLE OF A SET OF VALUES WHICH GIVES:

$m_{\nu_1} = 0.1 \text{ eV}$, $\delta m_{\text{atmos}}^2$, $\delta m_{\text{SOLAR-LMA}}^2$, $m_L/\Lambda \sim 10^{-9}$:

$$M_N = 1 \text{ TeV}, \quad m_{S_1} = 4 \text{ TeV}, \quad m_{S_2} = 5 \text{ TeV}$$

$$\lambda_S = 50 \text{ GeV}, \quad h_R = 10^{-4}, \quad \beta_{\text{eff}} = 0.2 \quad (= 3 \cdot 10^2 \beta_{\text{eff}} = 3 \cdot 10^5 \beta_{\text{eff}})$$

SUMMARY

- IN CONSTRAINED MODELS (SUSY SEESAW WITH UNIVERSALITY, GUT WITH TEXTURES) WE MIGHT GET A CHANCE TO TEST LEPTOGENESIS

↳ FROM ν MASS CONSTRAINTS, e^- AND μ^- EDM...

- THE OBSERVATION OF A Z' OR W' WOULD PROVIDE Δ INFORMATION FOR UNDERSTANDING THE ORIGIN OF LEPTOGENESIS

- WE ARE NOT OBLIGED TO GO TO A HIGH SCALE

↳ IF THE LAST L-VIOLATING DECAY WAS A 3-BODY DECAY AT 1-10 TeV, LEPTOGENESIS COULD OCCUR NATURALLY AND BE TESTED AT THIS SCALE

↳ THIS PROVIDES A WAY TO RECONCILE THE RADIATIVE ν MASS MODELS WITH LEPTOGENESIS