

Non-standard Physics probed by Neutrinos



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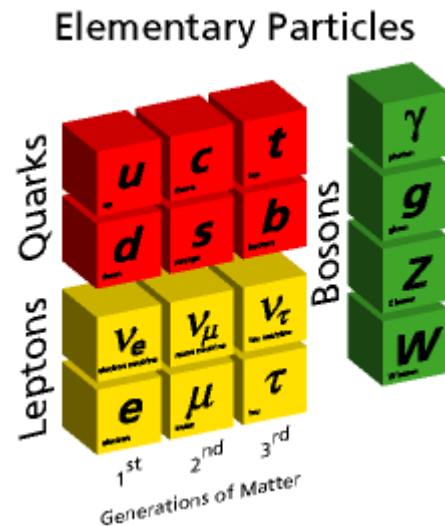
XXIIIrd rencontres de Blois, chateau de Blois, May 29 – June 3, 2011

Physics Beyond the Standard Model

QED \rightarrow QCD \rightarrow SM

$U(1)_{em}$ $SU(3)_C$ $SU(3)_C \times SU(2)_L \times U(1)_Y$

Success story of
d=4 renormalizable
QFTs



TOE solving all problems does not (yet) exist \rightarrow solve some problems \rightarrow increasing levels of speculation:

- 1) new fields
- 2) extend gauge group
- 3) new concepts (SUSY, ...)
- 4) wild speculations

\rightarrow BSM from a neutrino perspective

Theoretical reasons for BSM:

SM does not exist without cutoff (triviality)

Higgs-doublet = only simplest extension

Gauge hierarchy problem

Gauge unification, charge quantization

Strong CP problem

Unification with gravity

Why: 3 generations, which representations

Many parameters (9+? masses, 4+? mixings)

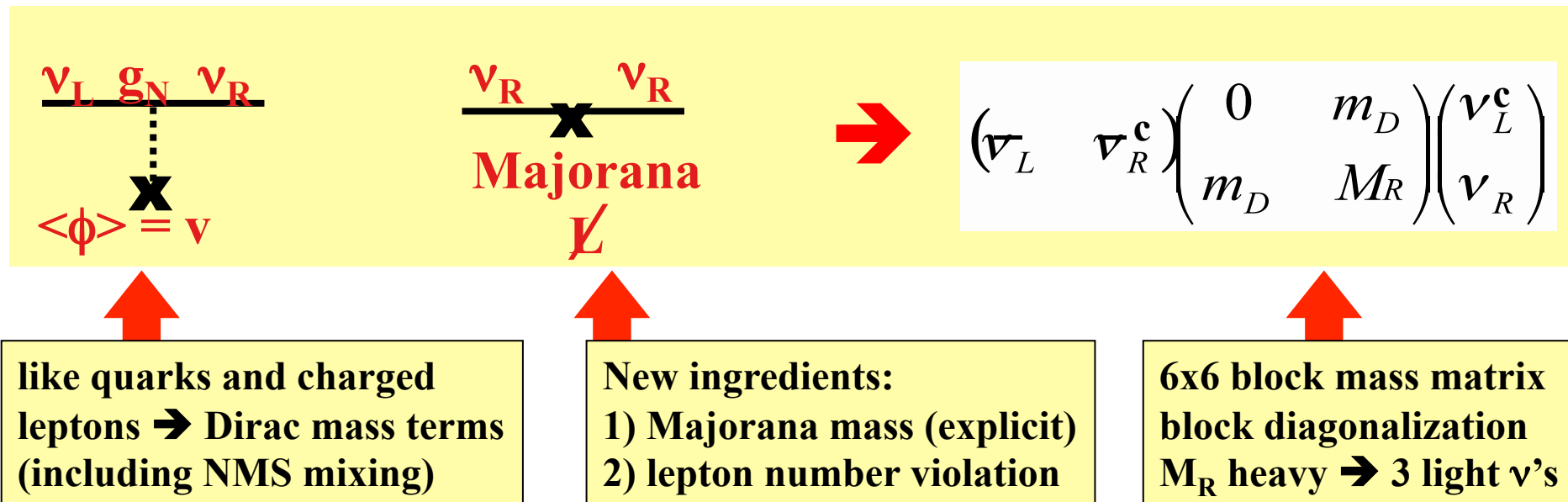
Experimental BSM facts:

- Electro weak scale \ll Planck scale
- Gauge couplings almost unify
- Neutrinos have masses & large mixings
- Baryon asymmetry of the Universe
- Dark Matter, Dark Energy, few $\geq 2\sigma$ hints?

New Physics: Neutrino Mass Terms

1) Simplest possibility:
add 3 right handed
neutrino fields

Field	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
$L_Q = \begin{pmatrix} l_u \\ l_d \end{pmatrix}$	3	2	1/3
r_u	3	1	4/3
r_d	3	1	-2/3
$L_L = \begin{pmatrix} l_\nu \\ l_e \end{pmatrix}$	1	2	-1
$r_\nu ???$	1	1	0
r_e	1	1	-2



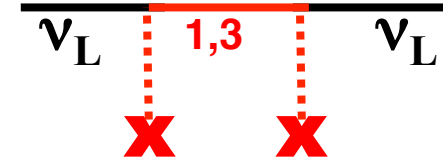
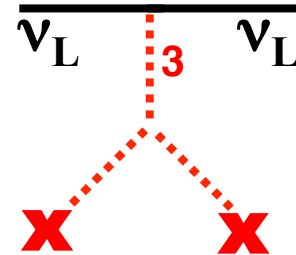
NEW ingredients, 9 parameters \rightarrow SM+

2) Maybe 3+N right handed neutrino fields

→ (6+N) x (6+N) mass matrix

→ how many of the 6+N eigenvalues are light (also for N=0)

3) new: scalar triplets ($\underline{3}_L$) or fermionic $\underline{1}_L$ or $\underline{3}_L$



→ left-handed Majorana mass term:

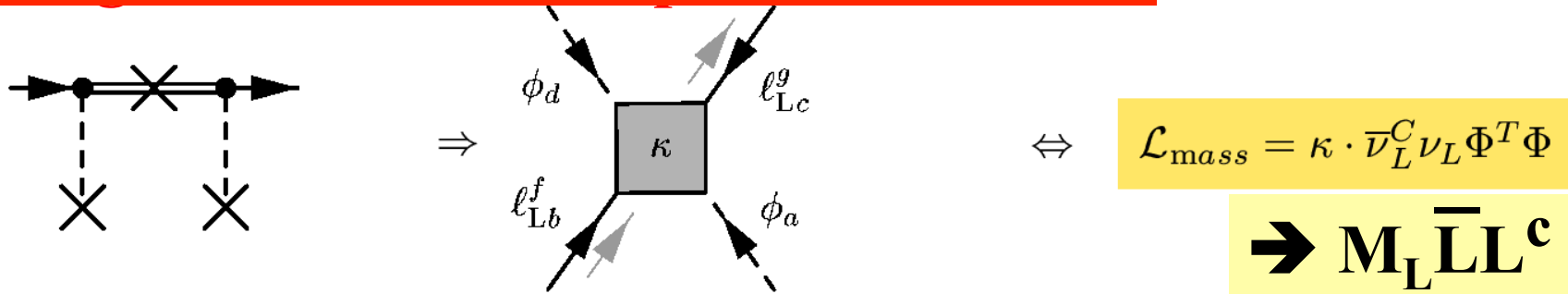
$$\rightarrow M_L \bar{L} L^c$$

4) Both $\underline{\nu}_R$ and new singlets / triplets:

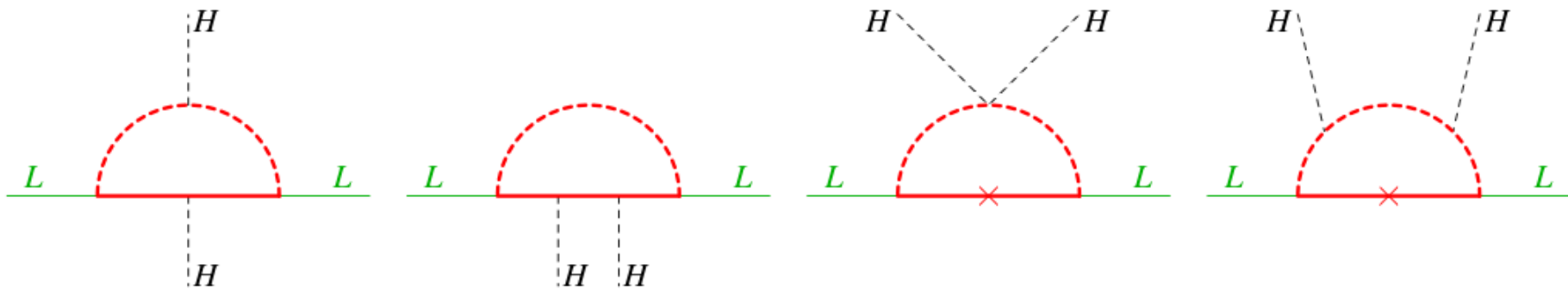
→ see-saw type II, III

$$m_\nu = M_L - m_D M_R^{-1} m_D^T$$

5) Higher dimensional operators: d=5, ...



6) Radiative neutrino mass generation



7-N) SUSY, extra dimensions, ...

