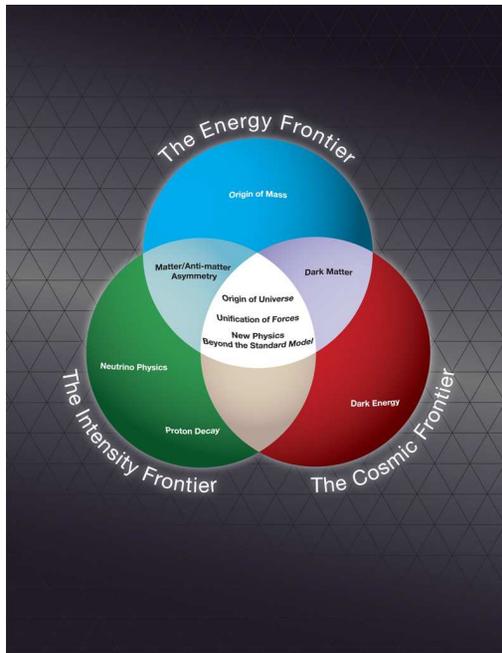


# STERILE NEUTRINOS: Baryogenesis, dark matter, primordial magnetic fields...



Oleg RUCHAYSKIY



24th Rencontres de Blois  
May 31, 2012



# Standard Model today

	I	II	III		
mass →	2.4 MeV	1.27 GeV	171.2 GeV	0	0
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
name →	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	
Quarks	4.8 MeV $-\frac{1}{3}$ <b>d</b> down	104 MeV $-\frac{1}{3}$ <b>s</b> strange	4.2 GeV $-\frac{1}{3}$ <b>b</b> bottom	0 0 <b><math>\gamma</math></b> photon	
	0 eV 0 <b><math>\nu_e</math></b> electron neutrino	0 eV 0 <b><math>\nu_\mu</math></b> muon neutrino	0 eV 0 <b><math>\nu_\tau</math></b> tau neutrino	Bosons (Forces) spin 1	91.2 GeV 0 <b>Z<sup>0</sup></b> weak force
	0.511 MeV -1 <b>e</b> electron	105.7 MeV -1 <b><math>\mu</math></b> muon	1.777 GeV -1 <b><math>\tau</math></b> tau		>114 GeV 0 0 <b>H</b> Higgs boson
Leptons				80.4 GeV $\pm 1$ <b>W<sup>±</sup></b> weak force	spin 0

## Future?

- SM Higgs gives a good fit to data.  
Reduced  $gg \rightarrow h$  and enhanced  $h \rightarrow \gamma\gamma$  improves the fit.  
Too good: is this just over-fitting fluctuations?
- SUSY: at the weak scale, or one loop above, or much above.
- $m_h \approx 125$  GeV corresponds to  $\lambda = 0$  at the Planck scale? Almost, but NO.  
 $\lambda$  gets slightly negative and the SM vacuum is meta-stable.

Implications for European Strategy for Particle Physics:  
The Higgs could be the last particle. Carpe diem.

From the talk of A. Strumia at CERN workshops "*Implication of the latest LHC results for new physics*"

## Beyond the Standard Model

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**BUT!** already now we know a number of observational **beyond the Standard Model** phenomena:

- **Neutrino oscillations**: transition between neutrinos of different flavours ( $\nu_e, \nu_\mu, \nu_\tau$ ) means violation of lepton flavour symmetries (but not total lepton number!)
- existence of **dark matter** (why observed gravity of galaxies and clusters is so strong?)
- the **absence of anti-matter** in the Universe
- **(Probably)** inflation (homogeneity of the observed Universe seem to require correlated initial conditions for causally non-connected regions)
- **(Maybe)** dark energy (If it will be shown that accelerated expansion of the Universe is caused not by a small cosmological constant, but by some other unknown substance – what is this substance?)

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What should we do with  
**beyond-the-Standard-Model**  
**problems** if the “nightmare  
scenario” becomes true?

# Neutrinos: left-only particles

	I	II	III		
mass →	2.4 MeV	1.27 GeV	171.2 GeV	0	0
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
name →	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	
Quarks	4.8 MeV $-\frac{1}{3}$ <b>d</b> down	104 MeV $-\frac{1}{3}$ <b>s</b> strange	4.2 GeV $-\frac{1}{3}$ <b>b</b> bottom	0 0 <b><math>\gamma</math></b> photon	
	0 eV 0 <b><math>\nu_e</math></b> electron neutrino	0 eV 0 <b><math>\nu_\mu</math></b> muon neutrino	0 eV 0 <b><math>\nu_\tau</math></b> tau neutrino	Bosons (Forces) spin 1	91.2 GeV 0 <b>Z<sup>0</sup></b> weak force
	0.511 MeV -1 <b>e</b> electron	105.7 MeV -1 <b><math>\mu</math></b> muon	1.777 GeV -1 <b><math>\tau</math></b> tau		>114 GeV 0 0 <b>H</b> Higgs boson
Leptons				80.4 GeV $\pm 1$ <b>W<sup>±</sup></b> weak force	spin 0

# Right-handed neutrinos: sterile particles

	I	II	III		
mass →	2.4 MeV	1.27 GeV	171.2 GeV	0	
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	
name →	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	
	Left Right	Left Right	Left Right		
	4.8 MeV	104 MeV	4.2 GeV	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
<b>Quarks</b>	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
	Left Right	Left Right	Left Right		
	<0.0001 eV ~10 keV	~0.01 eV ~GeV	~0.04 eV ~GeV	91.2 GeV	>114 GeV
	0	0	0	0	0
	<b><math>\nu_e</math></b> <b><math>N_1</math></b>	<b><math>\nu_\mu</math></b> <b><math>N_2</math></b>	<b><math>\nu_\tau</math></b> <b><math>N_3</math></b>	<b>Z</b> <sup>0</sup>	<b>H</b>
	electron neutrino sterile neutrino	muon neutrino sterile neutrino	tau neutrino sterile neutrino	weak force	Higgs boson
	Left Right	Left Right	Left Right		
	0.511 MeV	105.7 MeV	1.777 GeV	80.4 GeV	
	-1	-1	-1	$\pm 1$	
<b>Leptons</b>	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>W</b> <sup>±</sup>	spin 0
	Left Right	Left Right	Left Right	weak force	

Bosons (Forces) spin 1

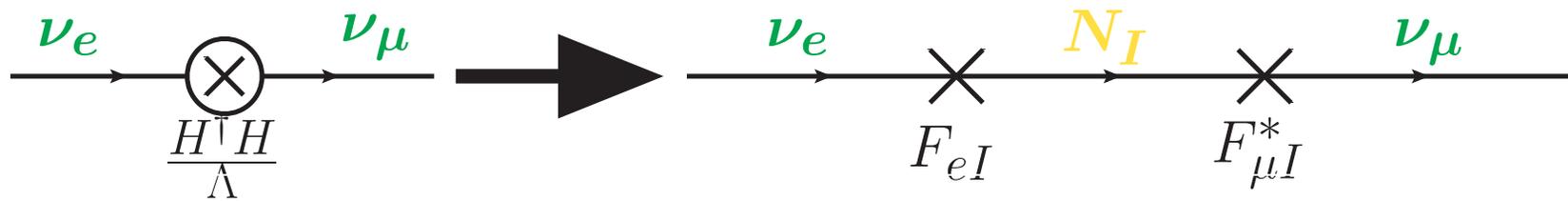
spin 0

# Neutrino oscillations and sterile neutrinos

mass →	2.4 MeV	1.27 GeV	171.2 GeV	0	0	0	0
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0	0	0
name →	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>γ</b> photon	<b>Z</b> <sup>0</sup> weak force	<b>H</b> Higgs boson
<b>Quarks</b>	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom				
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$				
	Left	Left	Left				
	Right	Right	Right				
	$<0.0001$ eV	$\sim 0.01$ eV	$\sim 0.04$ eV				
	$\sim 10$ keV	$\sim \text{GeV}$	$\sim \text{GeV}$				
	<b><math>\nu_e</math></b>	<b><math>\nu_\mu</math></b>	<b><math>\nu_\tau</math></b>				
	electron neutrino	muon neutrino	tau neutrino				
	<b><math>N_1</math></b>	<b><math>N_2</math></b>	<b><math>N_3</math></b>				
	sterile neutrino	sterile neutrino	sterile neutrino				
<b>Leptons</b>							
	0	0	0				
	Left	Left	Left				
	Right	Right	Right				
	0.511 MeV	105.7 MeV	1.777 GeV				
	<b>e</b> electron	<b>μ</b> muon	<b>τ</b> tau				
	-1	-1	-1				
	Left	Left	Left				
	Right	Right	Right				
	80.4 GeV						
	<b>W</b> <sup>±</sup> weak force						
	±1						
	Left						
	Right						

Bosons (Forces) spin 1

>114 GeV  
spin 0



## See-saw Lagrangian

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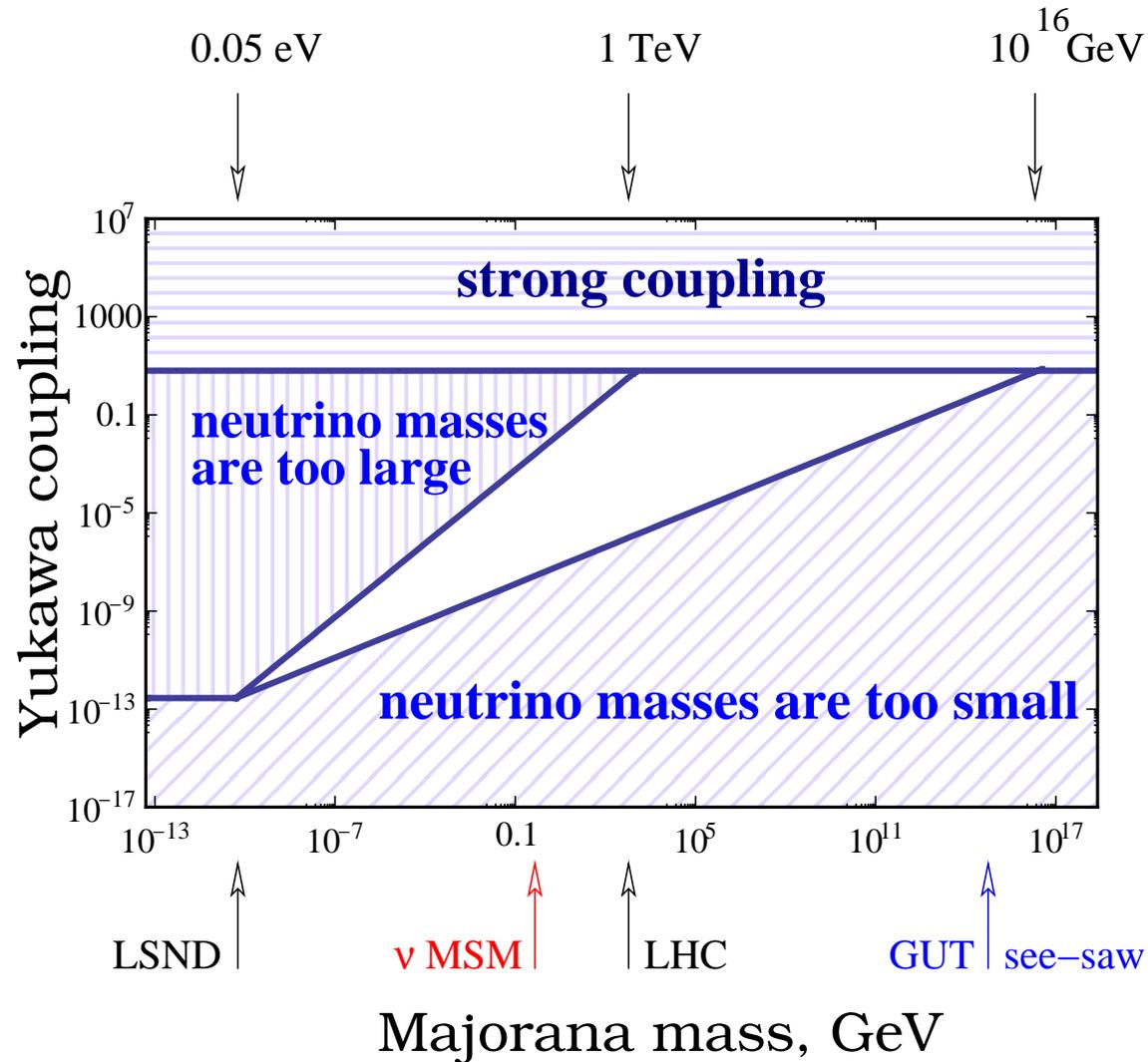
- Right-chiral neutrinos  $N_I$  **carry no charge** under the Standard Model interactions  $\Rightarrow$  **sterile neutrinos**

$$\mathcal{L}_{\text{see-saw}} = i\bar{N}\not{\partial}N + M_D\bar{\nu}N + \frac{M}{2}\bar{N}^cN$$

- **Dirac** mass term:  $(M_D)_{\alpha I} = \langle \text{Higgs} \rangle F_{\alpha I}$
- Two mass differences mean that there are **at least two** sterile neutrinos
- Neutrino masses are given by **see-saw formula**:

$$\text{Neutrino mass matrix} = -M_{\text{Dirac}} \frac{1}{M_{\text{Majorana}}} M_{\text{Dirac}}^T$$

# Scale of sterile neutrino masses



See the next talk by Pilar Hernandez

# Neutrino Minimal Standard Model

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- Alternative choice: make the masses of sterile neutrinos of the same order as those of quarks and leptons (keV–GeV) — **Neutrino Minimal Standard Model** ( $\nu$ MSM for short)
- The  $\nu$ MSM solves several beyond the Standard Model problems and **provides a complete cosmic history from inflation till today**

Asaka & Shaposhnikov (2005) and many subsequent works

- ✓ ... explains neutrino oscillations
- ✓ ... generates matter-antimatter asymmetry of the Universe
- ✓ ... generates cosmic magnetic fields

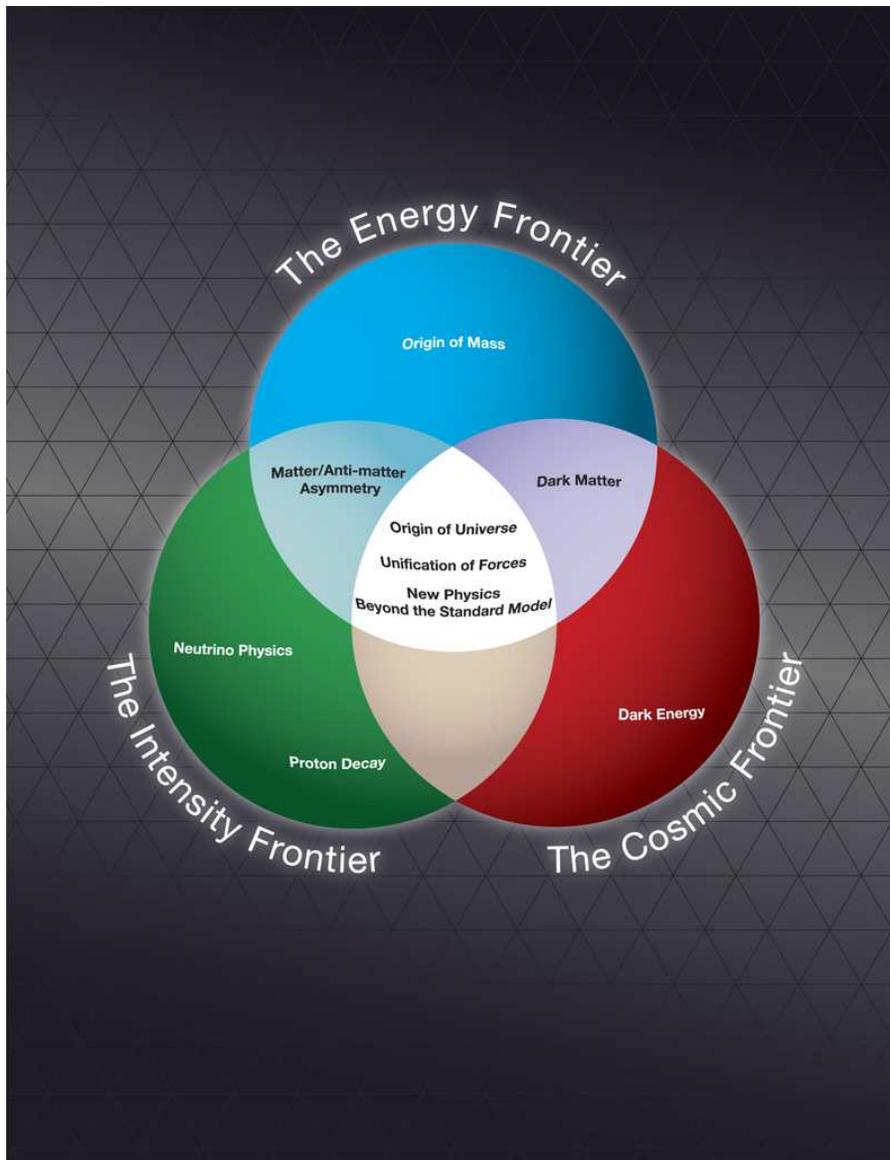
## Two sterile neutrinos with MeV–GeV masses