Other effective Operators Beyond the SM

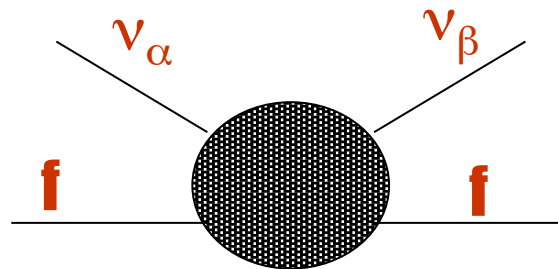
→ effects beyond 3 flavours

→ **Non Standard Interactions = NSIs** → effective 4f operators

$$\mathcal{L}_{NSI} \simeq \epsilon_{\alpha\beta} 2\sqrt{2}G_F (\bar{\nu}_{L\beta} \gamma^\rho \nu_{L\alpha}) (\bar{f}_L \gamma_\rho f_L)$$

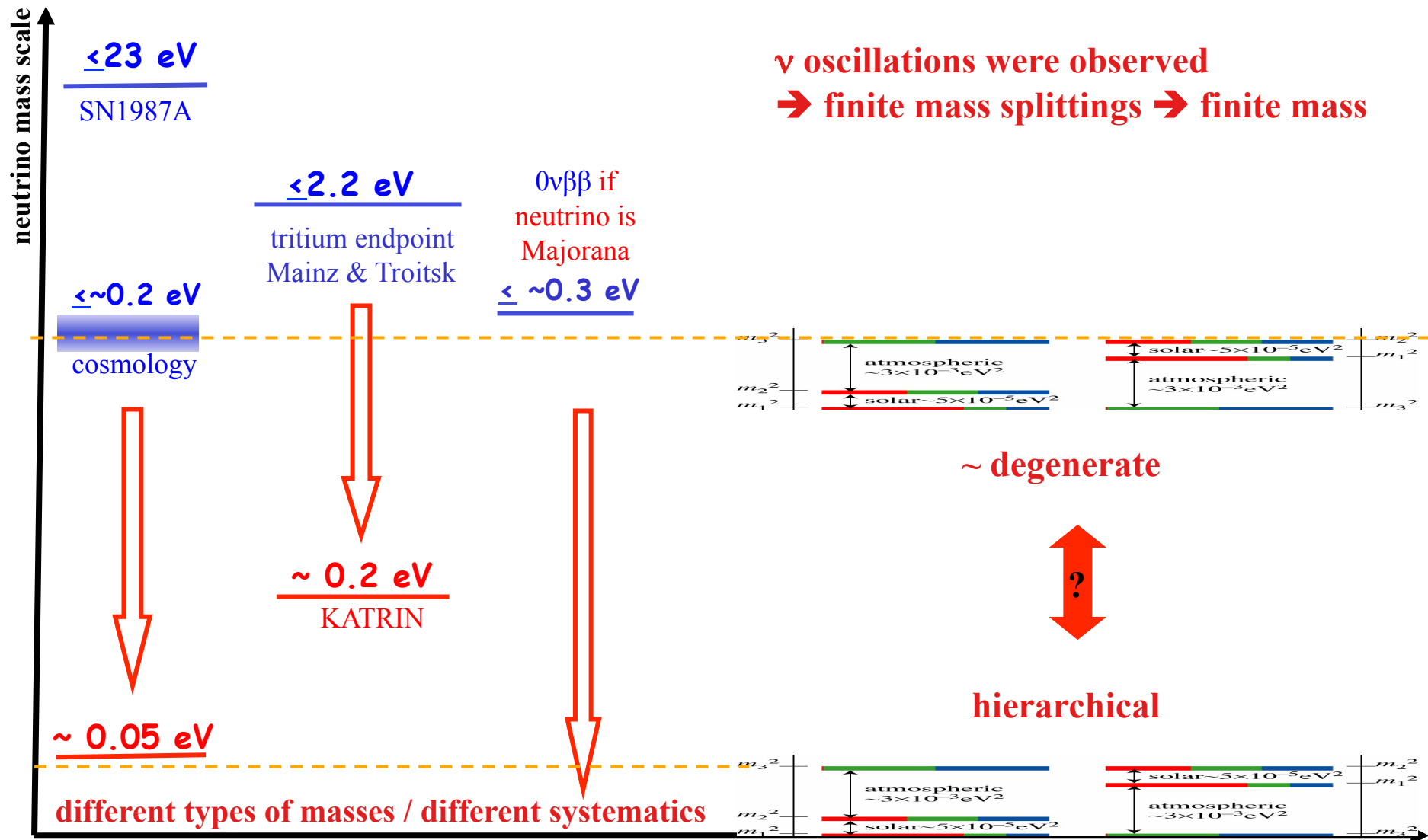
• **integrating out heavy physics (c.f. $G_F \leftrightarrow M_W$)**

$$|\epsilon| \simeq \frac{M_W^2}{M_{NSI}^2}$$



Grossman, Bergmann+Grossman, Ota+Sato, Honda et al., Friedland+Lunardini, Blennow+Ohlsson+Skrotzki, Huber+Valle, Huber+Schwetz+Valle, Campanelli+Romanino, Bueno et al., Barranco+Miranda+Rashba, Kopp+ML+Ota, ...

Overview of Neutrino Mass Knowledge



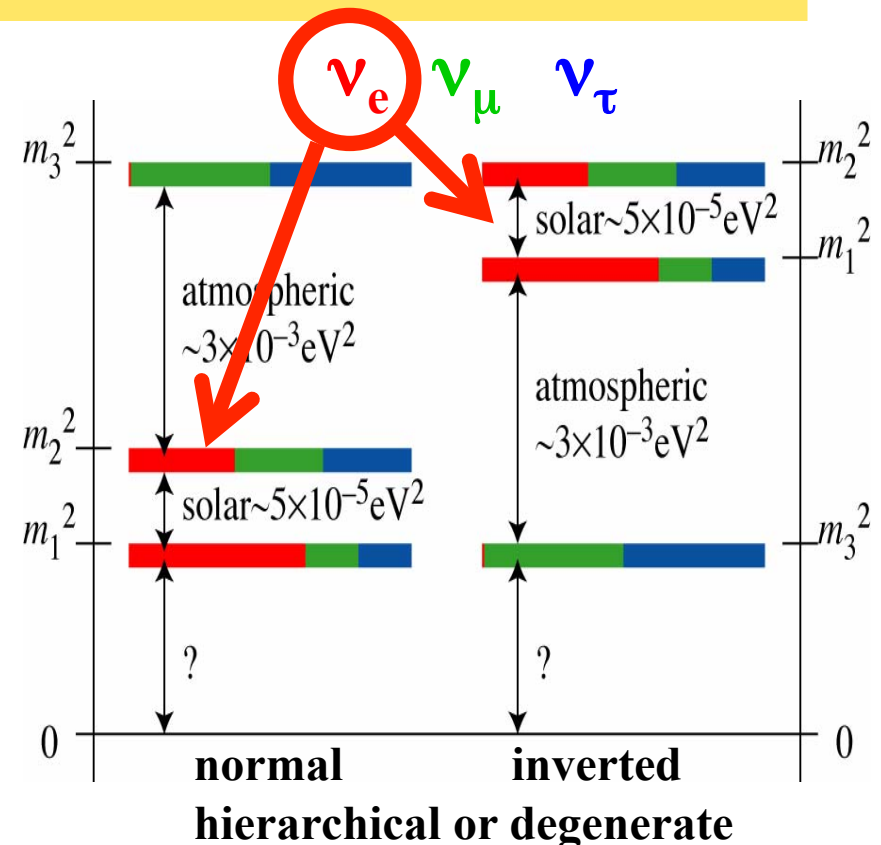
3 Light Neutrinos (...assumed)

Mass & mixing parameters: m_1 , Δm_{21}^2 , $|\Delta m_{31}^2|$, $\text{sign}(\Delta m_{31}^2)$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \text{diag}(e^{i\alpha}, e^{i\beta}, 1)$$

questions:

- Dirac / Majorana
- mass scale: m_1
- mass ordering: $\text{sgn}(\Delta m_{31}^2)$
- how small is θ_{13} , θ_{23} maximal?
- leptonic CP violation
- 3 flavour unitarity?
- why 3 generations, why $d=4$, ...



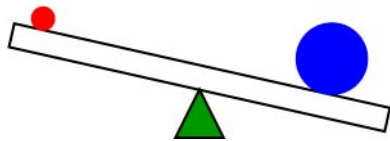
Suggestive Seesaw Features

QFT: natural value of mass operators \leftrightarrow scale of symmetry

$m_D \sim$ electro-weak scale

$M_R \sim$ L violation scale $\leftarrow? \rightarrow$ embedding (GUTs, ...)

See-saw mechanism (type I)



$$m_\nu = m_D M_R^{-1} m_D^T$$

$$m_h = M_R$$

Numerical hints:

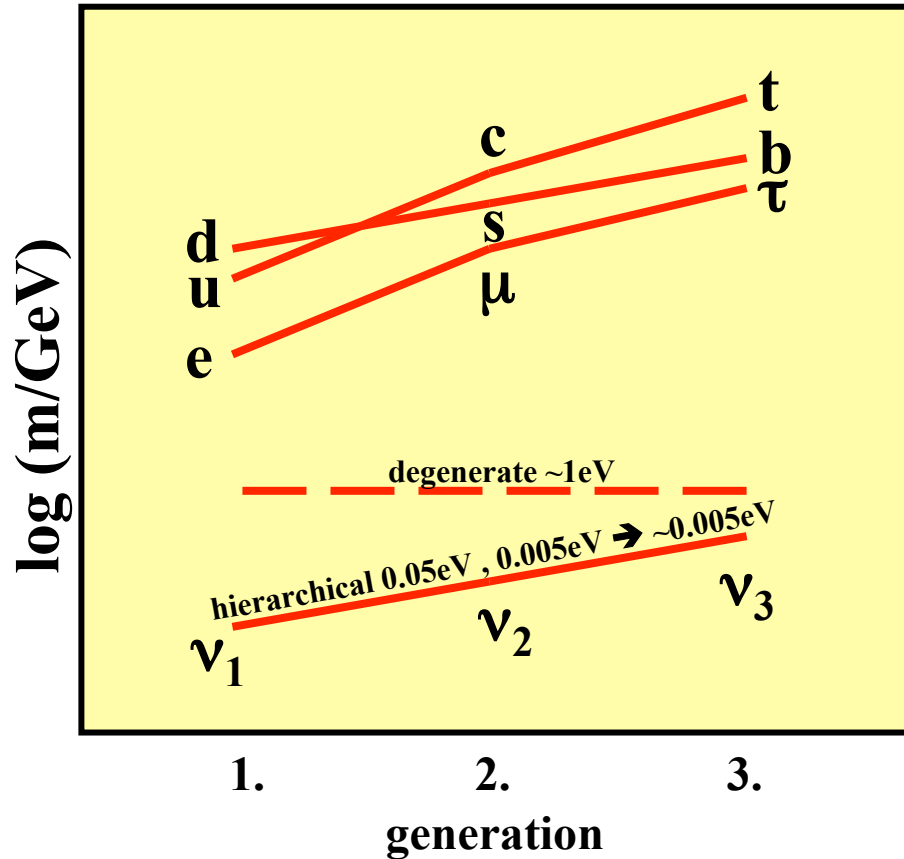
For $m_3 \sim (\Delta m_{\text{atm}}^2)^{1/2}$, $m_D \sim$ leptons $\rightarrow M_R \sim 10^{11} - 10^{16} \text{ GeV}$

$\rightarrow \nu$'s are **Majorana particles**, m_ν probes \sim GUT scale physics!

\rightarrow smallness of $m_\nu \leftrightarrow$ high scale of L , symmetries of m_D, M_R

2nd Look Questions

Quarks & charged leptons → hierarchical masses → neutrinos?



Quarks and charged leptons:

$$m_D \sim H^n ; n = 0, 1, 2 \rightarrow H \geq 20 \dots 200$$

Neutrinos:

$$m_\nu \sim H^n \rightarrow H \leq \sim 10$$

See-saw:

$$m_\nu = -m_D^T M_R^{-1} m_D$$

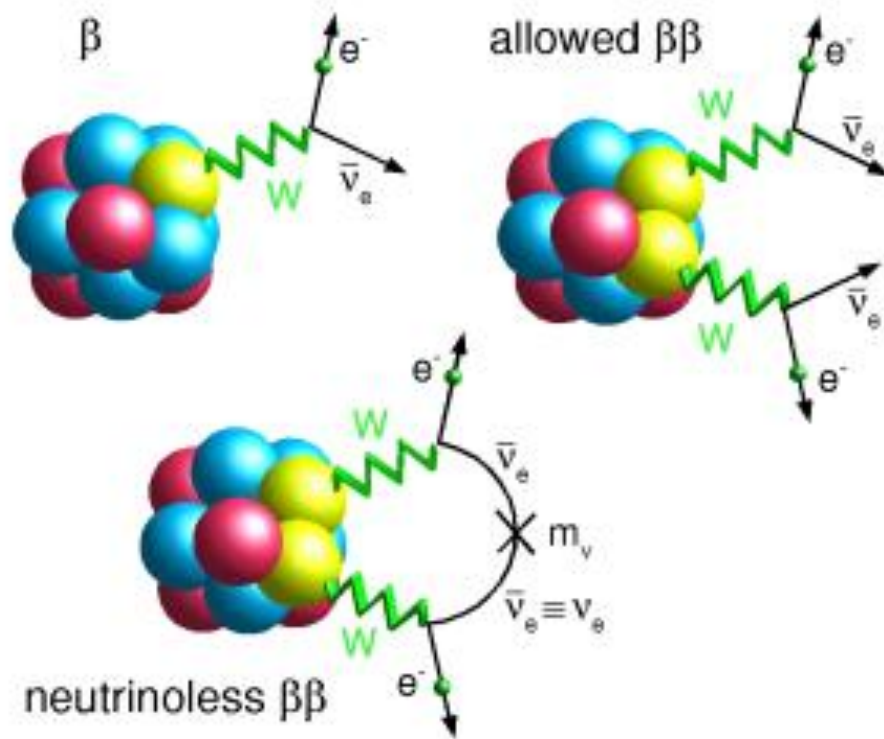


H	~ 10	≥ 20	?	≥ 20
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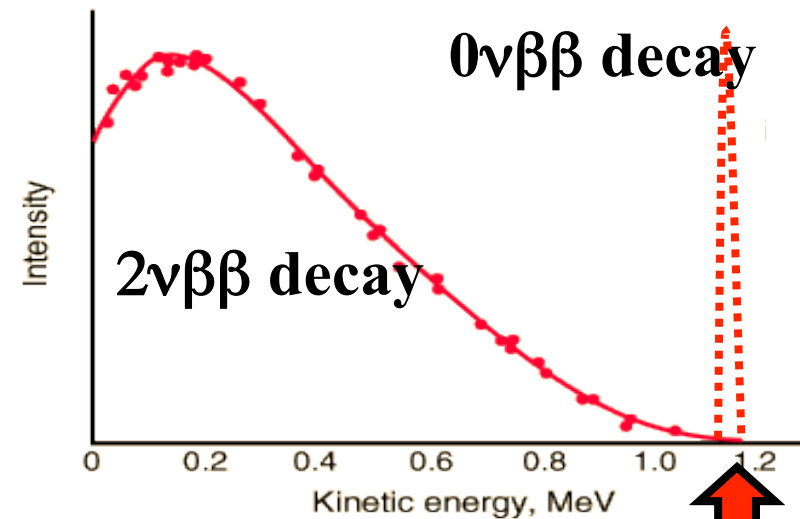
- » less hierarchy in m_D or corr. hierarchy in M_R ? → theoretically not connected!
- » other version of see-saw? → type II, III, ...?
- » Dirac masses?

What do we actually measure?

Neutrino-less Double Beta Decay



$2\nu\beta\beta$ decay of ^{76}Ge observed:
 $\tau = 1.5 \times 10^{21}$ y



Majorana $\nu \rightarrow 0\nu\beta\beta$ decay

warning: other lepton number violating processes may exist...

- signal at known Q-value
- $2\nu\beta\beta$ background (resolution)
- nuclear backgrounds
- ➔ use different nuclei

Neutrino Masses from Double β -Decay

Beta particle (electron)

Majorana $\nu \rightarrow 0\nu 2\beta$ decay

$\propto |\langle m_{ee} \rangle| = |\sum m_i U_{ei}^2| \leq 0.35 \text{ eV ?}$

Heidelberg-Moscow experiment

$$m_{ee} = |m_{ee}^{(1)}| + |m_{ee}^{(2)}| \cdot e^{i\Phi_2} + |m_{ee}^{(3)}| \cdot e^{i\Phi_3}$$

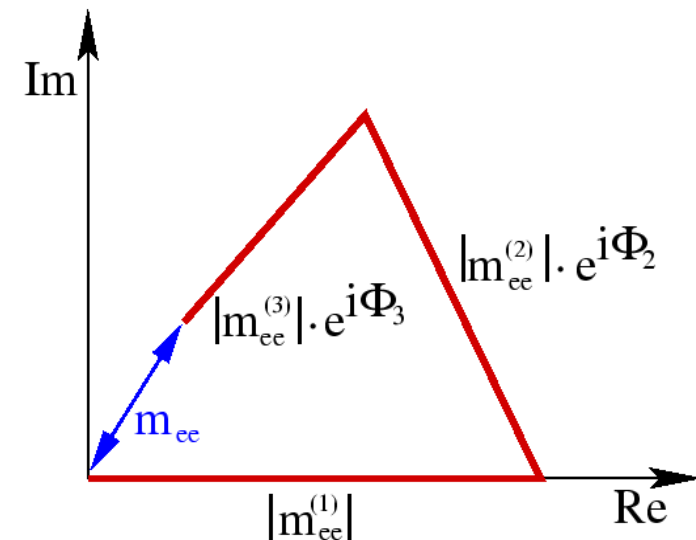
$$|m_{ee}^{(1)}| = |U_{e1}|^2 m_1$$

$$|m_{ee}^{(2)}| = |U_{e2}|^2 \sqrt{m_1^2 + \Delta m_{21}^2}$$

$$|m_{ee}^{(3)}| = |U_{e3}|^2 \sqrt{m_1^2 + \Delta m_{31}^2}$$

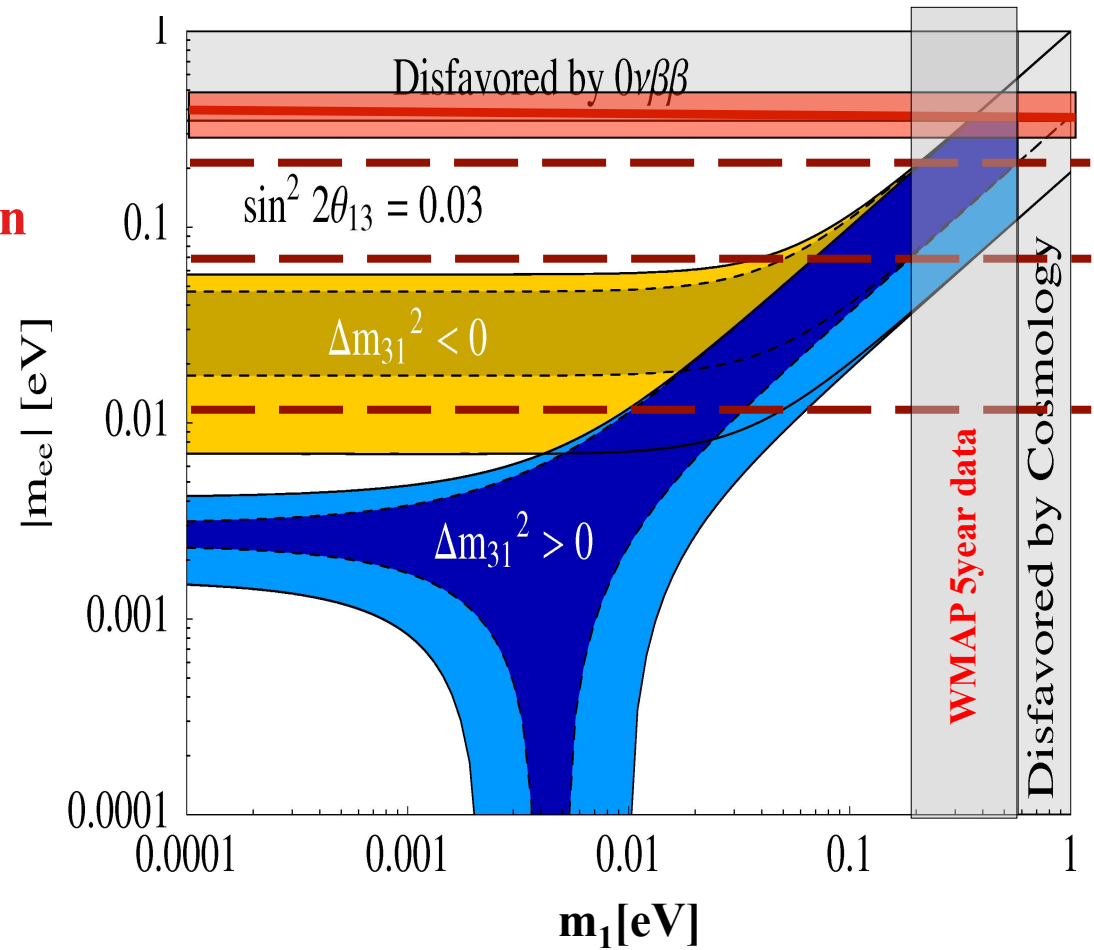
solar $\Rightarrow |U_{e1}|^2, |U_{e2}|^2, \Delta m_{21}^2$ atmosph. $\Rightarrow |\Delta m_{31}^2|$ CHOOZ $\Rightarrow |U_{e3}|^2 < 0.05$