- ✓ ... provides a dark matter particle (cold, **warm** or **mixed**)

## Third sterile neutrino with keV mass

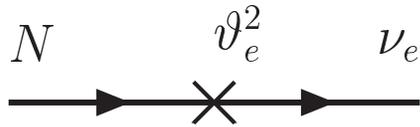
Review: Boyarsky, O.R., Shaposhnikov *Ann. Rev. Nucl. Part. Sci.* (2009), [0901.0011]

# Sterile neutrino can be searched:

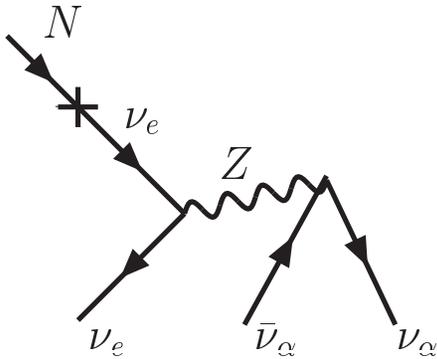


- At accelerators (“Intensity frontier experiments”)
- In the spectra of galaxies and galaxy clusters (“Cosmic frontier”)
- To make predictions for these searches the  $\nu$ MSM needs detailed early Universe computations, providing
  - correct baryon asymmetry
  - correct DM abundance
  - other early Universe signatures

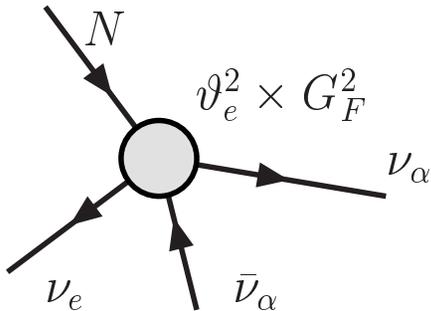
# Properties of sterile neutrino



Quadratic mixing  $N \leftrightarrow \nu$  of sterile neutrino to  $\nu$ ,  $\theta^2 = \frac{|M_{\text{Dirac}}|^2}{M_{\text{Majorana}}^2}$



Decay of sterile neutrino  $N \rightarrow \nu_e \nu_\alpha \bar{\nu}_\alpha$  through neutral current interactions



Fermi-like interaction with the “effective” Fermi constant  $\vartheta_e \times G_F$

**Sterile neutrinos behave as superweakly interacting heavy neutrinos with Fermi constant  $G_F \rightarrow \theta \cdot G_F, \theta \lll 1$**

## Baryo/leptogenesis in the $\nu$ MSM

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All three Sakharov conditions are satisfied if neutrinos are super-weakly interacting and light ( $M < M_W$ ):

Kuzmin,  
Rubakov,  
Shaposhnikov  
(1985)

**B-number violation:** sphalerons

**CP (and C) non-conservation:** phase of the CKM matrix **plus additional CP phases in the Dirac mass matrix of sterile neutrinos**

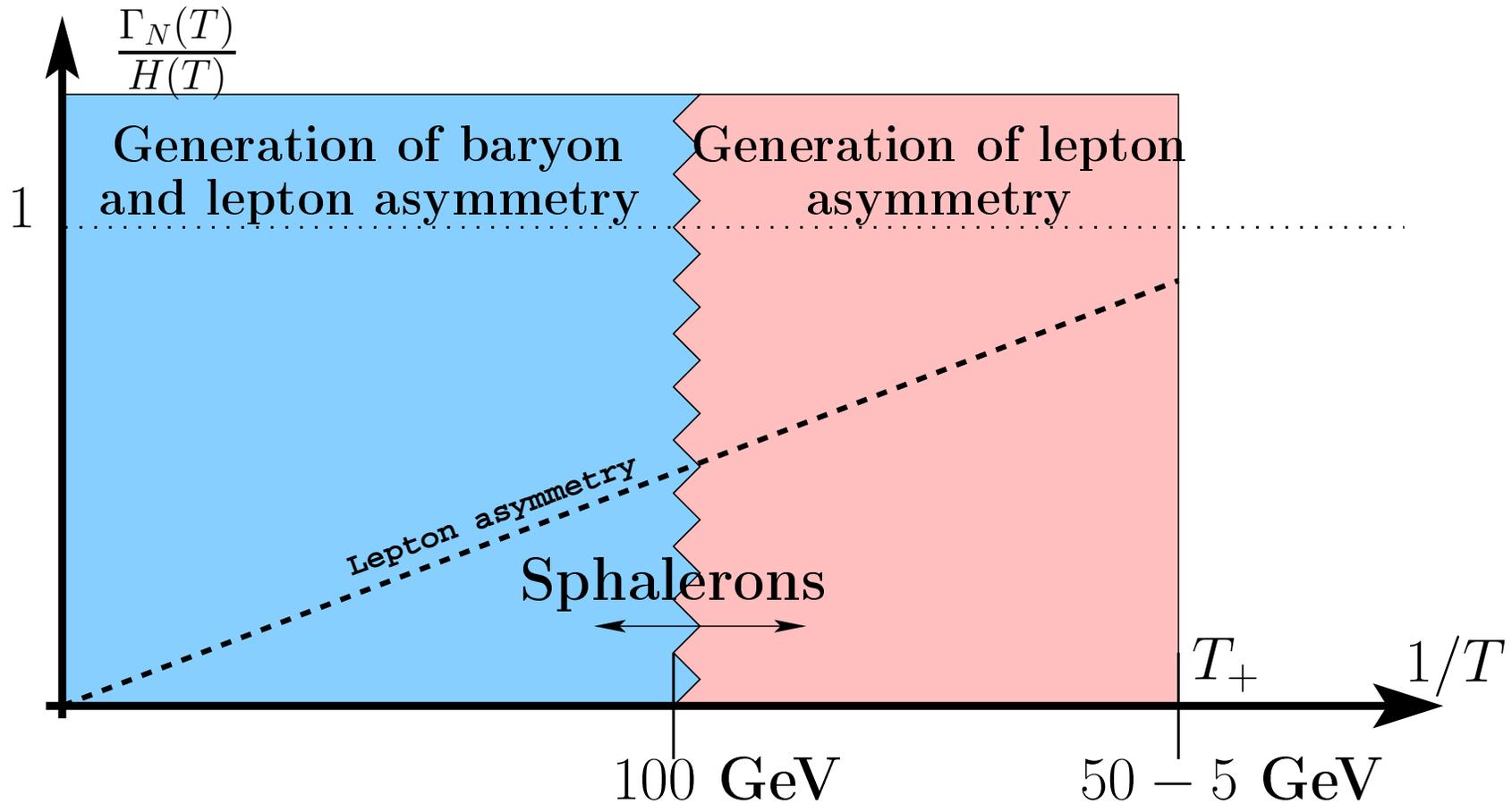
Farrar &  
Shaposhnikov  
(1994)

**Out-of-equilibrium processes:** no phase transition in the  $\nu$ MSM for  $m_H > 72$  GeV! **but Yukawa couplings of sterile neutrinos are small enough to keep them out of thermal equilibrium at  $T \sim 100$  GeV**

Kajantie et al.  
(1996)

Asaka,  
Shaposhnikov  
(2005)

## Baryo/leptogenesis in the $\nu$ MSM



In the  $\nu$ MSM lepton asymmetry is **much higher** than the baryon asymmetry

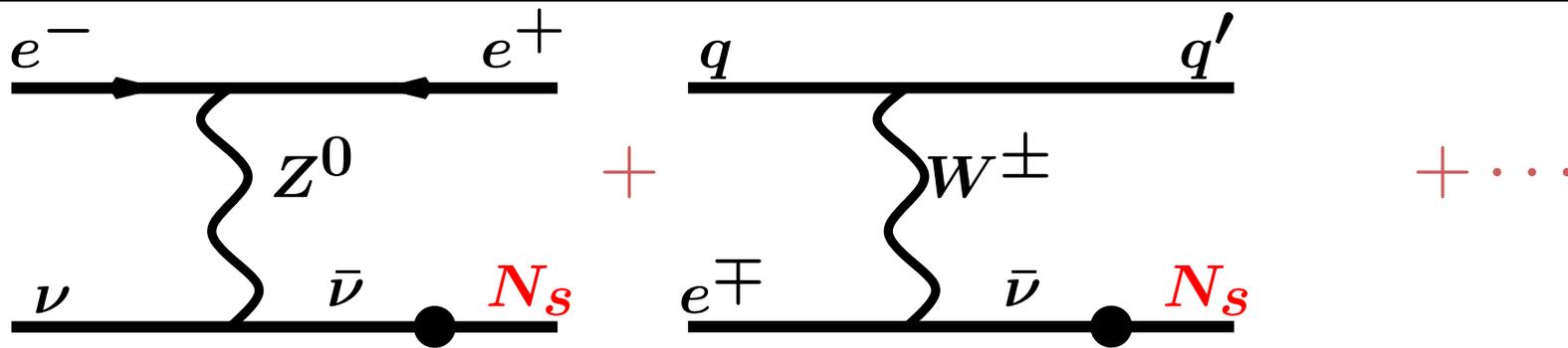
# Sterile neutrino dark matter

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- Two sterile neutrinos of the  $\nu$ MSM with the masses in MeV-GeV range are responsible for **neutrino oscillations**
- The same particle **generate baryon and lepton asymmetries of the Universe**
- The third sterile neutrino is **not required** to explain neutrino oscillations  $\Rightarrow$  can couple to the usual Standard Model neutrino sector arbitrarily weakly. This is a **dark matter** particle in the  $\nu$ MSM
- Two sterile neutrinos determine the properties of the DM sterile neutrino (abundance, primordial velocities, etc.)

Asaka,  
Shaposhnikov  
et al. 2005-...

# Production of sterile neutrino DM



Dolgov  
Hansen (2000)

Abazajian  
Fuller Tucker  
(2001)

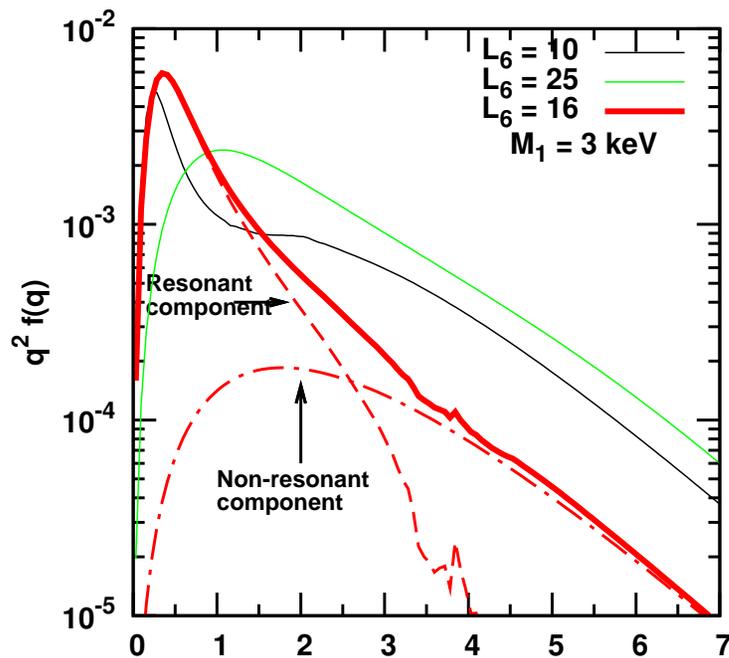
Asaka,  
Shaposhnikov  
Laine

Boyarsky, O.R.  
et al.  
(2006-2009)

Shi & Fuller  
(1998)

Laine &  
Shaposhnikov

In the presence of **lepton asymmetry** created by two other sterile neutrinos oscillations become **resonant**



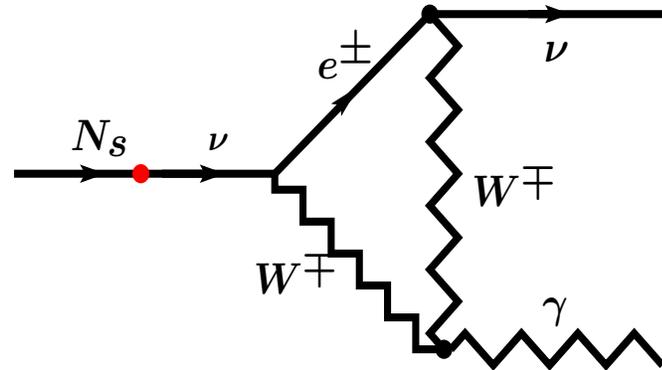
Sterile neutrino dark matter have

- Colder (**resonant**) component with  $\langle p \rangle \ll T_\nu$
- Warmer (**non-resonant**) component with  $\langle p \rangle \sim 3T_\nu$

# Signatures of sterile neutrino DM

- Sterile neutrino DM is produced **relativistic** in the early Universe — **warm** dark matter
- Erases primordial density fluctuations at scales below the **free streaming scale**

- Can **decay** with cosmological lifetime

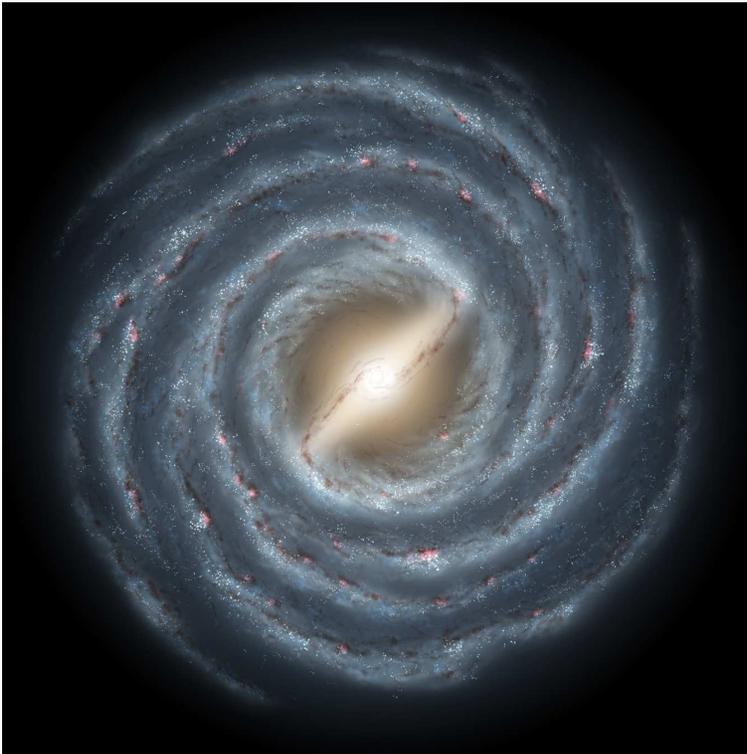


- Very characteristic signal: narrow line with  $E_\gamma = \frac{1}{2}m_{\text{DM}}c^2$
- The width of the decay line due to **Doppler broadening**  $\frac{\Delta E}{E_\gamma} \sim \frac{v_{\text{vir}}}{c} \sim 10^{-4} \div 10^{-3}$