\rightarrow free parameters: $m_1, \text{sign}(\Delta m_{31}^2), \text{CP-phases } \Phi_2, \Phi_3$



Claim of part of the original Heidelberg-Moscow collaboration
 \leftrightarrow cosmology \rightarrow ,tension‘

- aims of new experiments:**
- test HM claim
 - $(\Delta m_{31}^2)^{1/2} \simeq 0.05\text{eV} \pm \text{errors}$
 - \rightarrow reach 0.01eV
 - \rightarrow CUORE
 - \rightarrow GERDA phases I, II, (III)



Comments:

- cosmology: limitation by systematical errors \rightarrow ~another factor 5?
- $0\nu\beta\beta$ nuclear matrix elements ~factor 1.3-2 **theoretical** uncertainty in m_{ee}
- $\Delta m^2 > 0$ allows complete cancellation
 - \rightarrow $0\nu\beta\beta$ signal not guaranteed, but cancelation appears unlikely

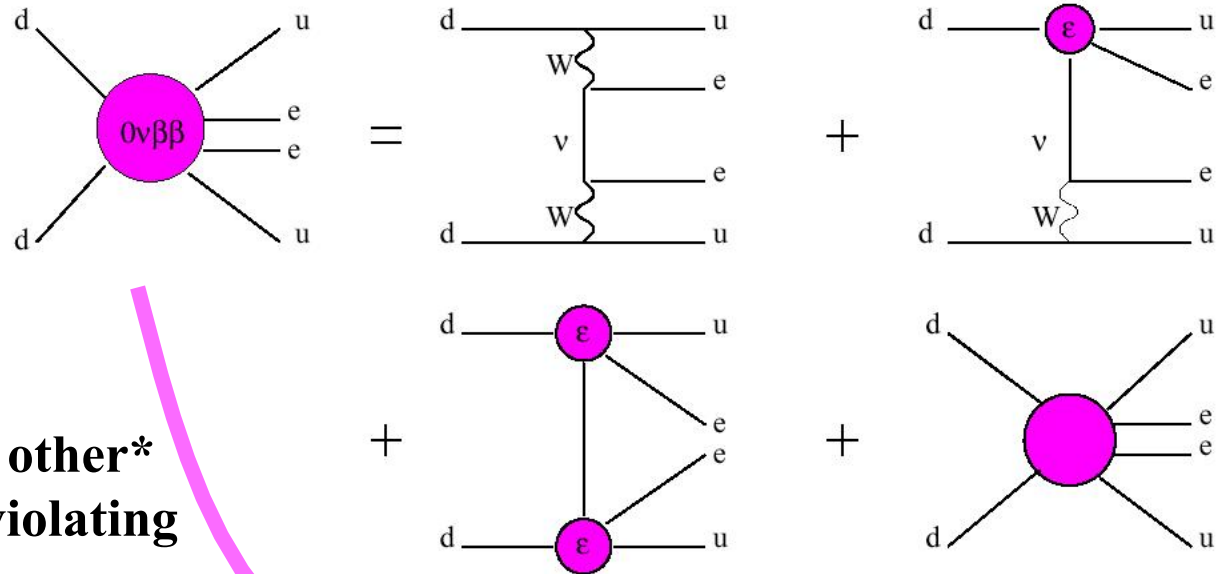
$0\nu\beta\beta$ from Alternative $\Delta L=2$ Operators

Various possibilities:

- LR symmetry
- SUSY (RPV)
- ...

→ $0\nu\beta\beta$ signal from *some other* new BSM lepton number violating operator

→ very promising interplay of neutrino mass determinations, cosmology, LHC, LNF experiments and theory

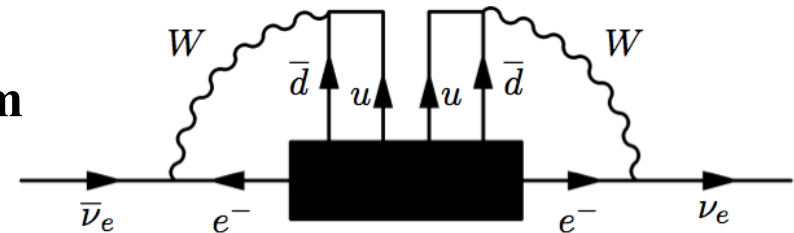


Schechter + Valle: Any $\Delta L=2$ violating operator

→ radiative generation of Majorana mass term

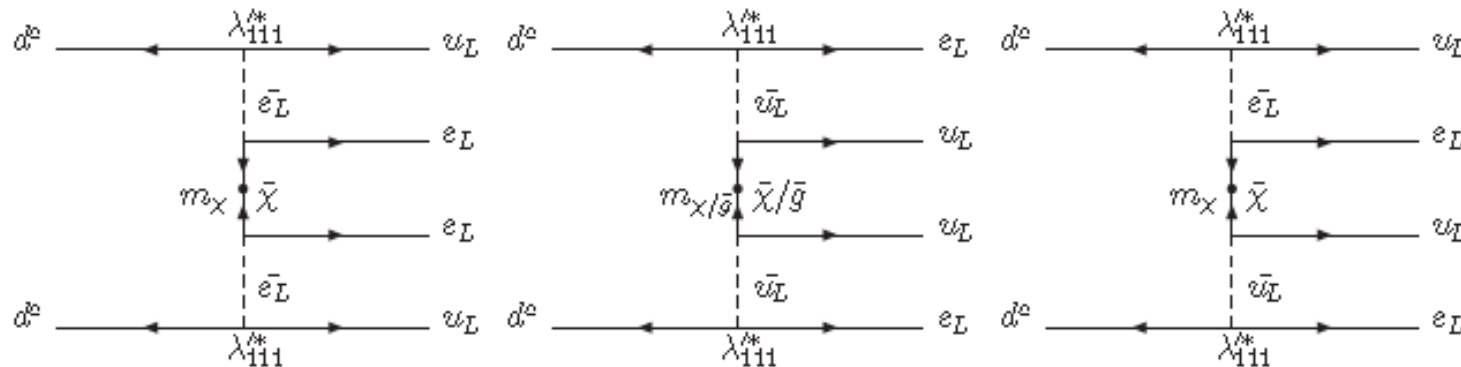
→ Majorana nature of ν 's guaranteed

→ but how big is the mass?



SUSY Example

Direct, TeV scale short range mediation w/o intermediate light ν , e.g.

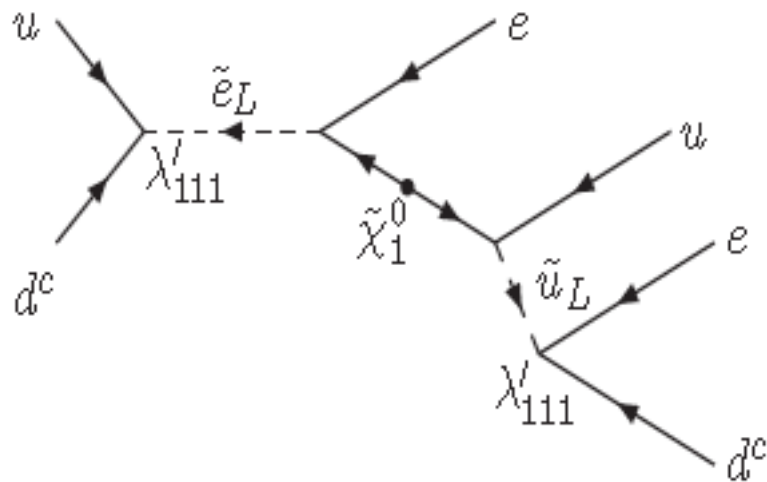


$$\mathcal{L}_{\lambda_{111}^* \lambda_{111}^*}^{eff, \Delta L_e=2}(x) = \frac{G_F^2}{2} m_p^{-1} [\bar{e}(1 + \gamma_5)e^c] \times \left[(\epsilon_{\tilde{g}} + \epsilon_\chi)(J_{PS} J_{PS} - \frac{1}{4} J_T^{\mu\nu} J_{T\mu\nu}) + (\epsilon_{\chi\tilde{e}} + \epsilon'_{\tilde{g}} + \epsilon_{\chi\tilde{f}}) J_{PS} J_{PS} \right]$$

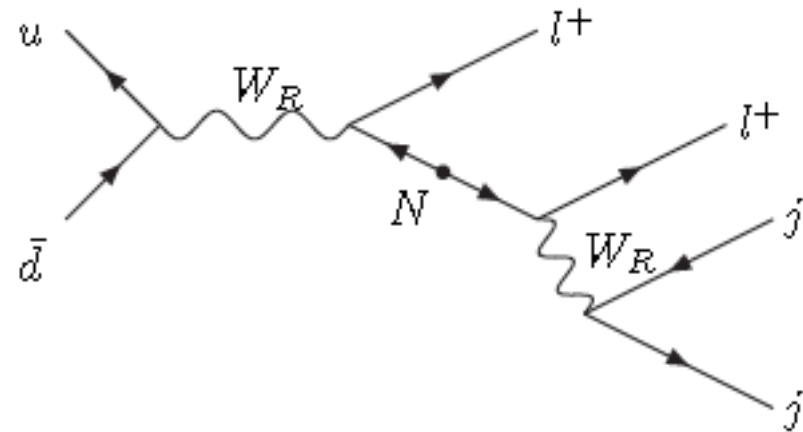
$$\epsilon_i \sim \pi \alpha_{(Strong, EW)} \frac{\lambda_{111}^2}{G_F^2} \frac{m_p}{m_{(\tilde{g}, \tilde{\chi})}} \frac{1}{m_{(\tilde{u}, \tilde{d}, \tilde{e})}^4}.$$

$\Delta L=2$ Operators and TeV Scale Physics

SUSY: direct test of λ'_{111}

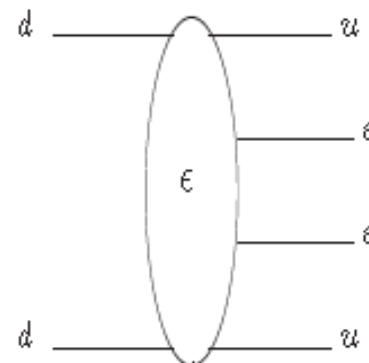
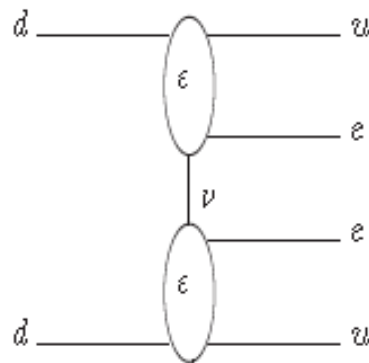


L-R symmetry: heavy N's



Relative strength of 'light' and 'heavy' $0\nu\beta\beta$ amplitudes:

$$M_{\text{light}} \sim G_F^2 \frac{m_{\beta\beta}}{\langle k^2 \rangle}$$



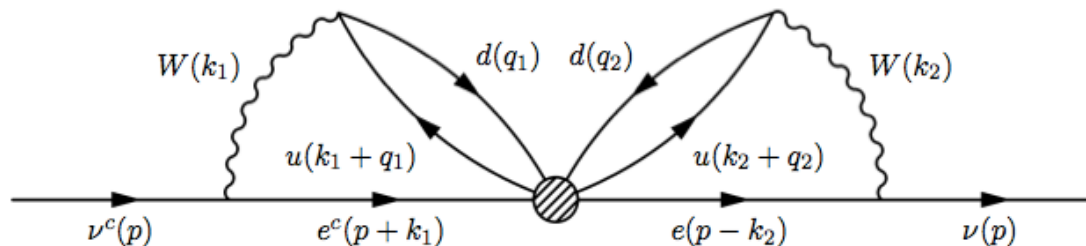
$$M_{\text{heavy}} \sim G_F^2 \left(\frac{\lambda}{g_2} \right)^4 \frac{M_W^4}{\Lambda^5}$$

SV-induced Neutrino Masses

General Lorentz-invariant Lagrangian for $0\nu\beta\beta$ (point operator)

$$\mathcal{L} = \frac{G_F^2}{2} m_p^{-1} (\epsilon_1 J J j + \epsilon_2 J^{\mu\nu} J_{\mu\nu} j + \epsilon_3 J^\mu J_\mu j + \epsilon_4 J^\mu J_{\mu\nu} j^\nu + \epsilon_5 J^\mu J j_\mu)$$

$$J = \bar{u} (1 \pm \gamma_5) d, \quad J^\mu = \bar{u} \gamma^\mu (1 \pm \gamma_5) d \text{ etc.}$$



Outcome:

M. Dürr, ML, A. Merle, arXiv:1105.0901

If other $\Delta L=2$ physics drives $0\nu\beta\beta \rightarrow$ SV gives $\delta m_\nu = 10^{-24}$ eV

\rightarrow mass correction too small to explain observed masses and splittings

\rightarrow explicit neutrino mass operators required

Dirac: $0\nu\beta\beta$ essentially unrelated to neutrino masses \leftrightarrow other BSM

Majorana: dominates over SV contribution

$0\nu\beta\beta$ may be a mixture of Majorana mass and other $\Delta L=2$ physics

\rightarrow mimics higher Majorana neutrino mass

Neutrinos Oscillation Surprises

... many untested assumptions: Majorana, 3 ν 's, mass mechanism

→ example: How NSI's can fool us in precision experiments:

Source	⊗	Oscillation	⊗	Detector
<ul style="list-style-type: none">- neutrino energy E- flux and spectrum- flavour composition- contamination- symmetric $\nu/\bar{\nu}$ operation		<ul style="list-style-type: none">- oscillation channels- realistic baselines- MSW matter profile- degeneracies- correlations		<ul style="list-style-type: none">- effective mass, material- threshold, resolution- particle ID (flavour, charge, event reconstruction, ...)- backgrounds- x-sections (at low E)

precision experiments might see new effects beyond oscillations!

NSI Operators

- **Good reasons for physics beyond the SM+ (with ν 's)**
 - expect effects beyond 3 flavours in many models
 - effective 4f interactions

$$\mathcal{L}_{NSI} \simeq \epsilon_{\alpha\beta} 2\sqrt{2}G_F (\bar{\nu}_{L\beta} \gamma^\rho \nu_{L\alpha}) (\bar{f}_L \gamma_\rho f_L)$$

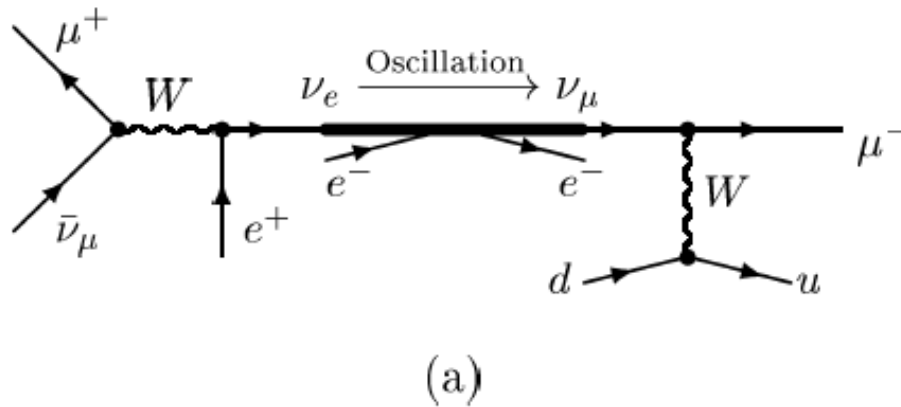
- **integrating out heavy physics (c.f. $G_F \leftrightarrow M_W$)**

$$|\epsilon| \simeq \frac{M_W^2}{M_{NSI}^2}$$

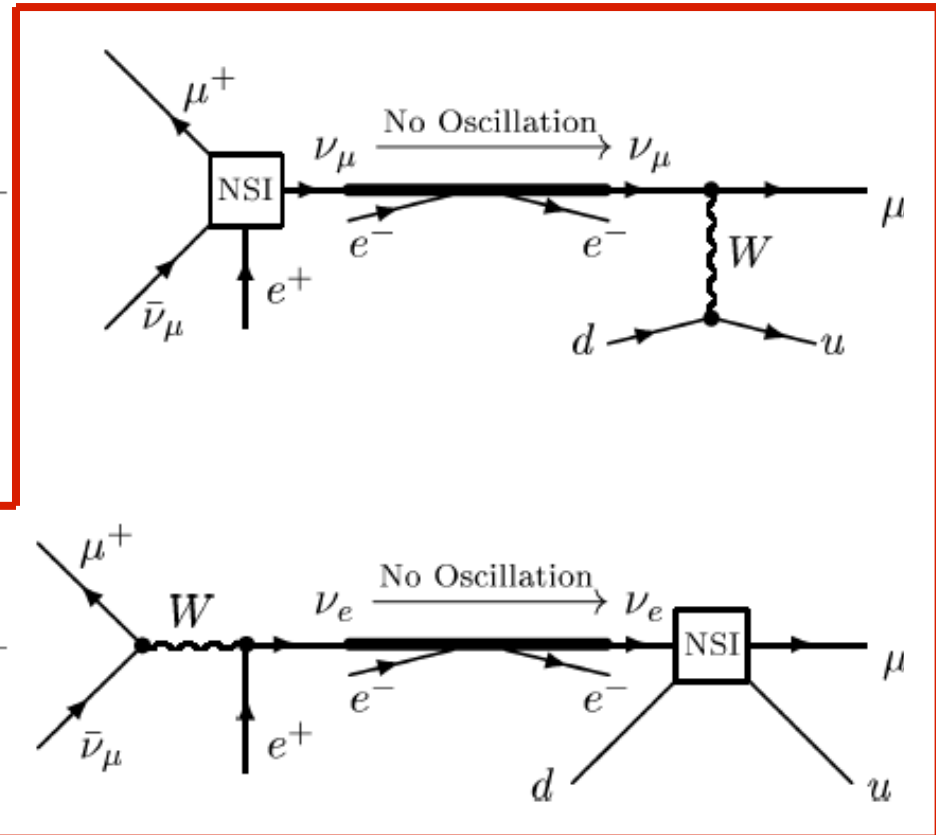
Grossman, Bergmann+Grossman, Ota+Sato, Honda et al., Friedland+Lunardini, Blennlow+Ohlsson+Skrotzki, Huber+Valle, Huber+Schwetz+Valle, Campanelli+Romanino, Bueno et al., Kopp+ML+Ota, ...