## Search for dark matter particles

---

- DM may be decaying with a cosmologically long life-time (age of the Universe or even longer). Can we detect such decay?
- **Yes!** if you multiply a small number (probability of decay) with a large number (typical amount of DM particles in a galaxy  $\sim 10^{70}-10^{100}$ )



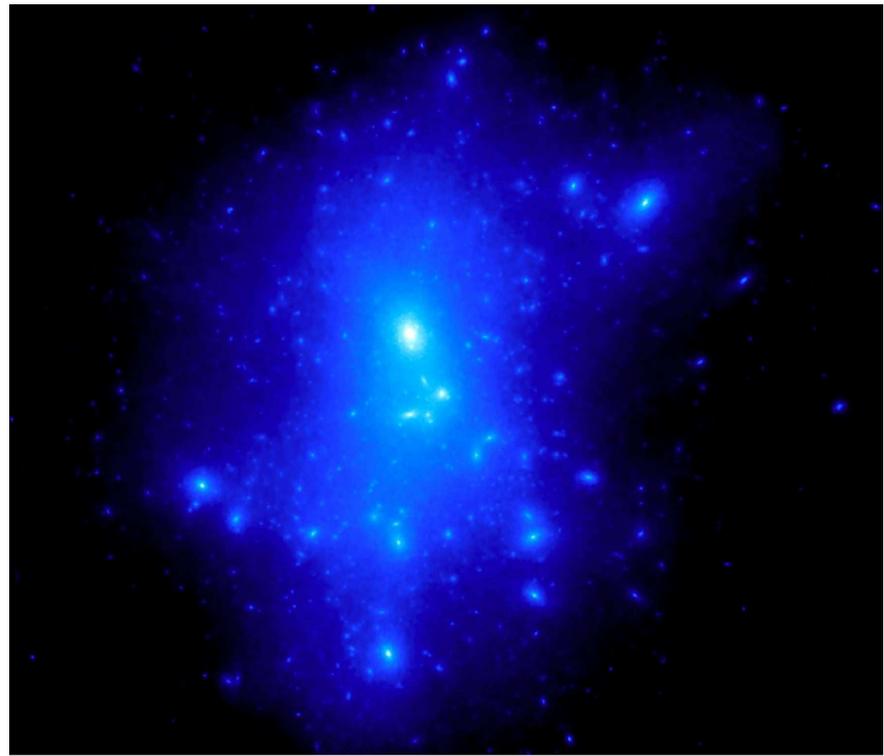
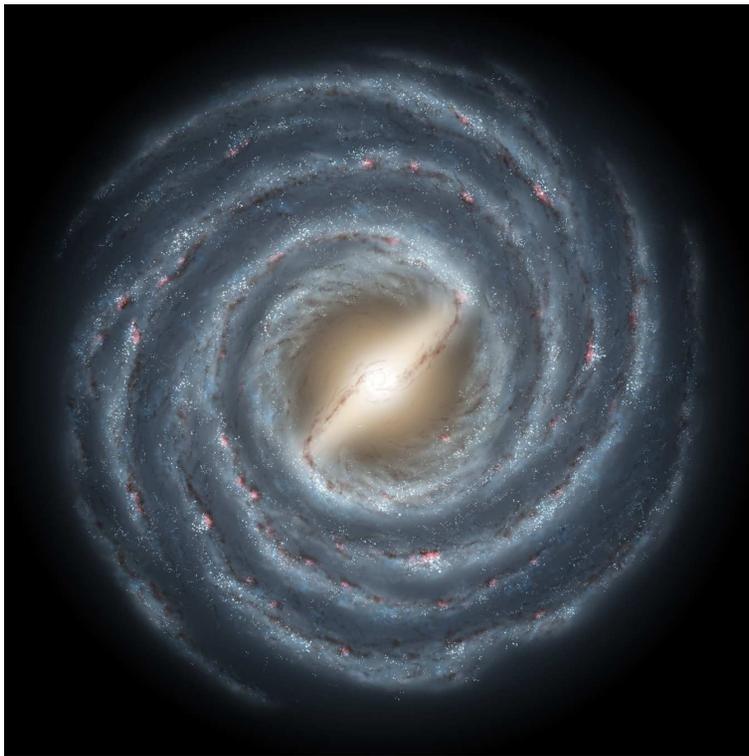
$$\text{Signal} \propto \int_{\text{line of sight}} \rho_{\text{DM}}(r) dl$$

Expected signal from the galaxy at a particular energy

## Search for dark matter particles

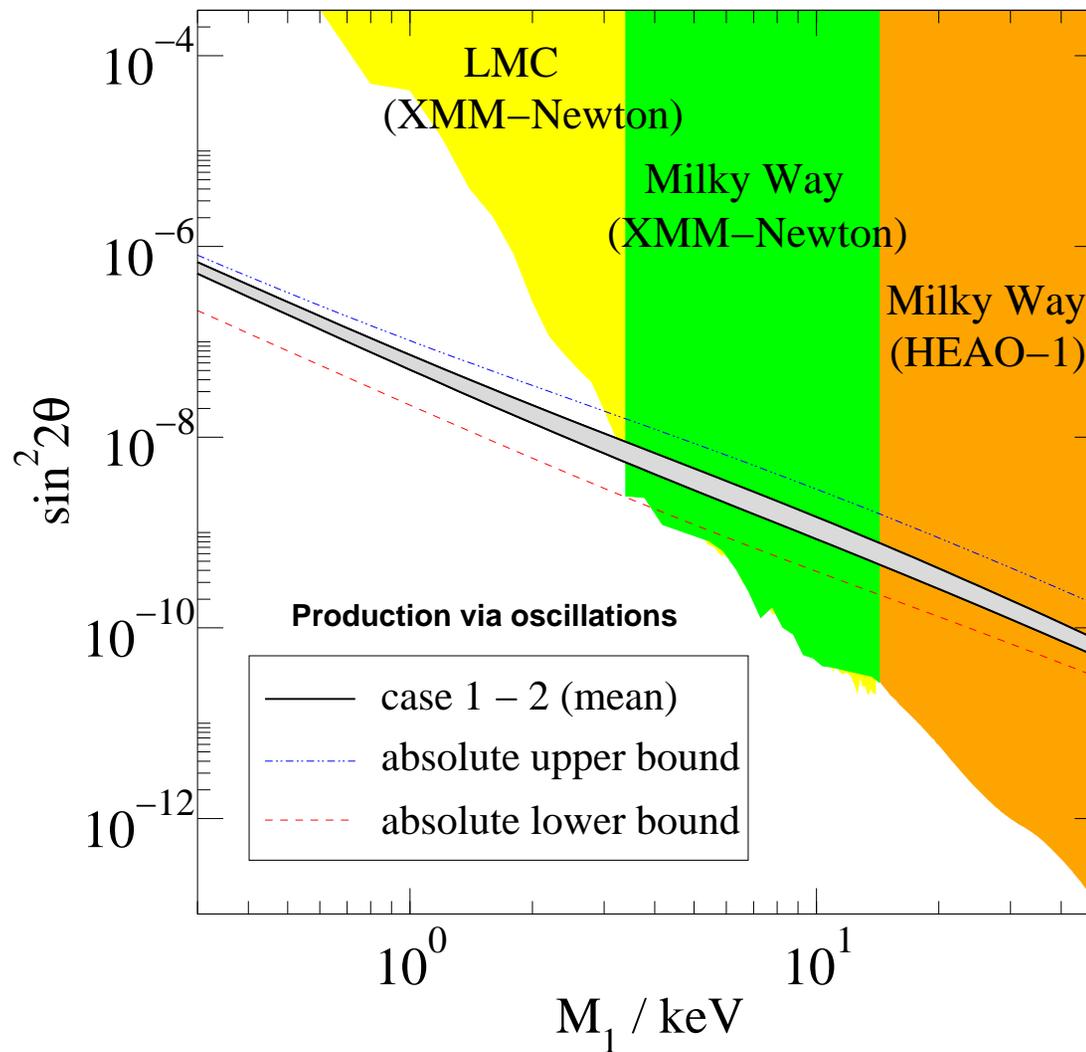
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Expected signal from a galaxy at a particular energy (simulation from B. Moore)

# Parameters of sterile neutrino DM



Production: Asaka, Laine, Shaposhnikov (2006)

MW (HEAO-1)  
Boyarsky, O.R.  
et al. 2005

Coma and  
Virgo clusters  
Boyarsky, O.R.  
et al.

Bullet cluster  
Boyarsky, O.R.  
et al. 2006

LMC+MW(XM  
Boyarsky, O.R.  
et al. 2006

MW Riemer-  
Sørensen et  
al.; Abazajian  
et al.

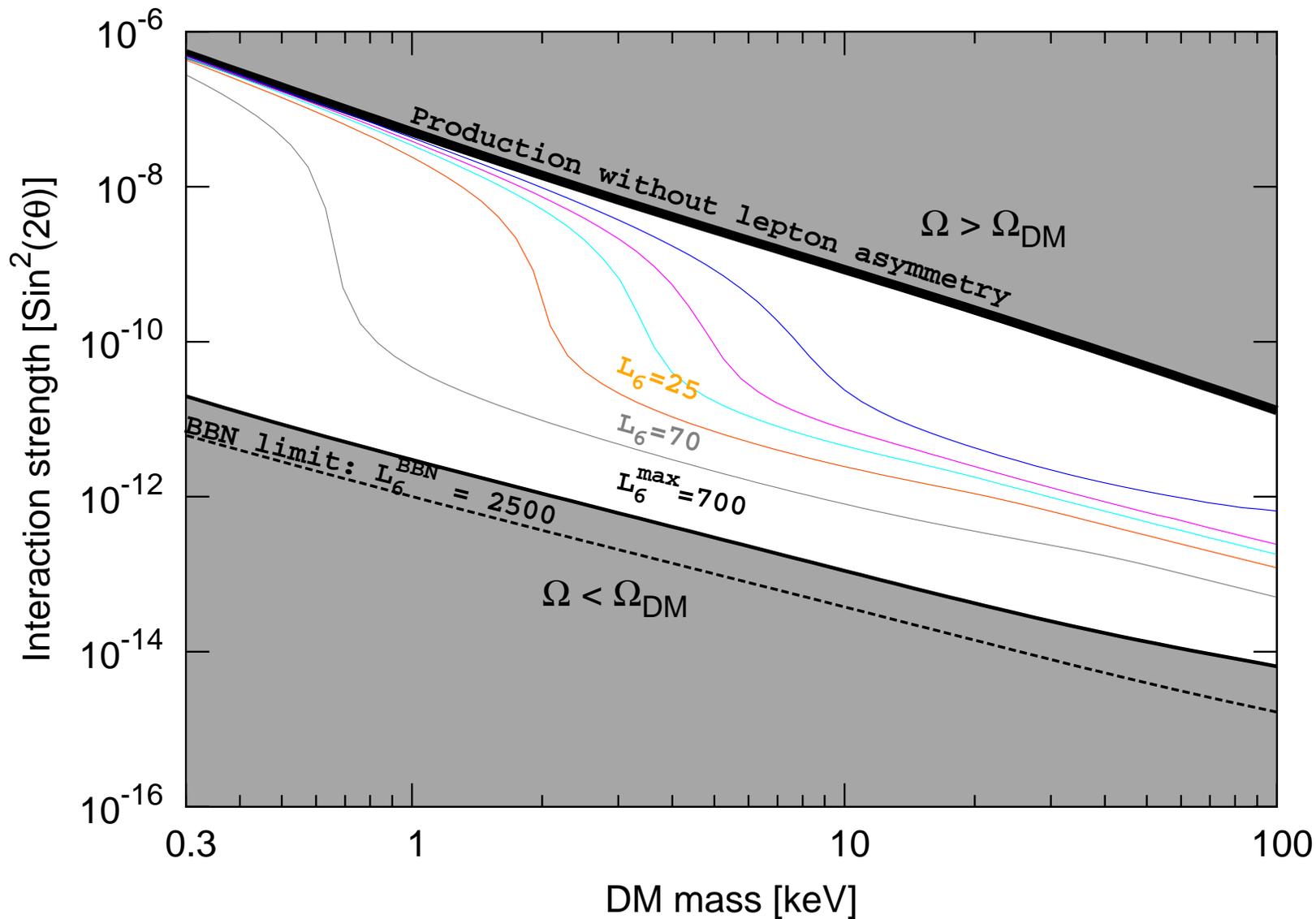
MW (XMM)  
Boyarsky, O.R.  
et al. 2007

M31 Watson  
et al. 2006;  
Boyarsky et al  
2007

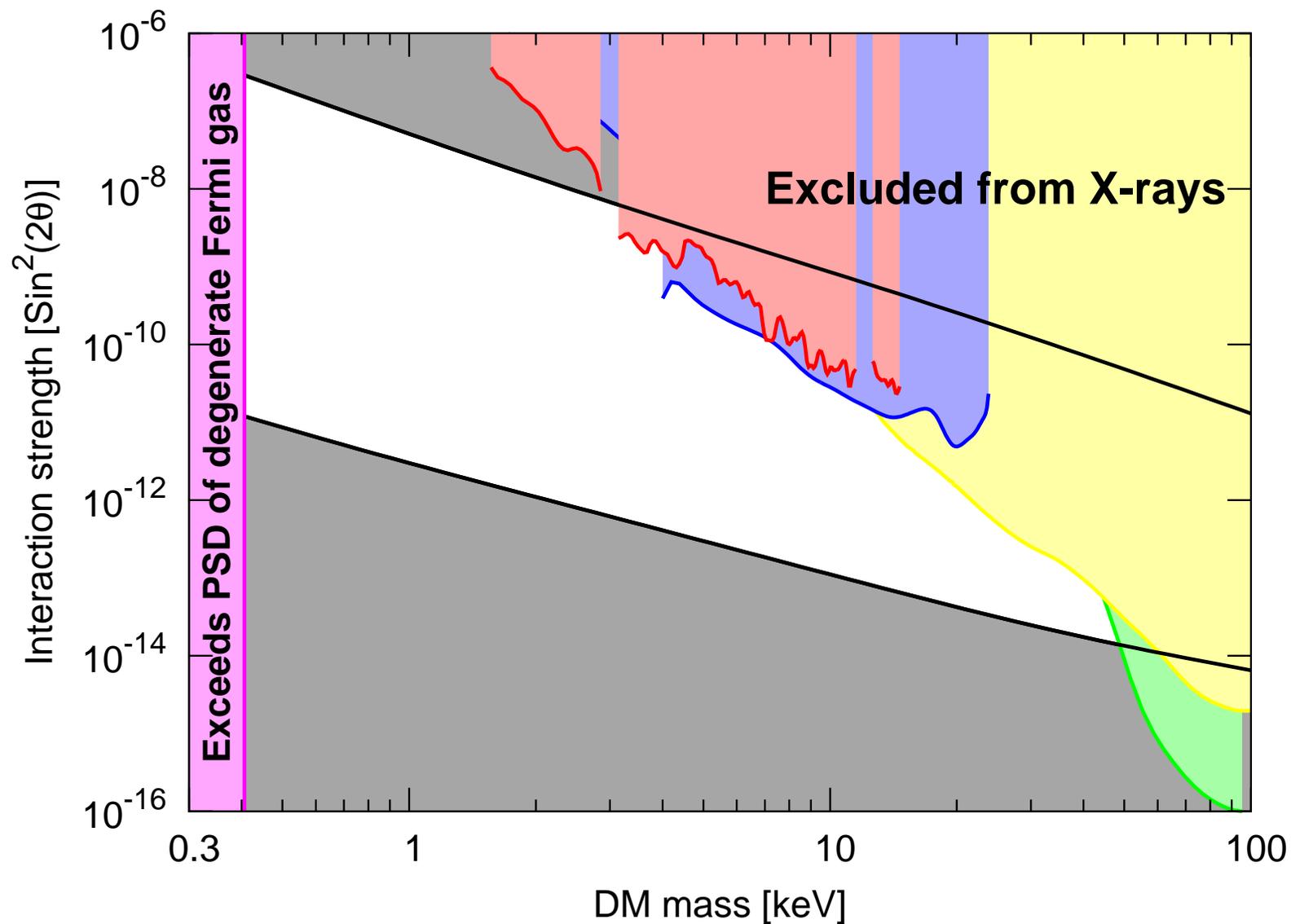
# Window of parameters of sterile neutrino DM