NSIs interfere with Oscillations

the “golden” oscillation channel

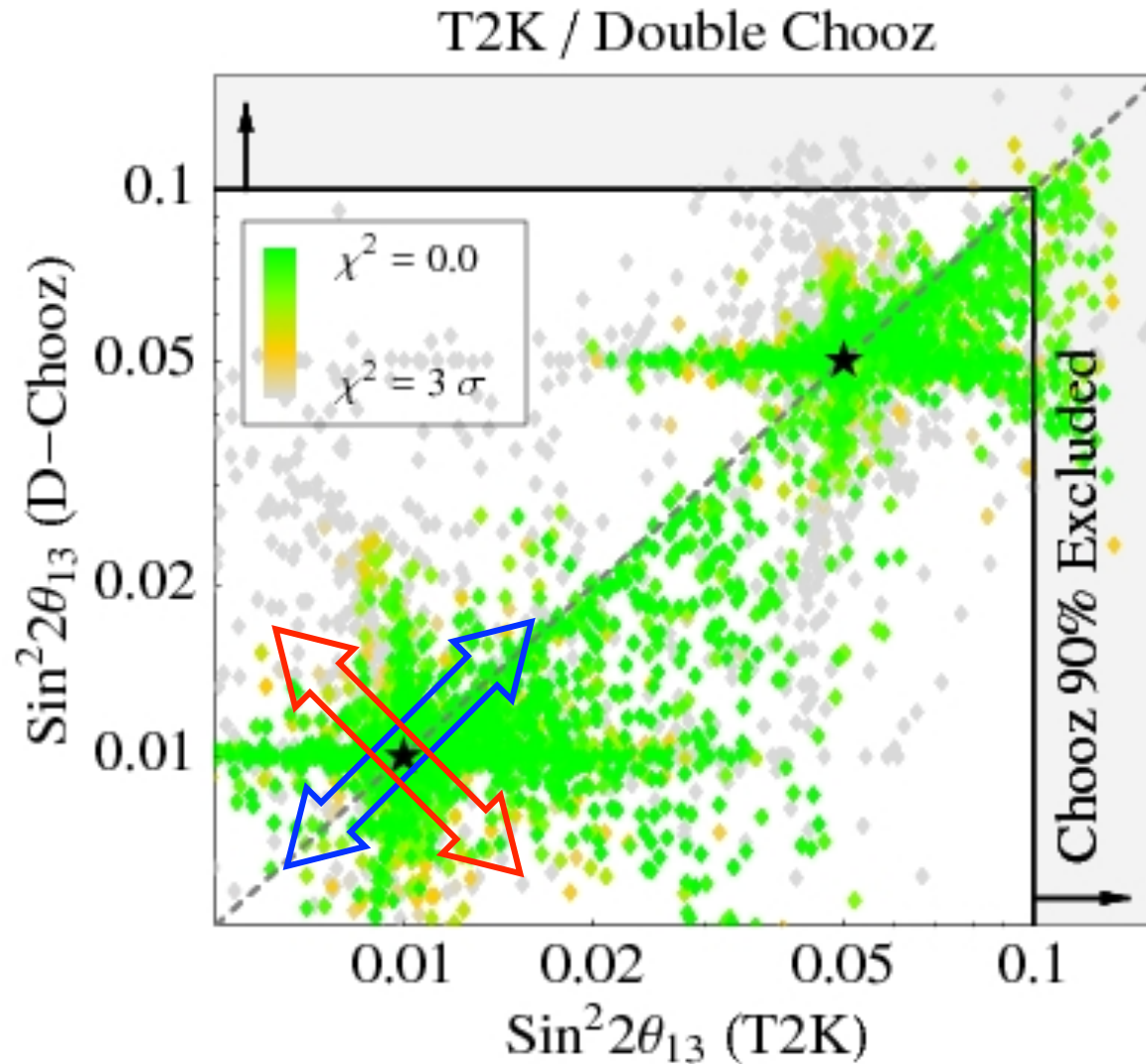


NSI contributions to the “golden” channel



note: interference in oscillations $\sim \epsilon$ \leftrightarrow FCNC effects $\sim \epsilon^2$

NSI: Offset and Mismatch in θ_{13}



Redundant measurements:

Double Chooz + T2K

***=assumed 'true' values of θ_{13}**

scatter-plot: ϵ values random

- below existing bounds

- random phases

NSIs can lead to:

- **offset**

- **mismatch**

➔ **redundancy**

➔ **interesting potential**

'natural magnitude' ...more natural for NuFact ➔ see talk by S. Pascoli

Evidence(s) for Sterile Neutrinos

LSND - MiniBooNE - MINOS – Gallex - ... → evidence for sterile ν 's?

→ see talk by E. Lisi

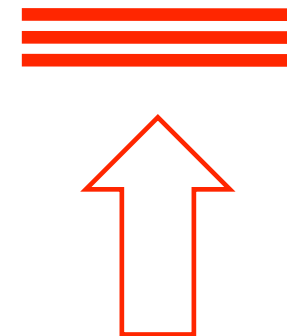
→ New and better data / experiments are needed to clarify the situation

→ maybe something exciting around the corner?

The standard picture:

3 heavy sterile neutrinos typ. $\geq 10^{13}$ GeV →

→ leptogenesis, role in GUTs, ...



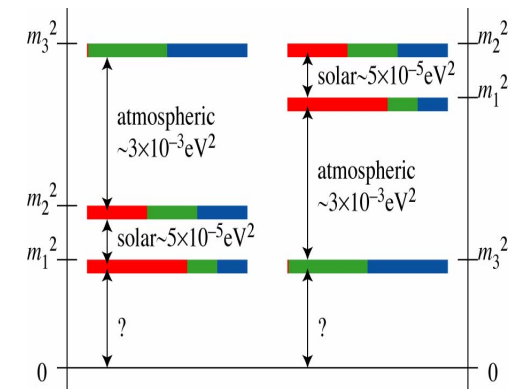
3 light active neutrinos →

→ this could easily be wrong

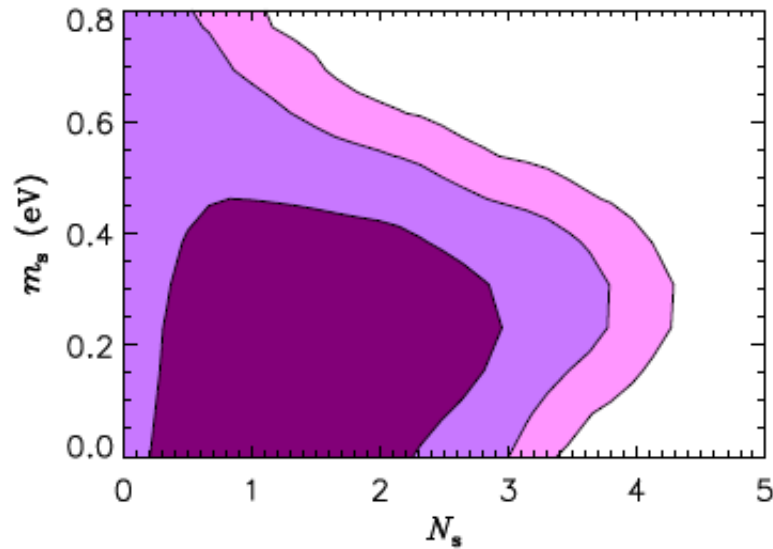
- more than 3 N_R states, ...

- M_R may have special eigenvalues, ...

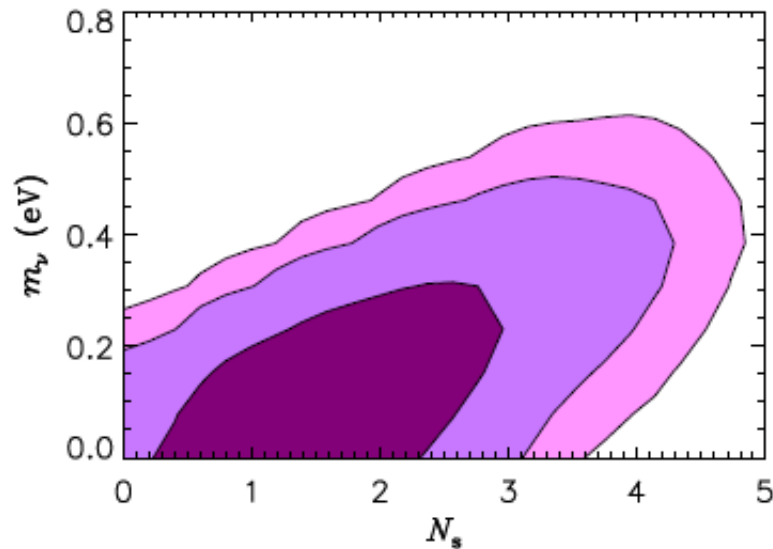
→ light sterile neutrinos ?!