Asaka, Laine,  
Shaposhnikov

Laine,  
Shaposhnikov



# Window of parameters of sterile neutrino DM

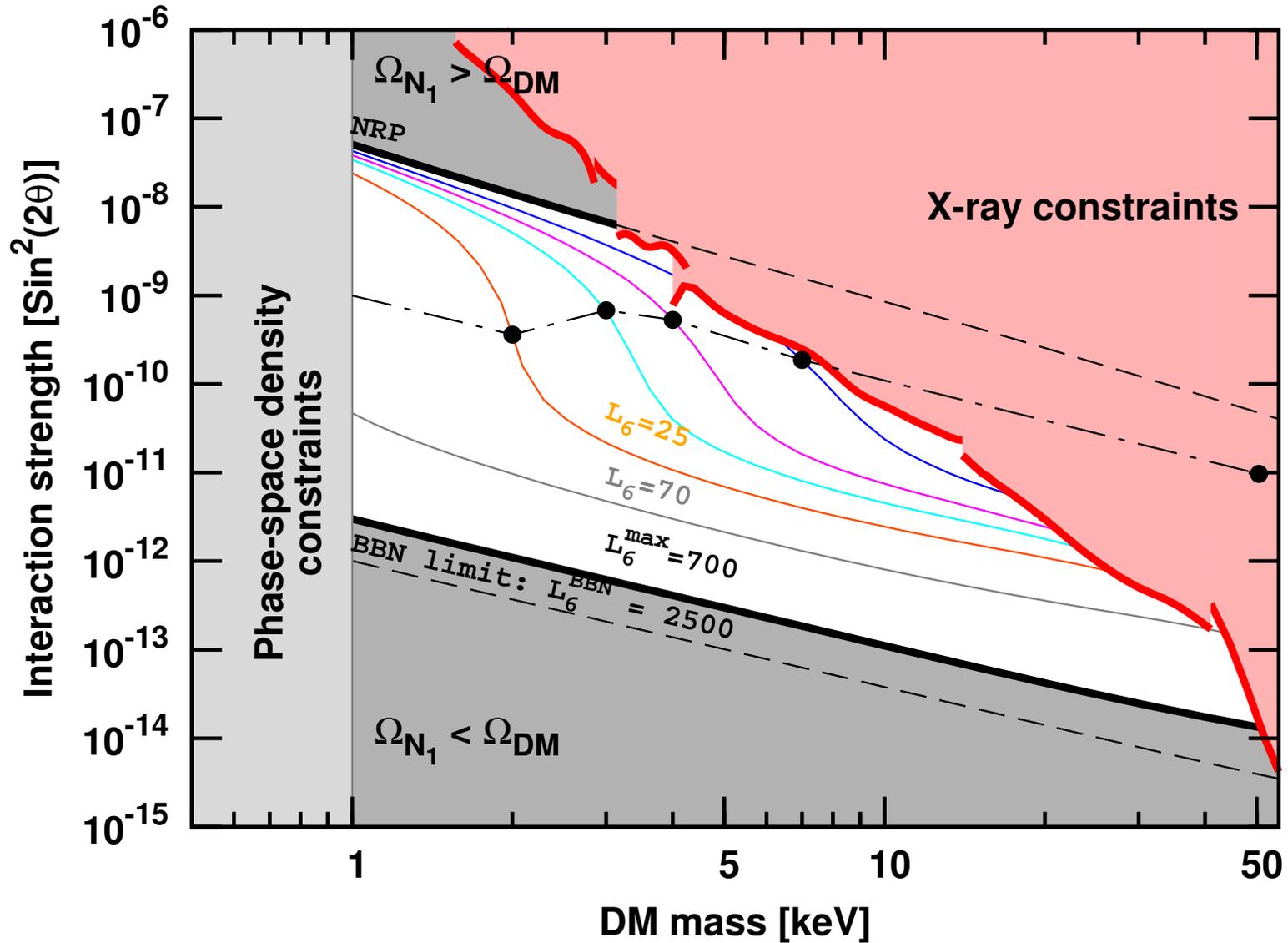


Asaka, Laine,  
Shaposhnikov

Laine,  
Shaposhnikov

O.R. and  
many others  
2005-2010

# Sterile neutrino DM in the $\nu$ MSM



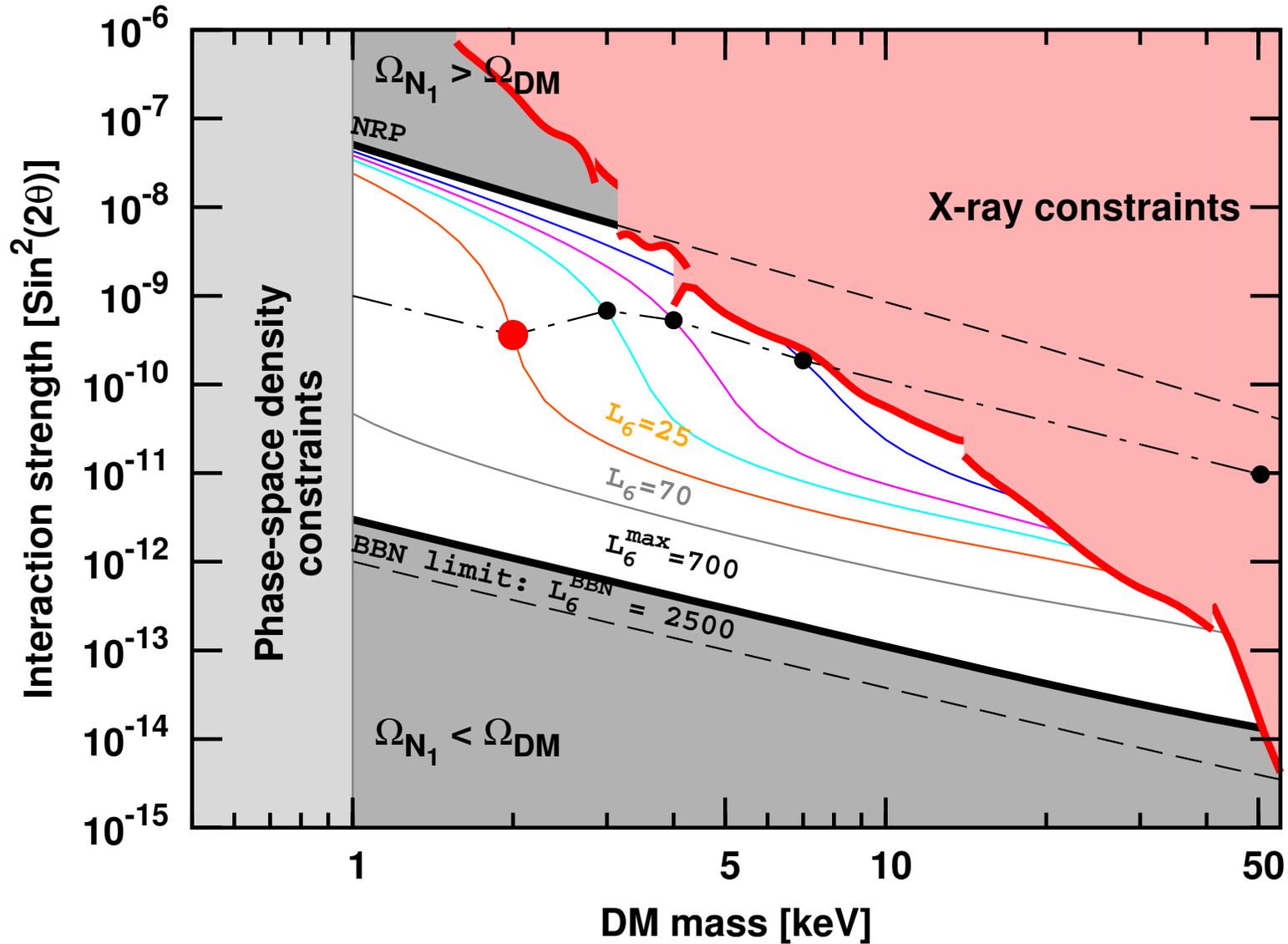
Boyarsky,  
O.R.,  
Lesgourgues,  
Viel PRL 2009

Review:  
[0901.0011]

# Sterile neutrino DM in the $\nu$ MSM

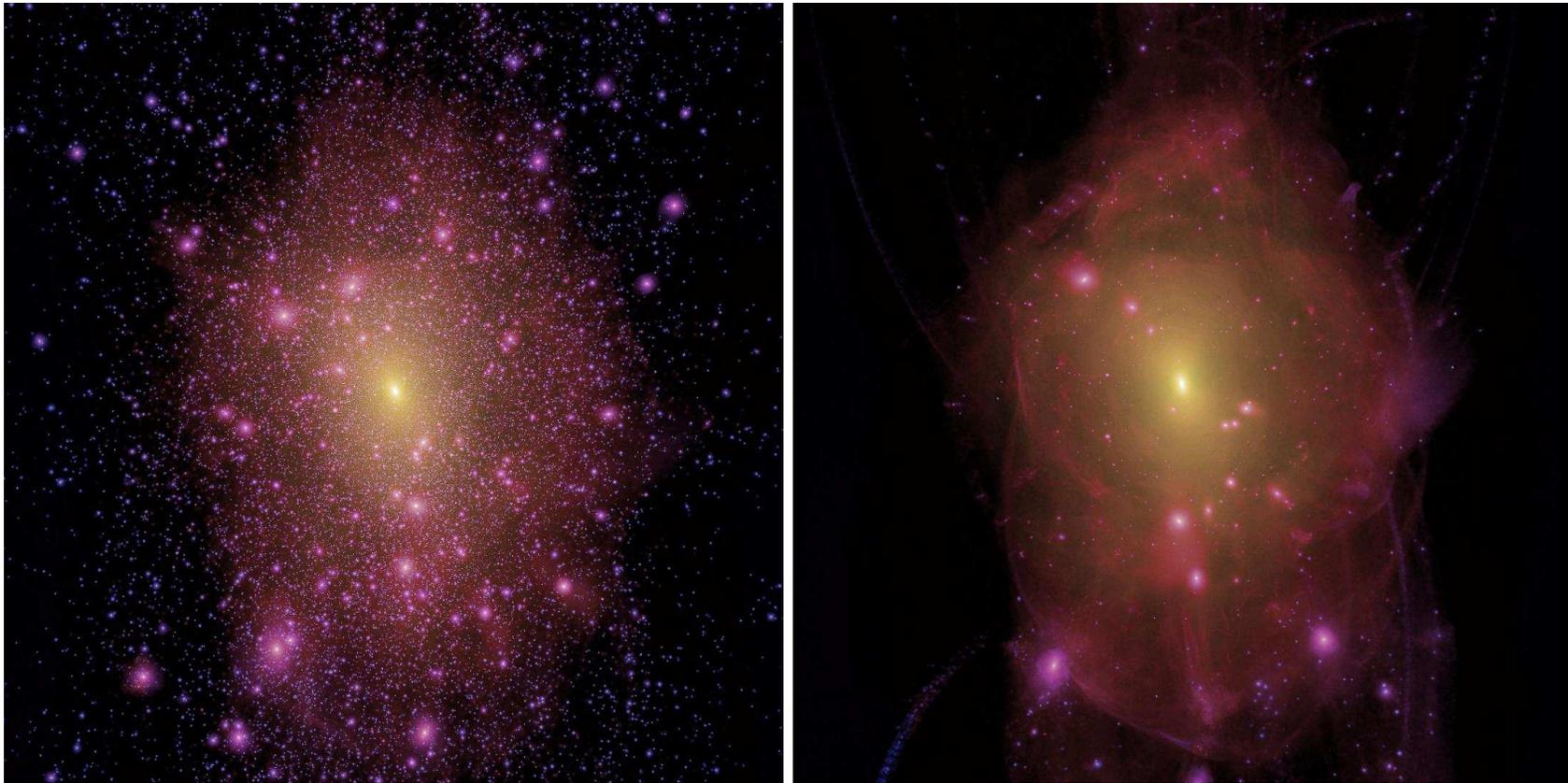
Boyarsky,  
O.R.,  
Lesgourgues,  
Viel PRL 2009

Review:  
[0901.0011]



# Halo substructure with sterile neutrino DM

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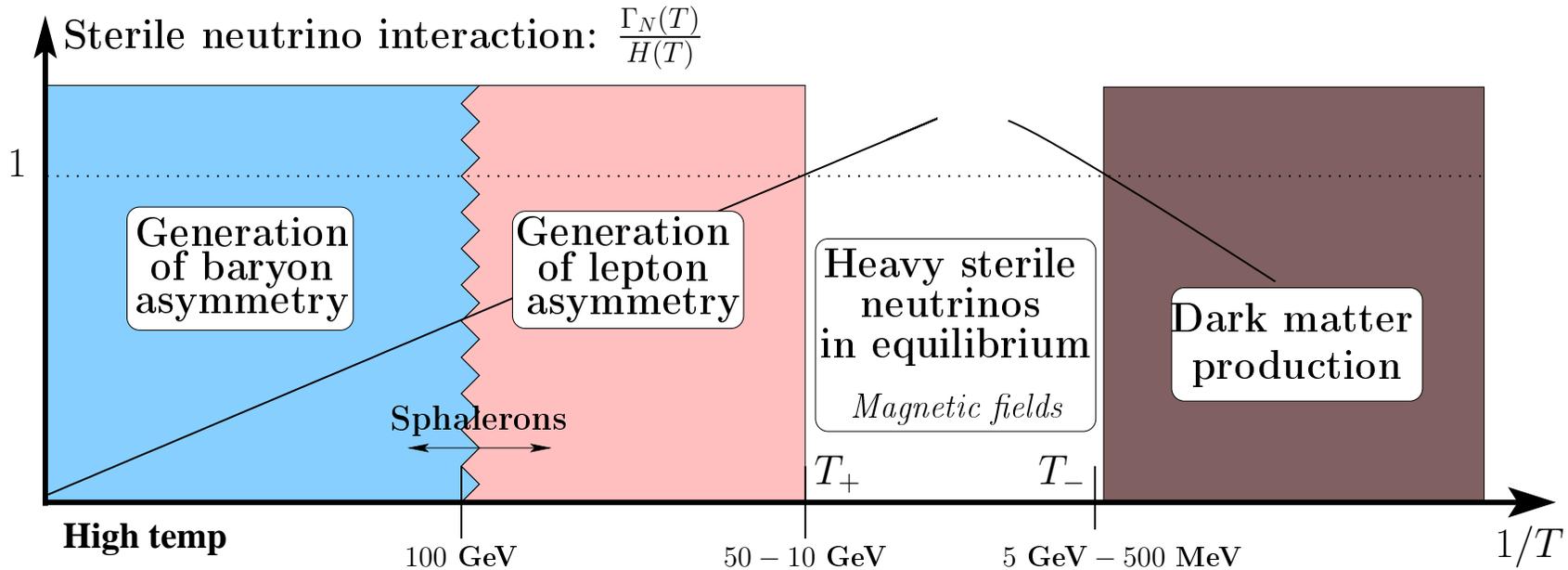
## *Aq-A-2 CDM halo*

Simulated sterile neutrino DM halo (right) is fully compatible with the Lyman- $\alpha$  forest data but provides a structure of Milky way-size halo different from CDM

*Aq-A-2 halo* made of sterile neutrino DM

(C. Frenk, T. Theuns, O.R., ...)

# History of the Universe in the $\nu$ MSM



- Magnetic fields in the plasma **relate** baryogenesis and sterile neutrino dark matter production
- Magnetic fields may be observable today in the intergalactic space

work in progress

## Magnetic fields in the $\nu$ MSM

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**Large lepton asymmetry** of the  $\nu$ MSM triggers **instability** in the Maxwell's equations and leads to the generation of large helical magnetic fields

- Magnetic fields, generated in the  $\nu$ MSM *below* 100 GeV are:
  - Maximally helical (sign of helicity determined by the sign of baryon asymmetry)
  - Energetic (magnetic energy density can be  $\sim$  total radiation density)
- Their survival until today is a matter of complicated evolution, described by **chiral magneto-hydrodynamics** work in progress

O.R.+  
Phys.Rev.Lett  
2012

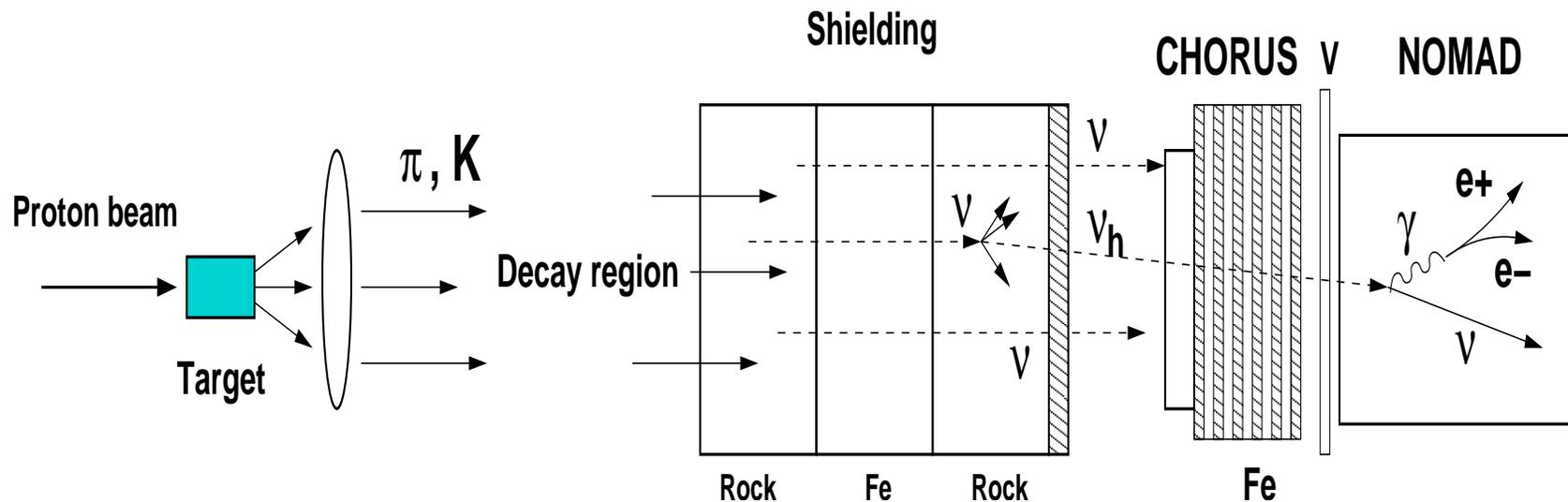
Boyarsky,  
O.R.,  
Shaposhnikov  
[1204.3604]

The leptogenesis in the  $\nu$ MSM leads to the **baryogenesis** and generation of **potentially observable** cosmological magnetic fields

**How can we search for these particles?**

# Peak searches and fixed-target experiments

$M_I < 1 \text{ MeV}$	$M_I \gtrsim 1 \text{ MeV}$	$M_I \gtrsim 140 \text{ MeV}$	...
$N_I \rightarrow \nu\nu\bar{\nu}$	$N_I \rightarrow \nu e^+ e^-$	$N_I \rightarrow \pi^\pm e^\mp$	
$N_I \rightarrow \nu\gamma$		$N_I \rightarrow \pi^0\nu$	



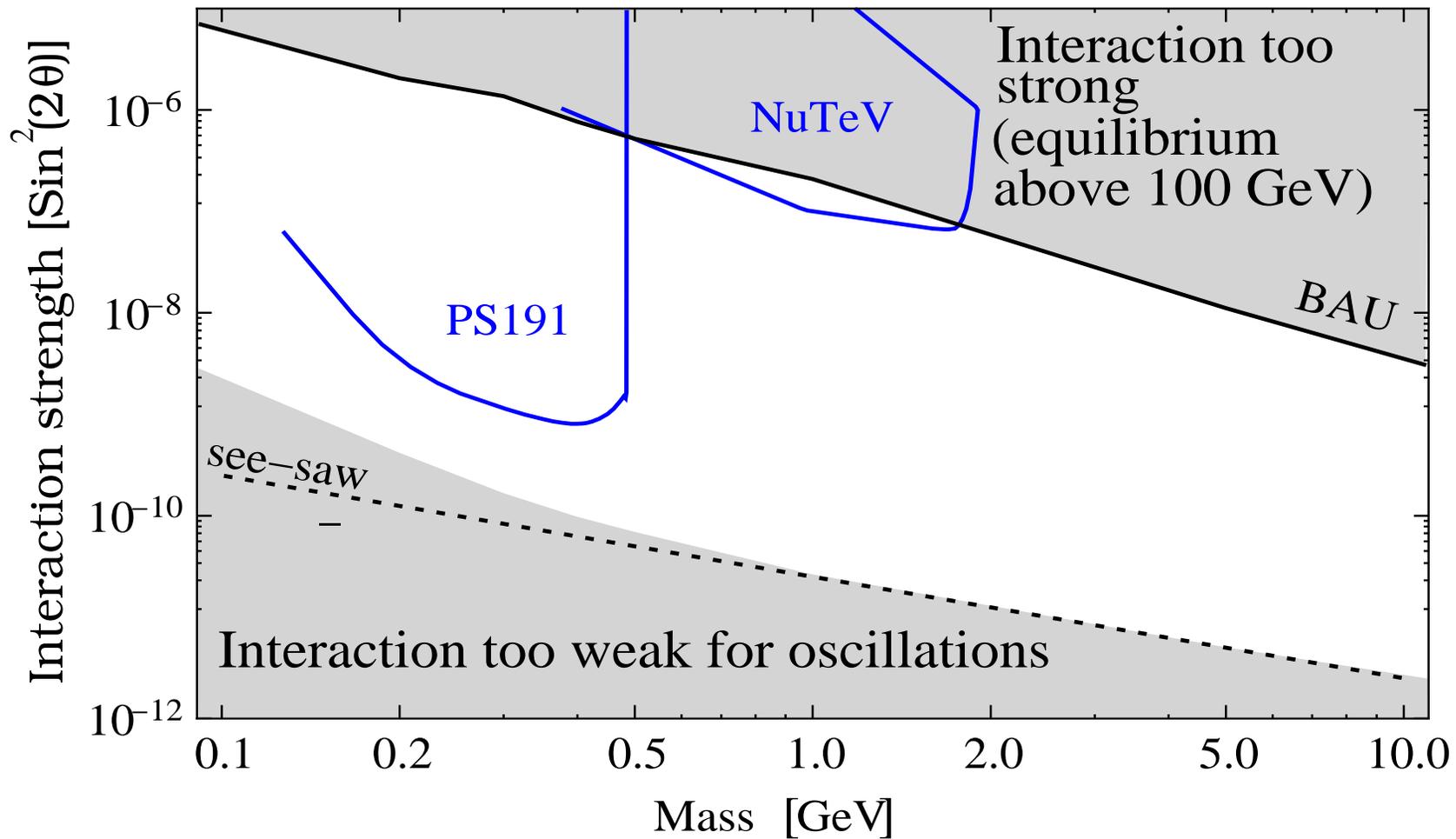
## Peak searches:

- SIN  $\pi M3$ , Switzerland – 1981
- KEK K3, Japan, 1982
- TRIUMF M13, Canada, 1992
- TRIUMF PIENU, Canada, 2011

## Fixed-target searches:

- PS191, CERN – 1984
- CHARM, CERN – 1985
- NuTeV, Fermilab – 1996-1997

# Parameter space of sterile neutrinos

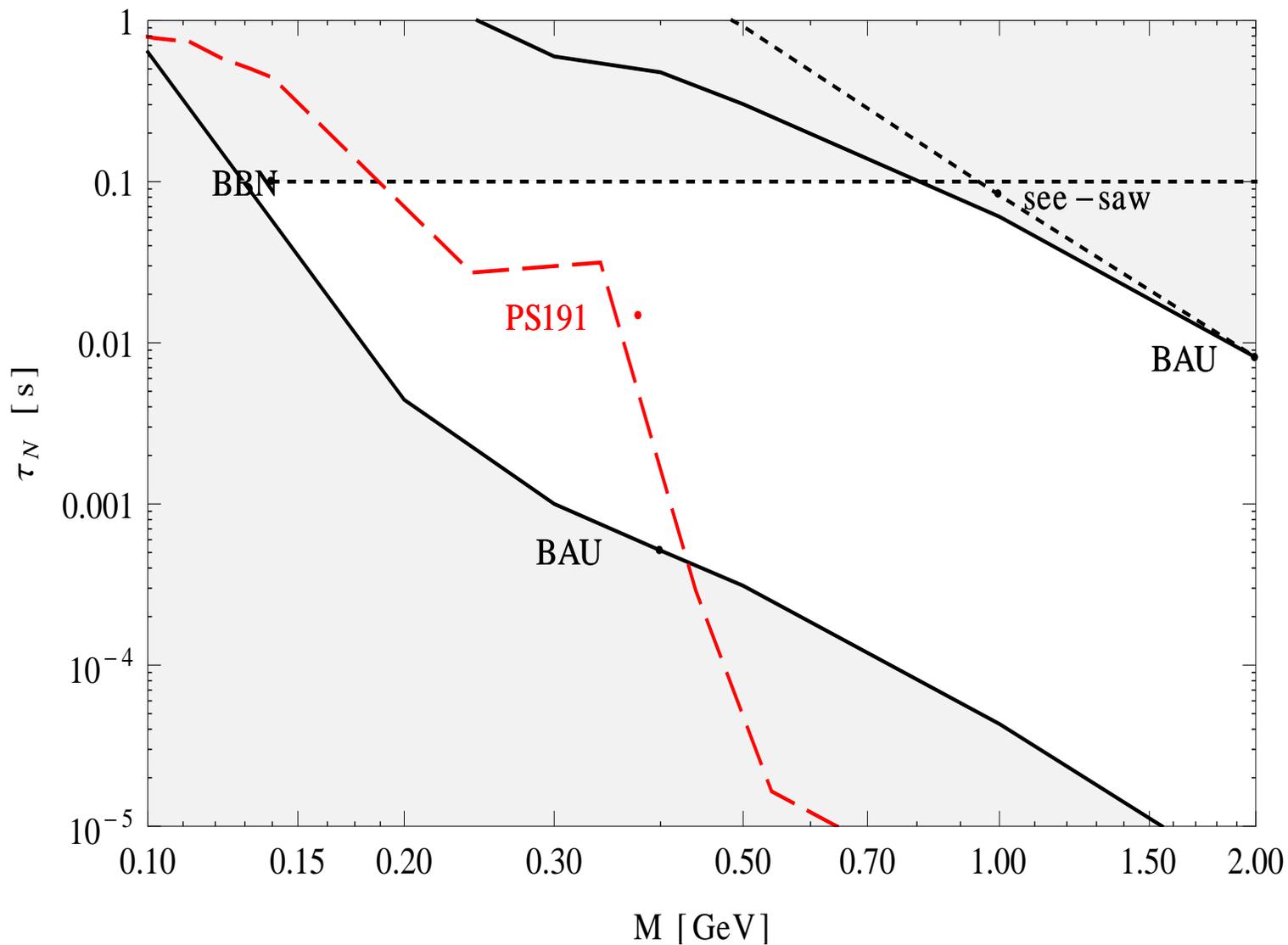


Asaka,  
Canetti,  
Gorbunov,  
Shaposhnikov,  
2005–2011;

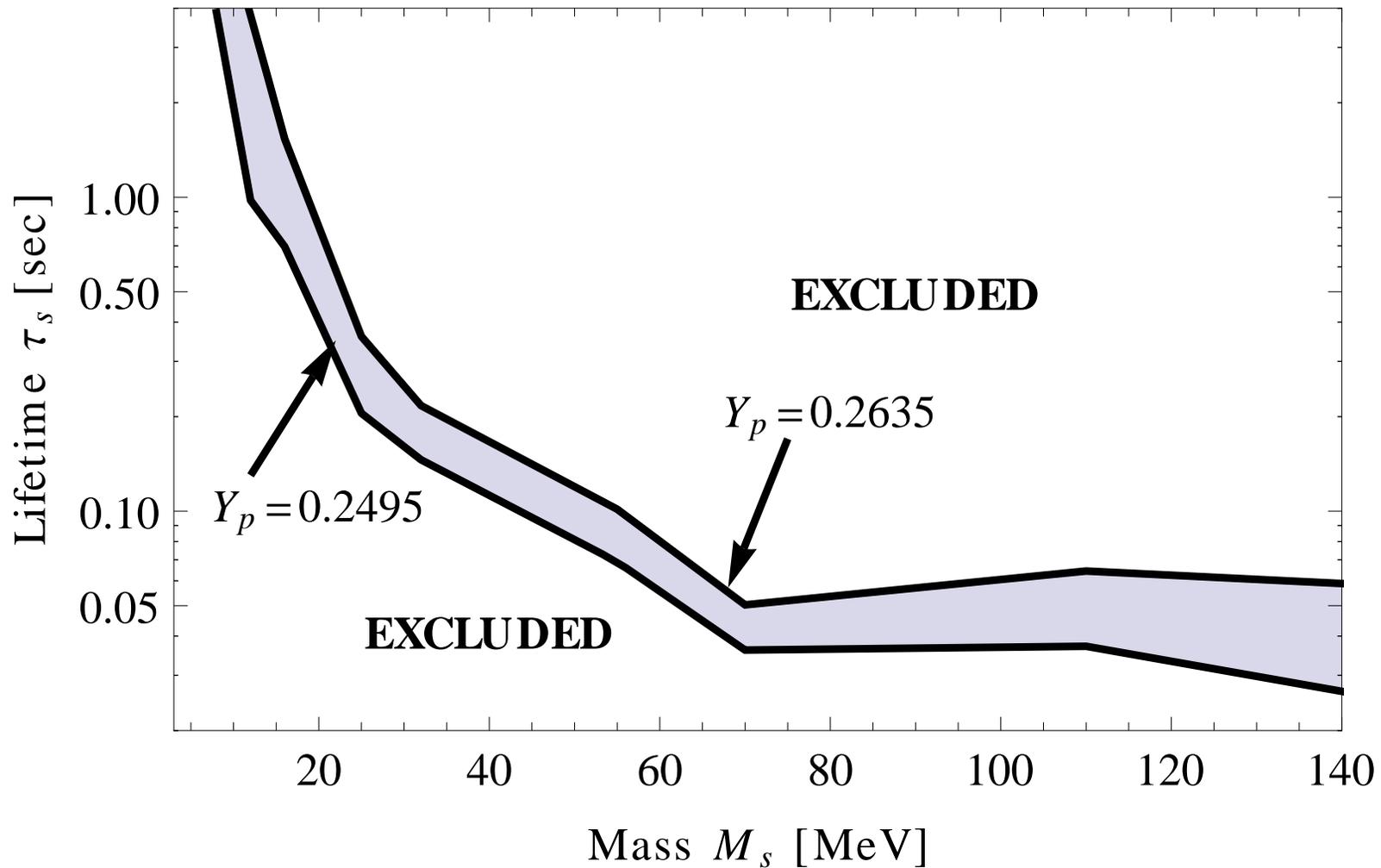
O.R & Ivashko  
[1112.3319] –  
revised  
accelerator  
bounds