Extra Sterile Neutrinos & CMB



**3 active massless neutrinos
+
 N_s massive neutrinos**



**3 active massive neutrinos
+
 N_s massless neutrinos**

J. Hamann et al

More Indications for light sterile Neutrinos

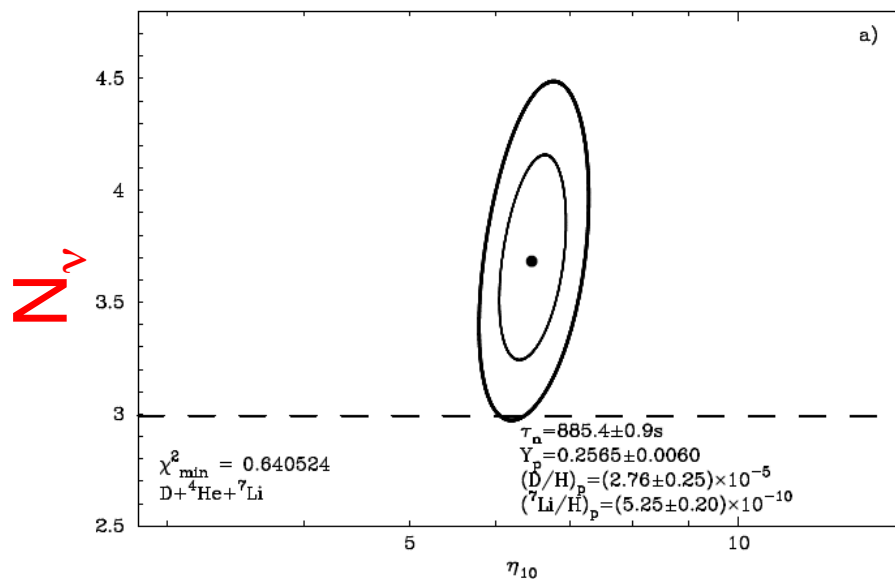
Astrophysics:

- Effects of keV-ish sterile neutrinos on pulsar kicks

Kusenko, Segre, Fuller, Mocioiu, Pascoli

- ...

Cosmology: BBN – ‘feels’ extra light particles:



$$N_{\nu} \simeq 3.7 \pm 1$$

E. Aver, K. Olive, E. Skillman (2010)

Y. Izotov, T. Thuan(2010)

Could Neutrinos be Dark Matter?

- **Active neutrinos = perfect Hot Dark Matter** → ruled out:
 - destroys small scale structures in cosmological evolution
 - required neutrino masses much too small → useful HDM component
 - **keV sterile neutrinos: Warm Dark Matter** → workes very well:
 - relativistic at decoupling
 - non-relativistic at radiation to matter dominance transition
 - OK for $M_\chi \simeq \text{few keV}$
 - reduced small scale structure → smoother profile, less dwarf satellites
 - scenario where one sterile neutrino is keV-ish, the others heavy

 - right handed neutrinos probably exist
 - only a mechanism to make one state light required
 - observational hints from astronomy
 - hints that a keV sterile particle may exist → right-handed neutrino?
- Biermann, Kusenko & Segre, Fuller et al., Biermann & Kusenko, Stasielak et al., Loewenstein et al., Dodelson, Widrow, Dolgov, ...

The ν MSM

Asaka, Blanchet, Shaposhnikov, 2005 Asaka, Shaposhnikov, 2005

Particle content:

- Gauge fields of $SU(3)_c \times SU(2)_W \times U(1)_Y$: g, W_{\pm}, Z, g
- Higgs doublet: $F=(1,2,1)$

• Matter

	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$	$U(1)_{em}$
$\begin{pmatrix} u \\ d \end{pmatrix}_L$	3	2	+1/3	$\begin{pmatrix} +2/3 \\ -1/3 \end{pmatrix}$
u_R	3	1	+4/3	+2/3
d_R	3	1	-2/3	-1/3
$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	1	2	-1	$\begin{pmatrix} 0 \\ -1 \end{pmatrix}$
e_R	1	1	-2	-1
N	1	1	0	0

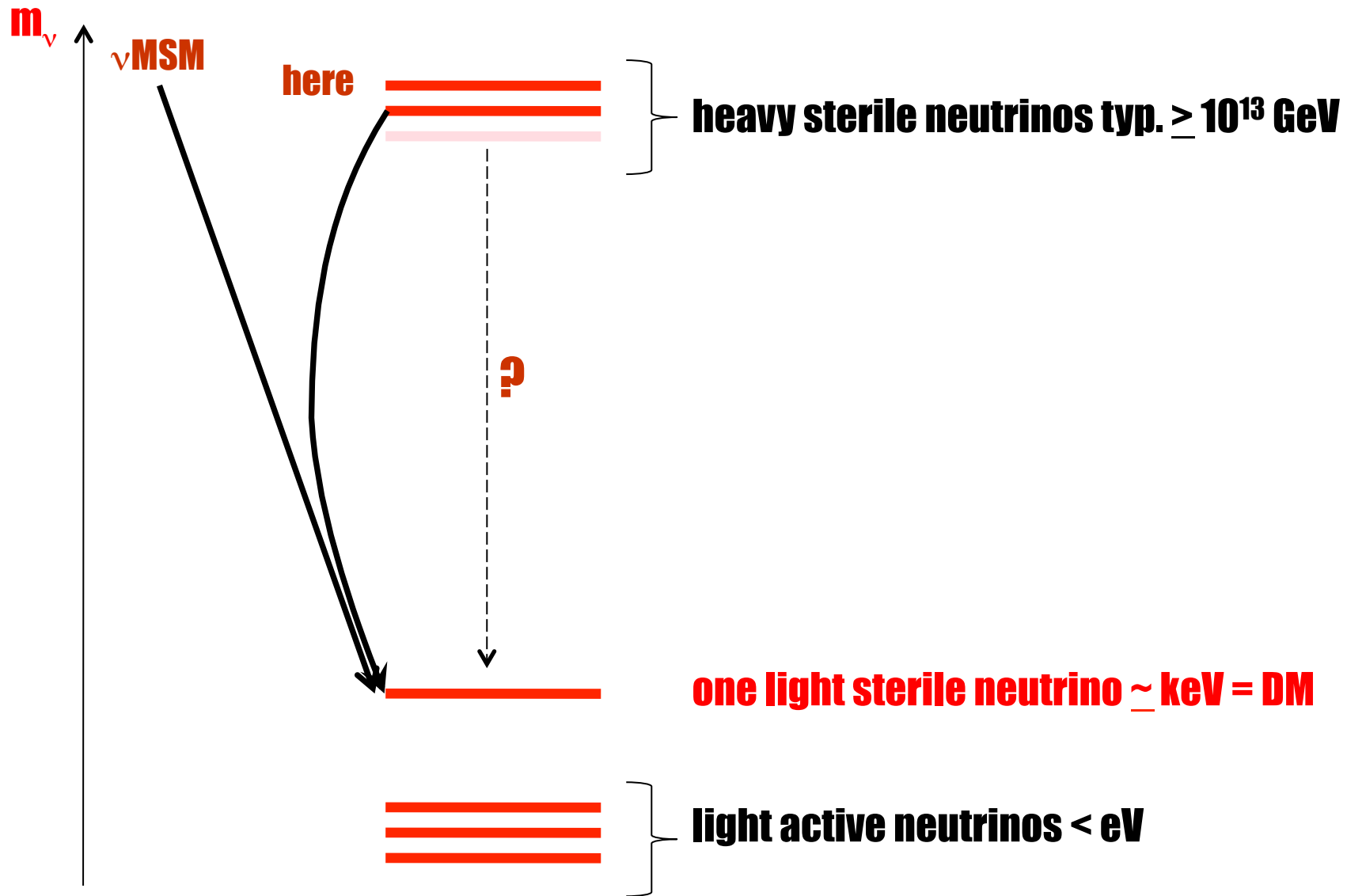
x3 generations

- lepton sector more symmetric to the quark sector
- Majorana masses for N
- choose for one sterile $n \sim \text{keV}$ mass → exceeds lifetime of Universe

Abundance in the ν MSM and in Alternatives

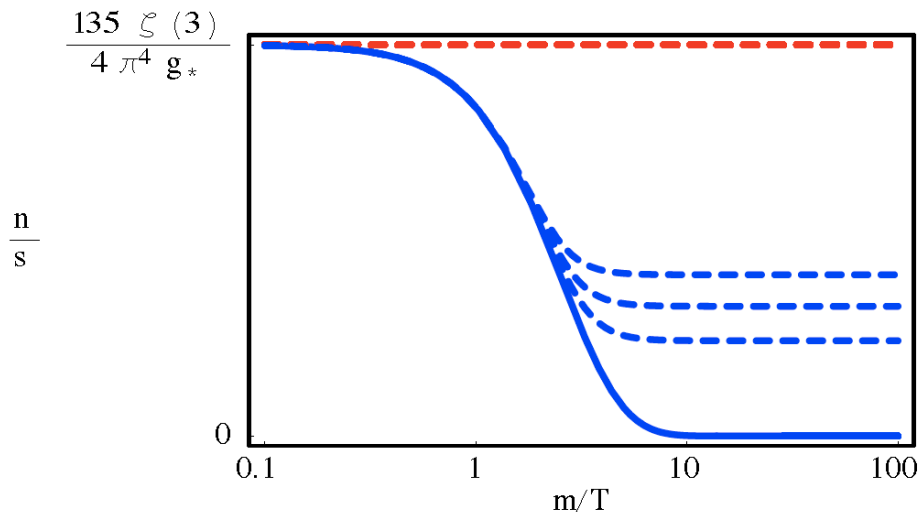
- **Virtue and problem of ν MSM:**
 - scenario with sterile ν and tiny mixing \rightarrow never enters thermal equilibrium
 - \rightarrow requires non-thermal production from other particles (avoid over-closure)
 - \rightarrow new physics before the beginning of the thermal evolution sets abundance
- **An alternative scenario: Bezrukov, Hettmannsperger, ML**
 - Three right-handed neutrinos N_1, N_2, N_3
 - Dirac and Majorana mass terms
 - **N Charged under some (BSM) gauge group \rightarrow scale M**
 - Specific LR example: $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$
- **everything follows nicely from sterile neutrinos**