# Lifetime of sterile neutrinos

Canetti &  
Shaposhnikov  
(2011)



# Sterile neutrinos and $^4\text{He}$ abundance



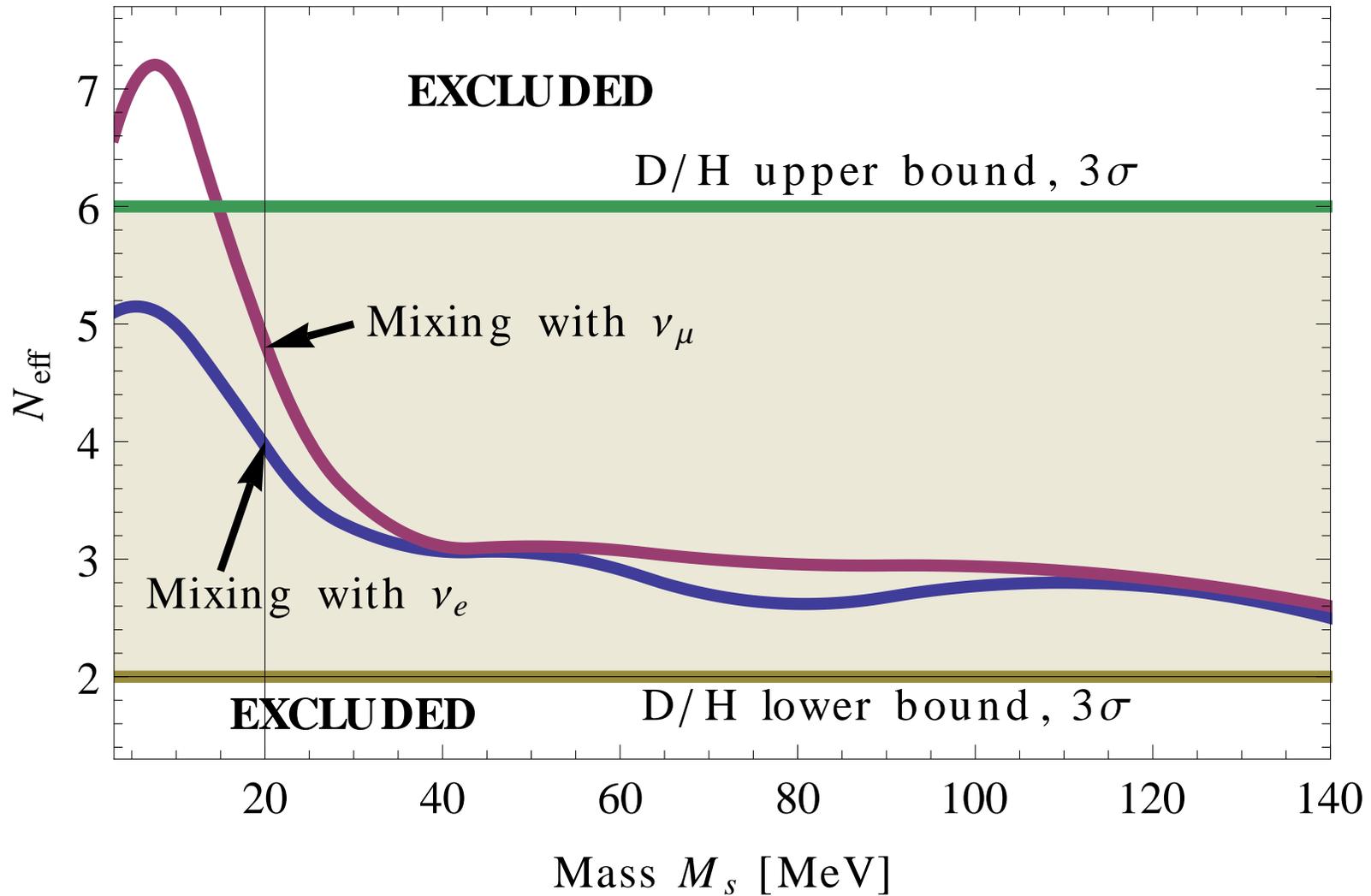
O.R. & Ivashko  
[1202.2841]

$2\sigma$  bounds  
based on  
Izotov &  
Thuan 2010

Decay of sterile neutrinos increases Helium-4 abundance

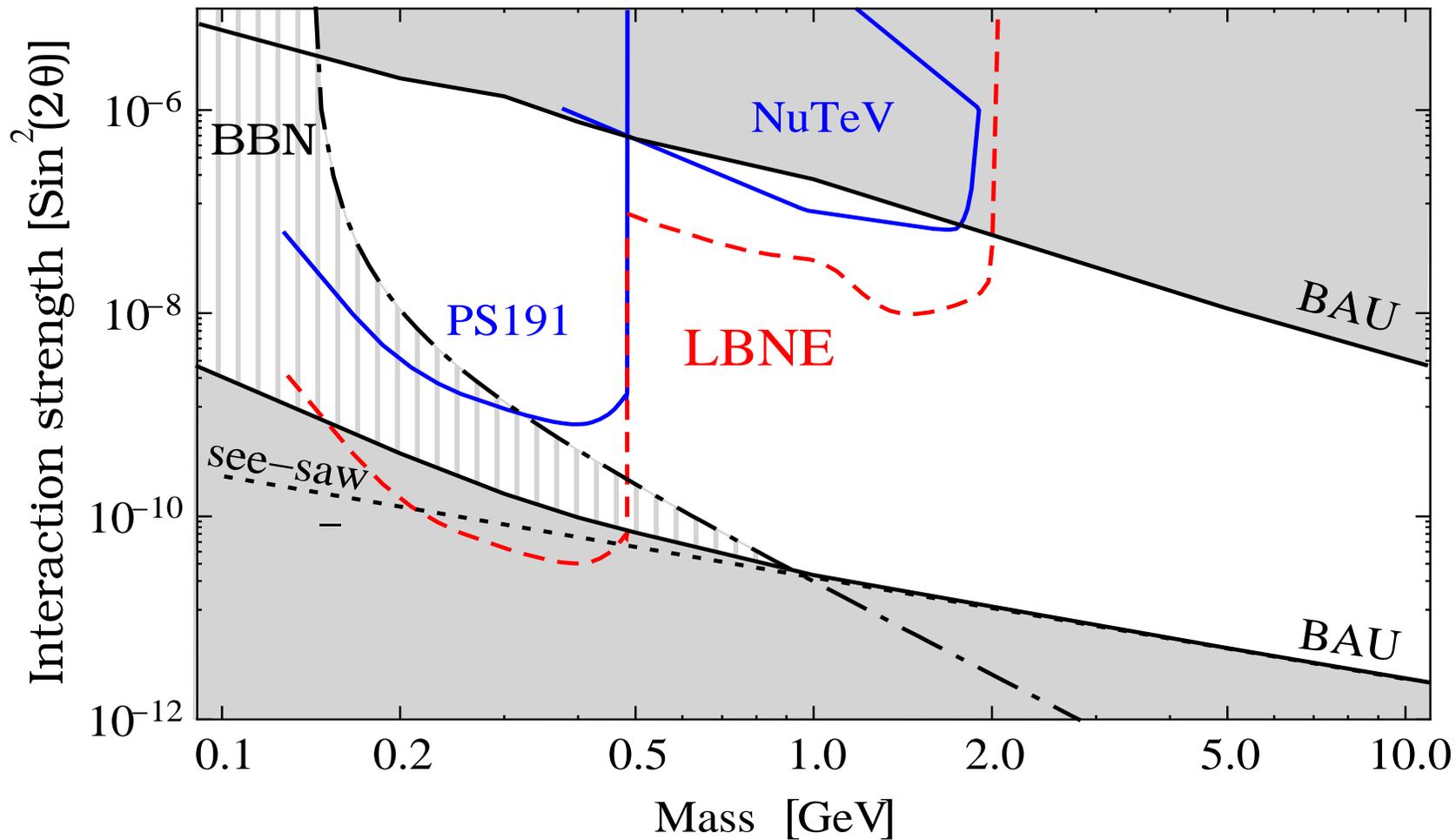
# Sterile neutrinos and $N_{\text{eff}}$

O.R. & Ivashko  
[1202.2841]



Decay of sterile neutrinos affects  $N_{\text{eff}}$

# Parameter space of sterile neutrinos



Gorbunov,  
Shaposhnikov  
(2009);

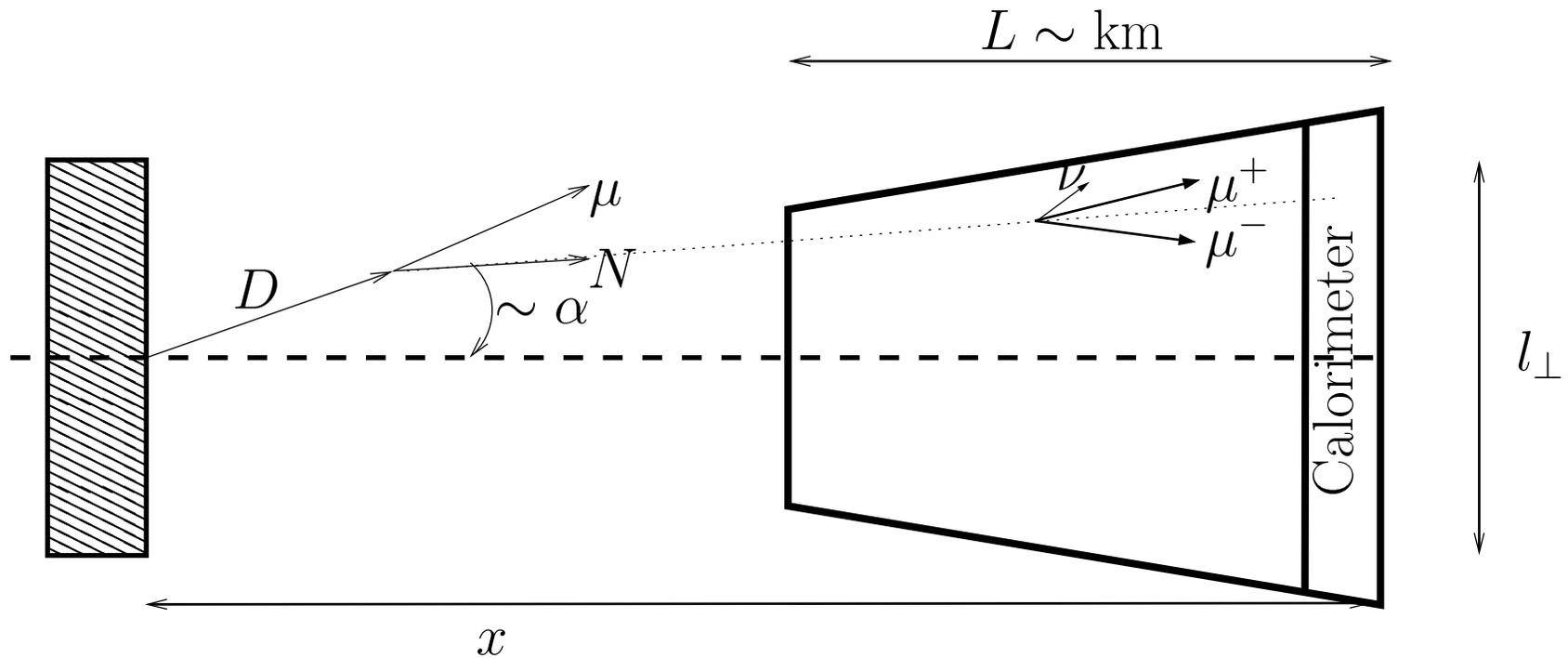
O.R. & Ivashko  
[1112.3319] –  
revised  
accelerator  
bounds

O.R. & Ivashko  
[1202.2841] –  
BBN bounds

LBNE white  
paper  
[1110.6249]

White paper  
on sterile  
neutrinos  
[1204.5379]

# Ultimate detector



To find MeV-GeV scale sterile neutrinos responsible for oscillations (and possibly playing an important role in the early Universe) we need:

- High energy/high intensity fixed-target experiments
- Large detector (length  $\sim \text{km}$ )

# "Nightmare scenario"

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## Predictions for the nearest future

- Standard Model Higgs with the mass above  $\sim 125$  GeV at LHC and no new physics otherwise
- Primordial spectral index  $n_s = 0.96 \dots$  correlated with the Higgs mass, measured by Planck
- Non-detection of tensor modes
- Sum of neutrino masses  $\sum m_\nu \approx (1 - 2)m_{\text{atm}}$   
(hard to see it until the next generation of CMB experiments)
- In the  $0\nu\beta\beta$  mass  $m_{\beta\beta}$  at the level  $1 - 10$  meV

## Within the next decade we can:<sup>1</sup>

- **Discover** two sterile neutrinos with the masses  $\mathcal{O}(100)$  MeV  $\div$  few GeV and mass splitting  $\sim m_{\text{atm}}$  in “intensity frontier” experiments  
(NA62 in CERN, LBNE, SLHCb)
- **Discover** dark matter decay line with mass/lifetime consistent with the parameters of two other sterile neutrinos  
(the first X-ray spectrometer of the new generation will fly in 2014).
- **Detect** the matter power spectrum suppression at  $k \sim 10$  h/Mpc  
(next round of weak lensing/Lyman- $\alpha$  forest experiments)
- **Find** the strength/correlation length of magnetic fields in voids consistent with params. of sterile neutrinos — **direct observational signature of baryogenesis**, *4th pillar of hot Big Bang*

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<sup>1</sup>By “can” I mean “experimental technologies are available for us today”

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THANK YOU FOR YOUR ATTENTION

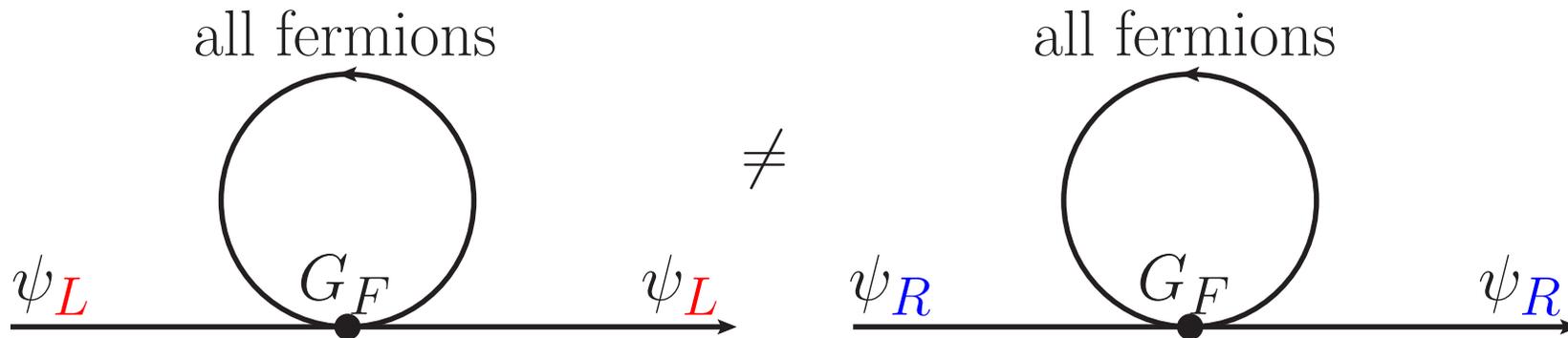
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# Additional slides

## Weak corrections

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- Chemical potentials of all left/right particles are in equilibrium at  $T < 80 \text{ TeV}$
- Weak corrections are proportional to the **asymmetry** of all fermions, running in the loops



- This diagram leads to the change of dispersion relations (shift of chemical potentials) of left/right particles

Boyarsky,  
O.R.,  
Shaposhnikov  
[1204.3604]

## Chern-Simons term

---

- At finite value of baryon and lepton asymmetry effective action for magnetic fields acquires a **Chern-Simons term**:

$$\mathcal{F}[\vec{A}] = \frac{1}{2} \int d^3x \left( \vec{B}^2 + \Pi_2 \vec{A} \cdot \vec{B} \right)$$

- The Chern-Simons coefficient is given by:

$$\Pi_2 = \frac{\alpha}{2\pi} G_F \times (c_1 \text{ baryon number} + c_2 \text{ lepton numbers}) \neq 0$$

Boyarsky,  
**O.R.**,  
Shaposhnikov  
[1204.3604]

- Its origin is **chiral anomaly** and parity-violating nature of weak interactions

## CS term and instability

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$$\mathcal{F}[\vec{A}] = \frac{1}{2} \int d^3x \left( \vec{B}^2 + \Pi_2 \vec{A} \cdot \vec{B} \right)$$

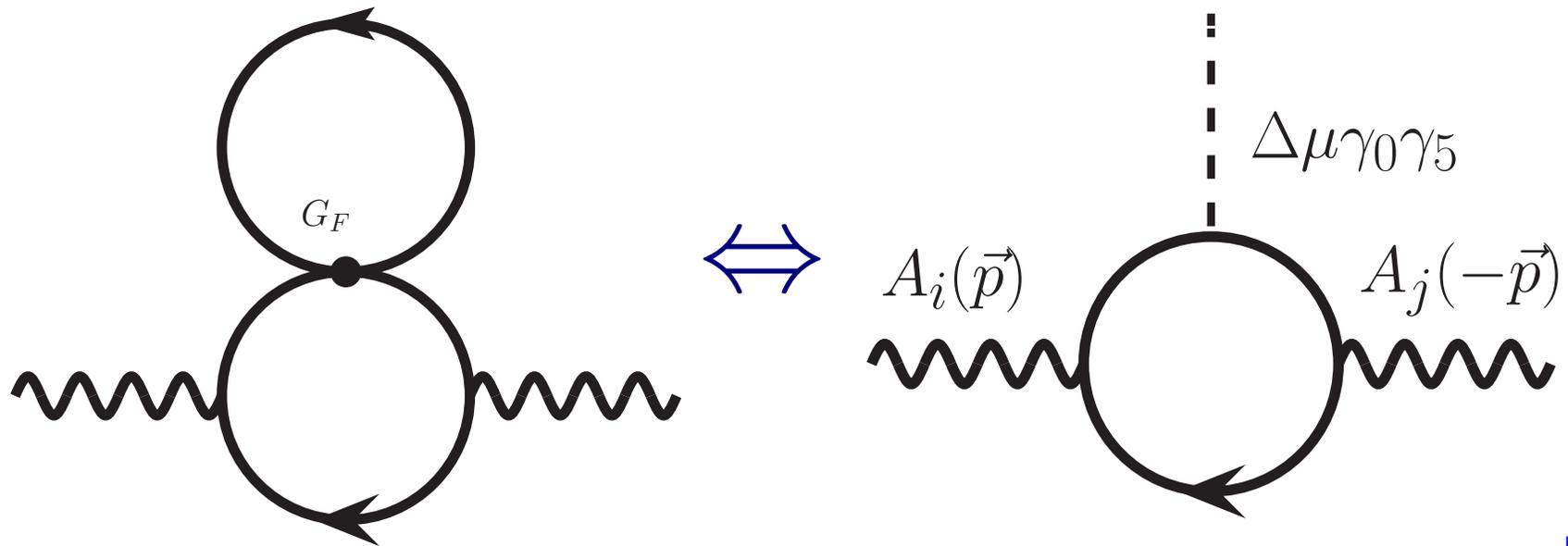
- The unstable mode will have a form  $\vec{A}(\vec{x}) = A_0 (\cos(pz), \sin(pz), 0)$
- On such a configuration  $\vec{B}^2 = -p \vec{A} \cdot \vec{B}$  and are homogeneous
- The effective action is unbounded from below for  $p < \Pi_2$

$$\mathcal{F}[A] = \frac{1}{2} \int d^3x \left( p^2 - p\Pi_2 \right) A_0^2 < 0$$

- **Long-range magnetic fields** will be spontaneously generated in the equilibrium

## Chern-Simons term

---



Boyarsky,  
O.R.,  
Shaposhnikov  
[1204.3604]

Weak interactions make propagators left-right asymmetric and chiral anomaly leads to the generation of parity-odd term

$$\Pi_2(0) = \frac{\alpha}{2\pi} G_F \times (c_1 \text{ baryon number} + c_2 \text{ lepton numbers}) \neq 0$$

## Maximally helical configuration

---

- The unstable mode will have a form

$$\vec{A}(\vec{x}) = A_0 \left( \cos(pz), \sin(pz), 0 \right)$$

- The magnetic field

$$\vec{B}(\vec{x}) = -p\vec{A}(\vec{x})$$

— is maximally helical

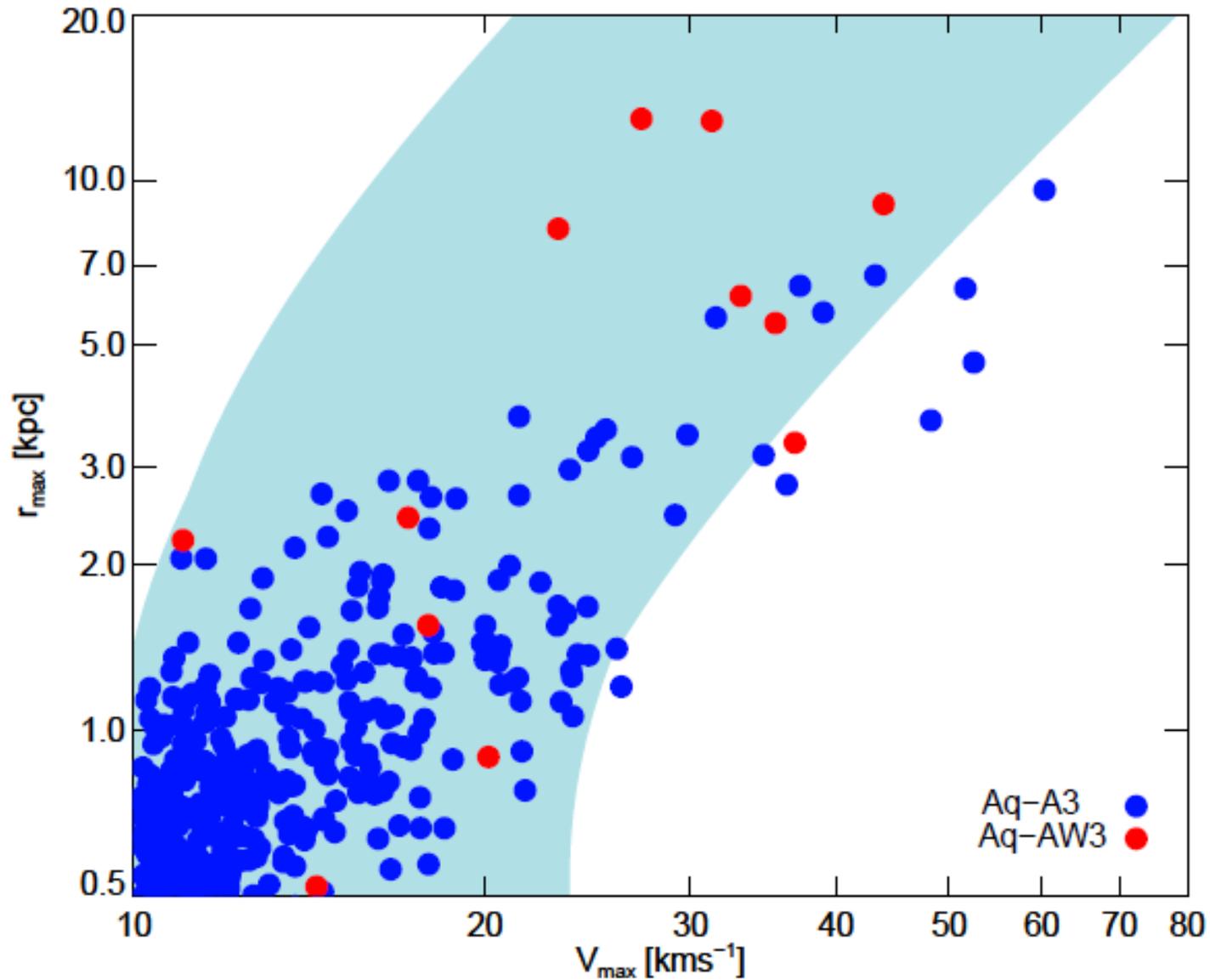
- On this configuration  $\vec{B}^2 = p\vec{A} \cdot \vec{B}$  and are homogeneous

- The effective action:

$$\mathcal{F}[A] = \frac{1}{2} \int d^3x \left( p^2 - p\Pi_2 \right) A_0^2 < 0$$

for  $p < \Pi_2$

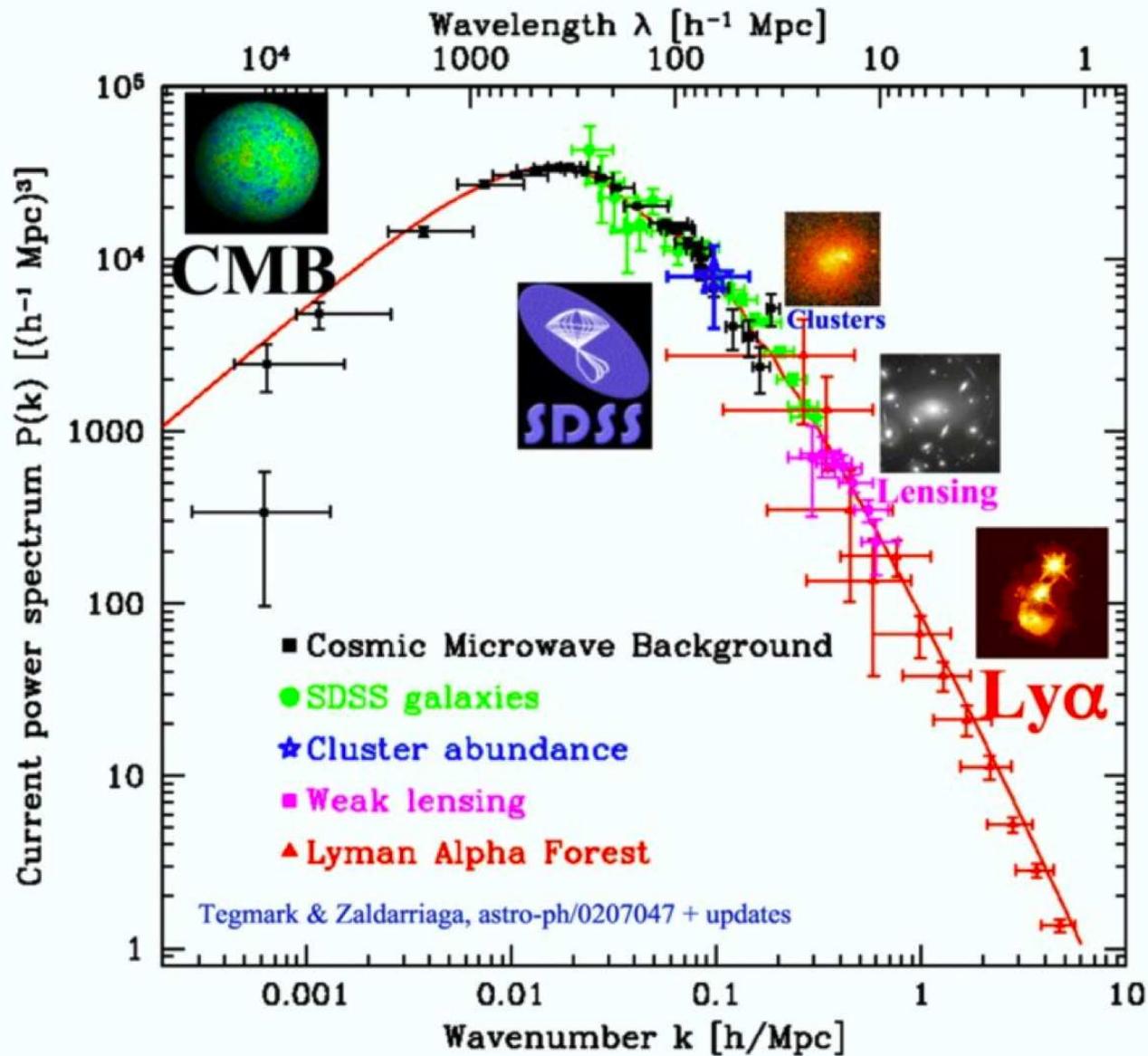
# Abundance of large satellites



Strigari, Frenk  
White (2011)

Lovell, Frenk,  
Eke, ...,  
Boyarsky, O.R.  
1104.2929  
[astro-ph.CO]

# How to measure power spectrum

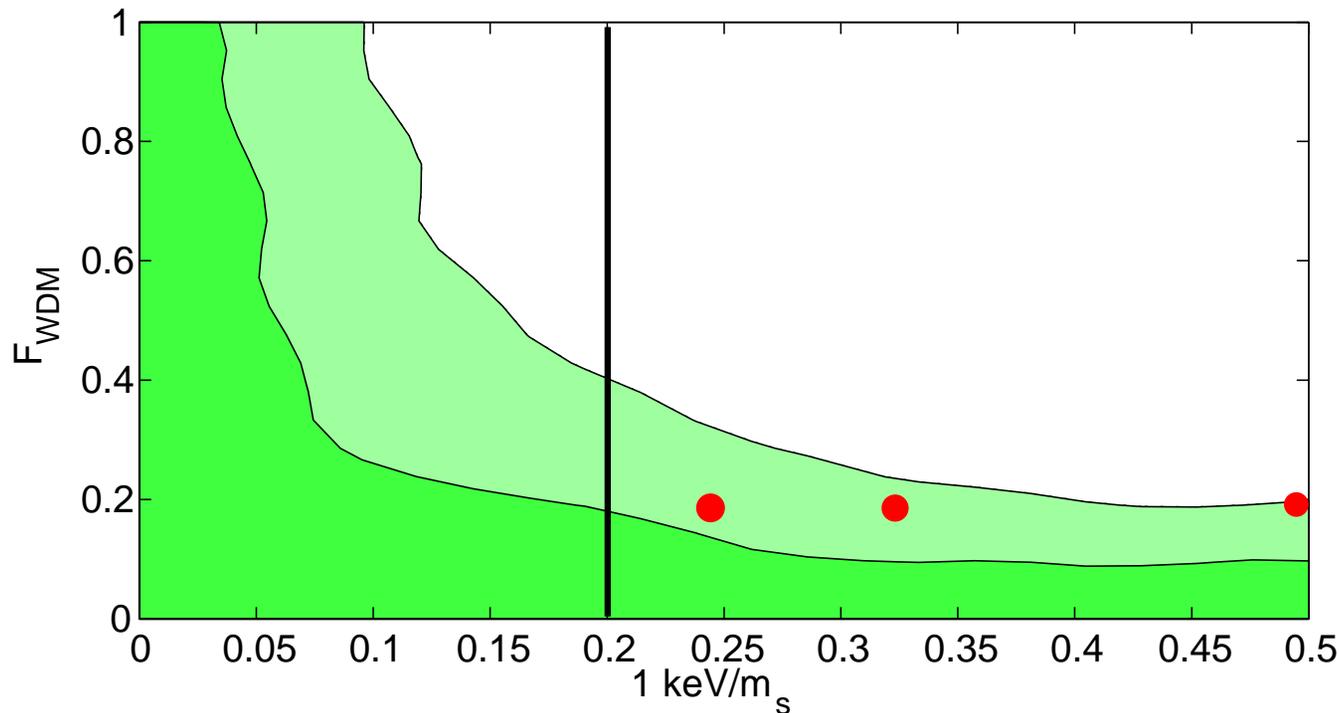


Max Tegmark  
Univ. of Pennsylvania  
max@physics.upenn.edu  
TAUP 2003  
September 5, 2003