The Scenario



Obtaining the right Abundance

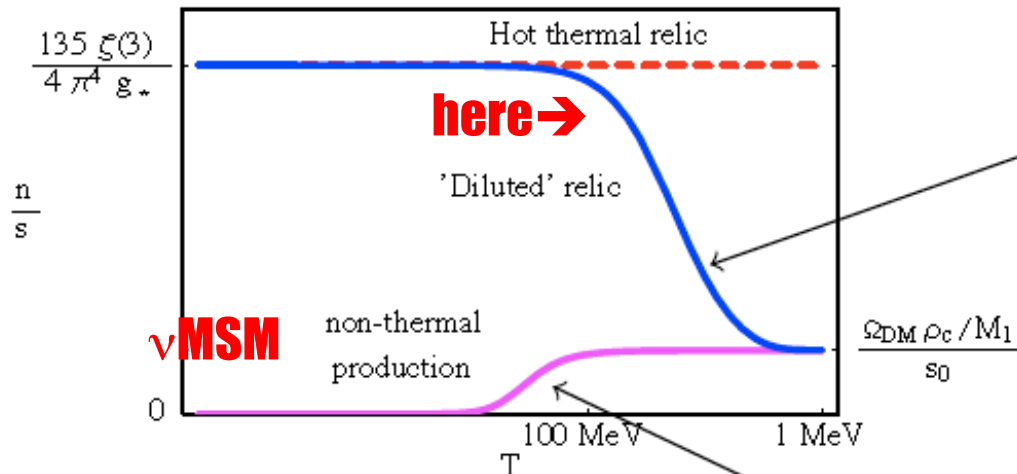
Usual thermal case:



HDM: $\frac{\Omega}{\Omega_{DM}} \simeq \left(\frac{10}{g_{*f}}\right) \left(\frac{M}{10\text{eV}}\right)$
Decoupled relativistic

CDM:
(M >> MeV) $\Omega \sim \Omega_{DM}$
Decoupled nonrelativistic

keV sterile neutrinos:



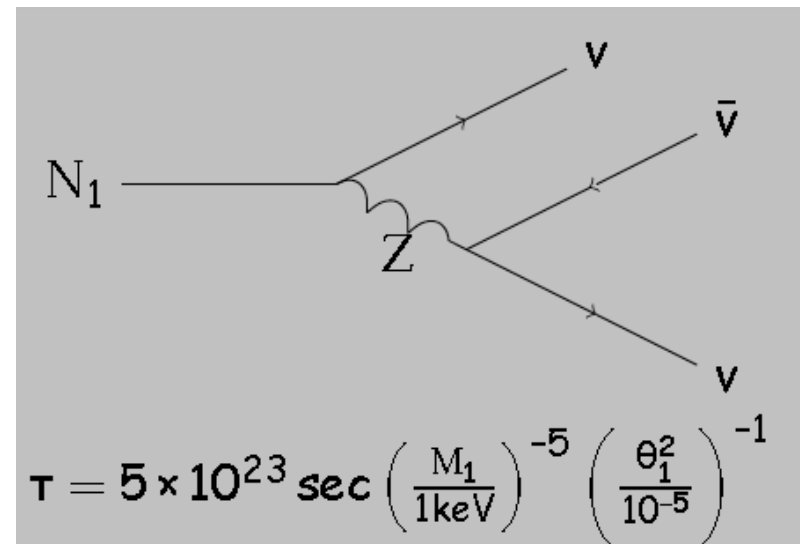
Diluted after decoupling
(entropy generated by other
particle decay)

$\Omega \sim \Omega_{DM}$

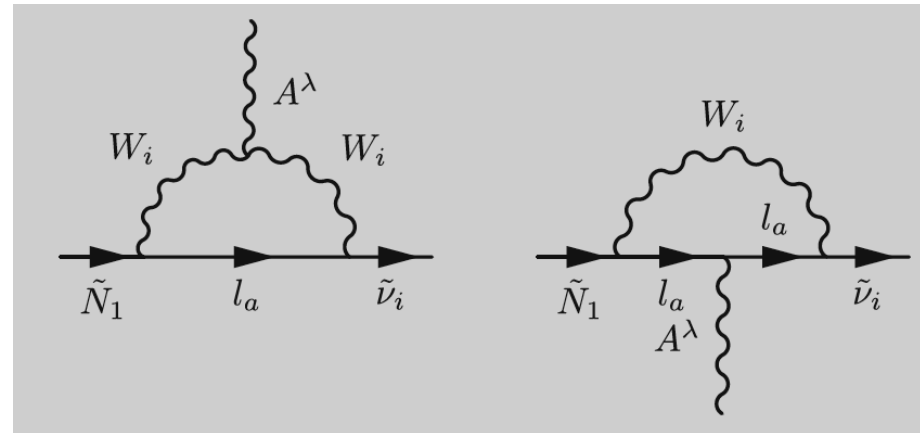
Never entered thermal equilibrium

Observing keV-ish Neutrino DM

- **LHC**
 - sterile neutrino DM is not observable
 - WIMP-like particles still possible – but not DM
- **direct searches**
 - sterile neutrino DM is not observable
- **astrophysics/cosmology** → at some level: keV X-rays
 - sterile neutrino DM is decaying into active neutrinos
 - decay $N_1 \rightarrow \bar{\nu}\nu\nu$, $N_1 \rightarrow \bar{\nu}\bar{\nu}\nu$
 - not very constraining since $\tau \gg \tau_{\text{Universe}}$



- radiative decays $N_1 \rightarrow \nu\gamma$



- so far: observations limit active-sterile mixing angle

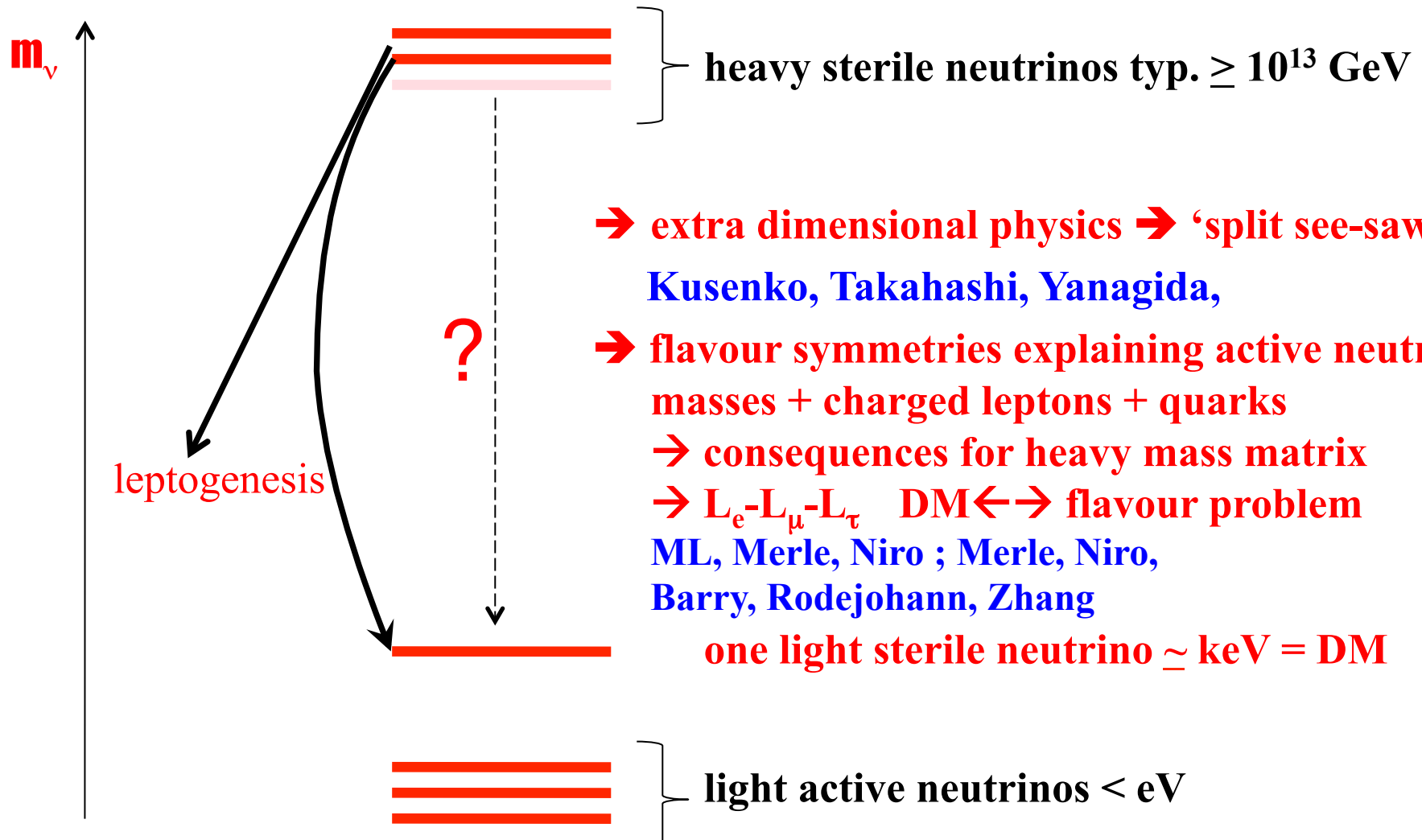
$$\Gamma_{N_1 \rightarrow \nu\gamma} \simeq 5.5 \times 10^{-22} \theta_1^2 \left(\frac{M_1}{1 \text{ keV}} \right)^5 \text{ s}^{-1}$$

$$\theta_1^2 \lesssim 1.8 \times 10^{-5} \left(\frac{1 \text{ keV}}{M_1} \right)^5$$

- mixing tiny, but naturally expected to be tiny: $O(\text{scale ratio})$

Generating keV-ish Sterile Neutrinos

Possible scenario: See-saw + a reason why 1 sterile ν is light



- extra dimensional physics → ‘split see-saw’
Kusenko, Takahashi, Yanagida,
- flavour symmetries explaining active neutrino masses + charged leptons + quarks
→ consequences for heavy mass matrix
→ $L_e - L_\mu - L_\tau$ DM \leftrightarrow flavour problem
ML, Merle, Niro ; Merle, Niro,
Barry, Rodejohann, Zhang

Conclusions

- There are very good reasons to go beyond the SM
- Neutrino masses are still the only solid direct evidence for BSM
- Many non-standard neutrino options & connections exist
 - the three neutrino picture may be incomplete
 - connections to EW-SB, DM, LFV, ...

→ neutrinos may easily surprise us again!