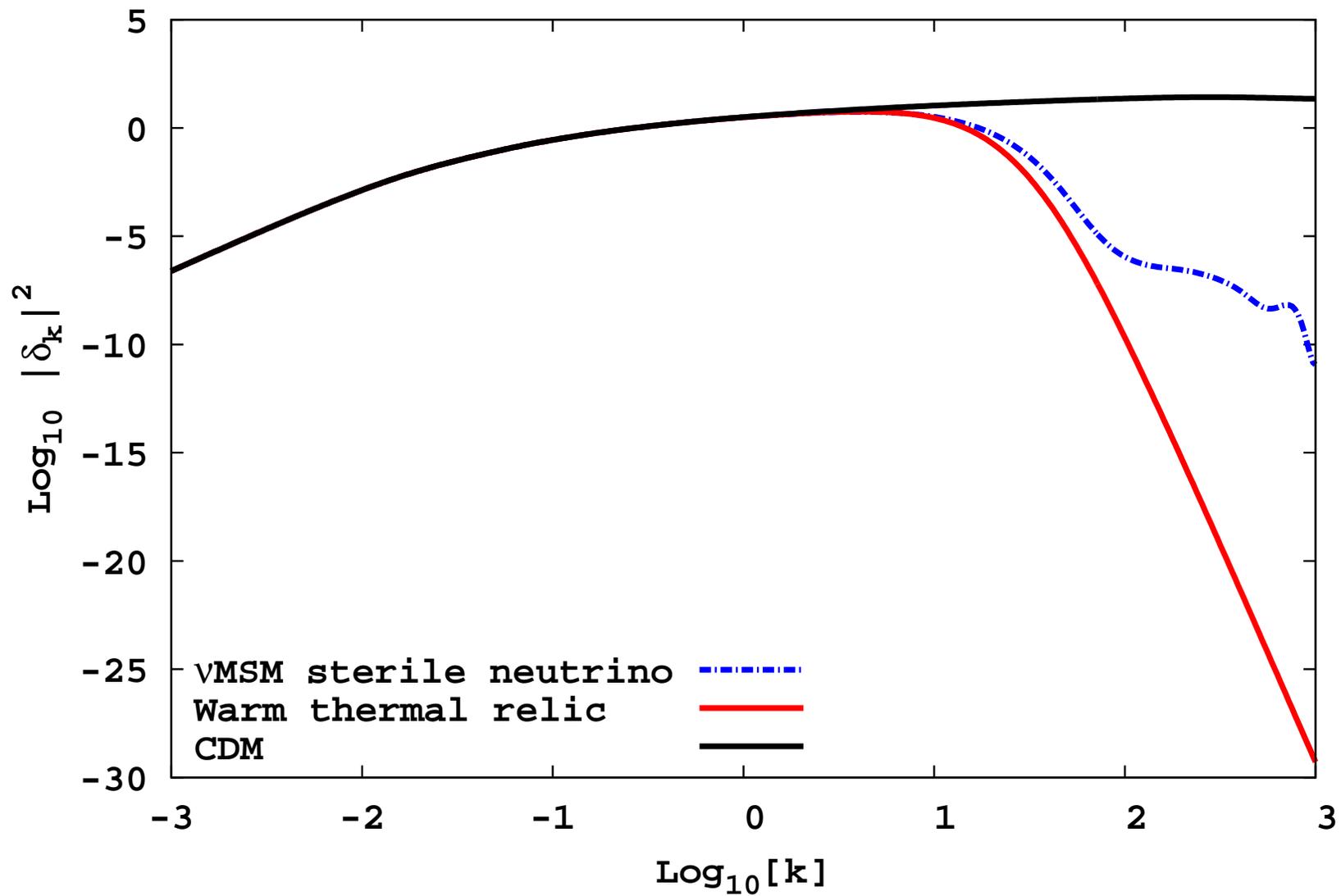
# Lyman- $\alpha$ bounds for sterile neutrinos

- Revised version of these bounds in CDM+WDM (mixed, CWDM) models demonstrates that
  - The primordial spectra **are not described by free-streaming**
  - There exist viable sterile neutrino DM models with the masses as low as 2 keV

Boyarsky,  
**O.R.**,  
Lesgourgues,  
Viel JCAP &  
PRL (2009)

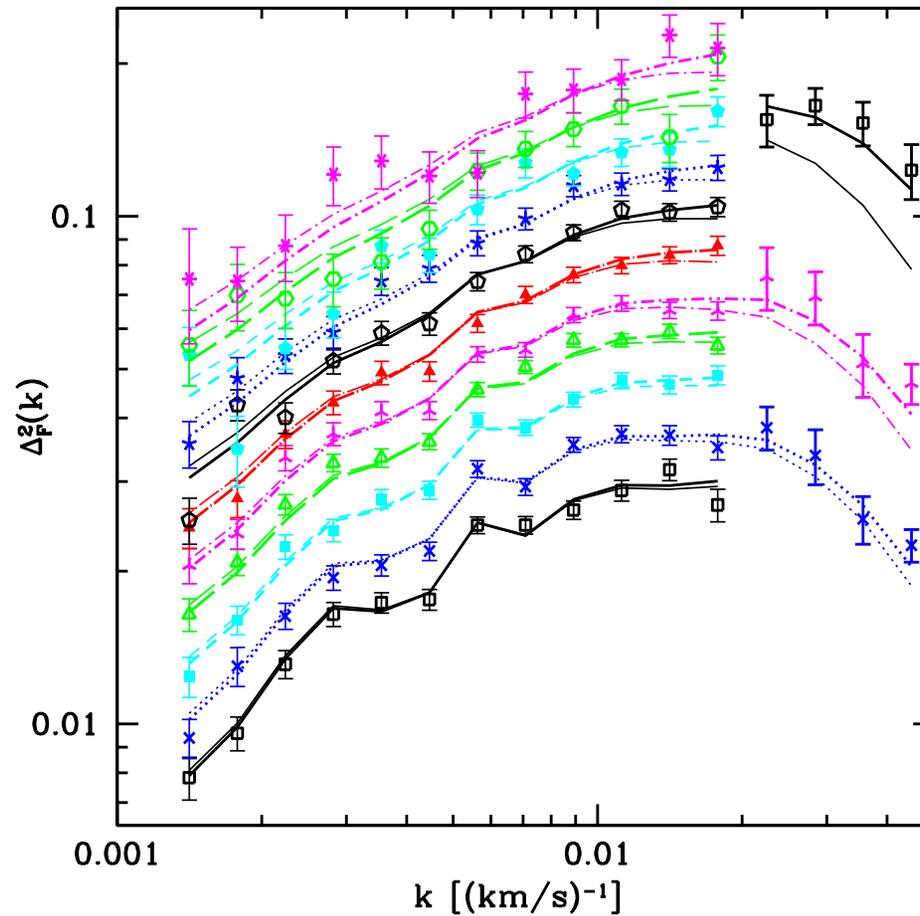


# Suppression of power spectrum



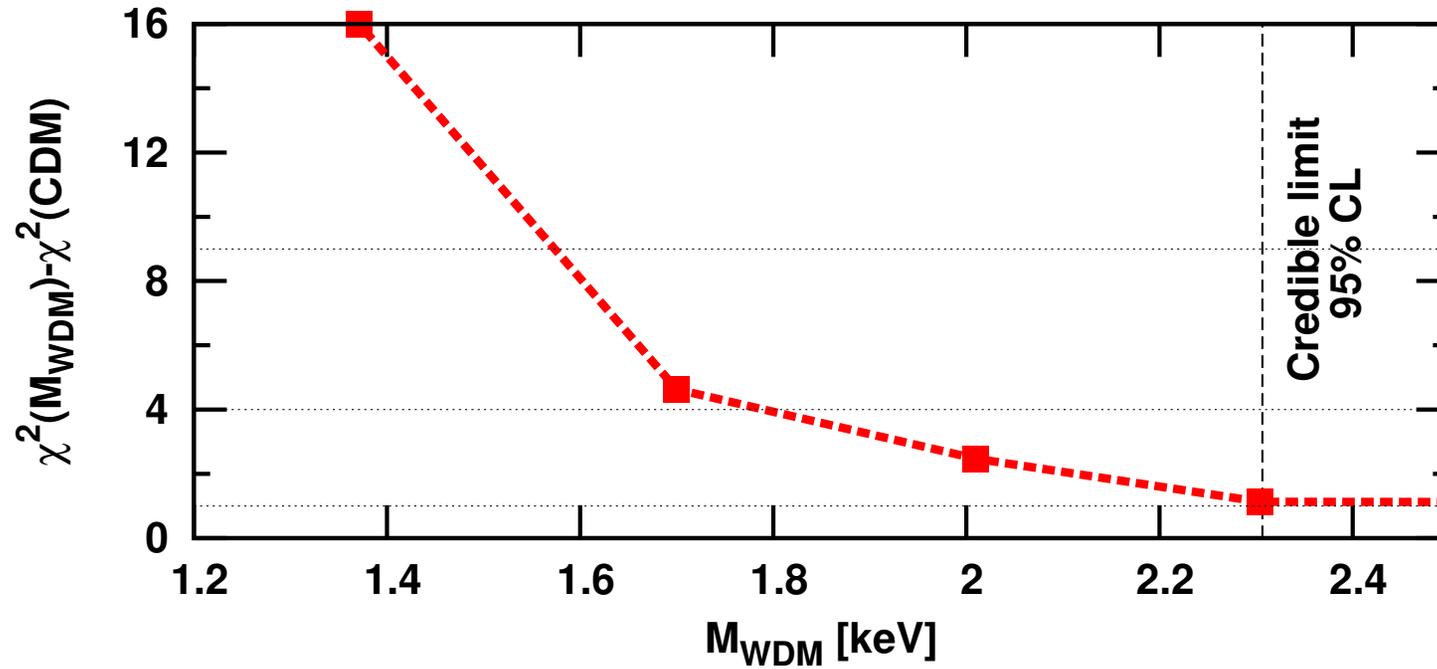
# Lyman- $\alpha$ forest flux power spectrum

Seljak et al.  
'06



Measured flux power spectrum is compared against CDM and non-CDM models

# Ly- $\alpha$ and thermal relics



Boyarsky,  
Lesgourgues,  
**O.R.**, Viel  
[0812.0010]  
(JCAP 2009)

Also Viel et al.  
2005-2007;

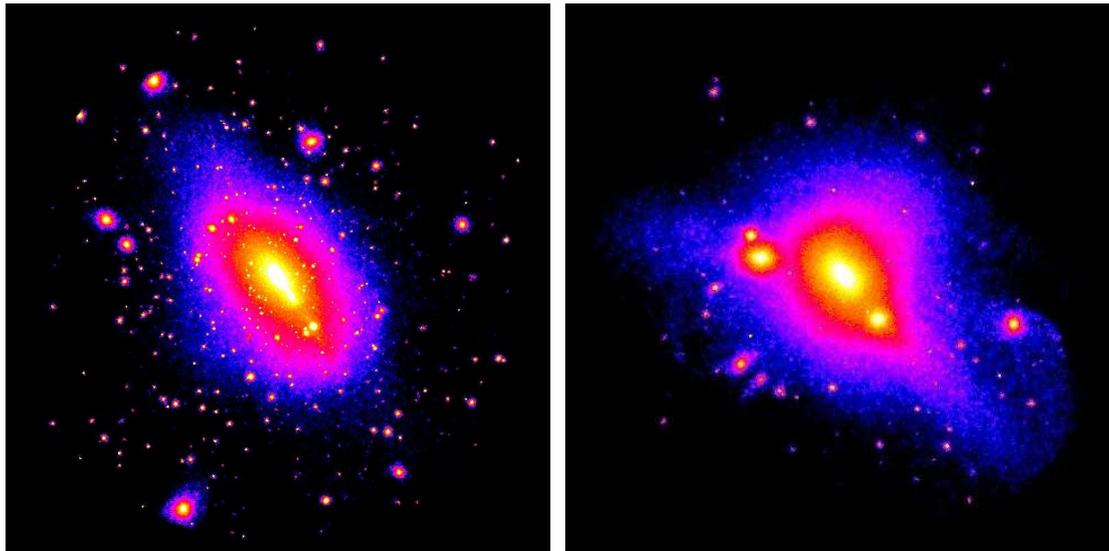
Seljak et al.  
(2006)

*These bounds are for **thermal relics** only!*

## Lyman- $\alpha$ forest and warm DM

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- Previous works put bounds on free-streaming  $\lambda_{FS} \lesssim 150$  kpc (“WDM mass”  $> 2.3$  keV) Viel et al. 2005-2007; Seljak et al.(2006)
- The simplest **WDM** with such a free-streaming would not modify visible substructures:



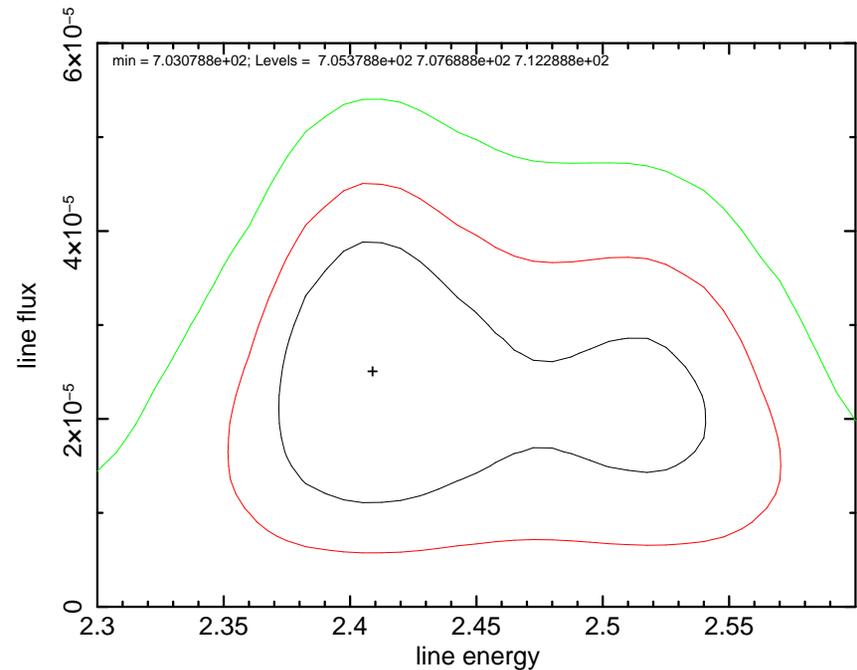
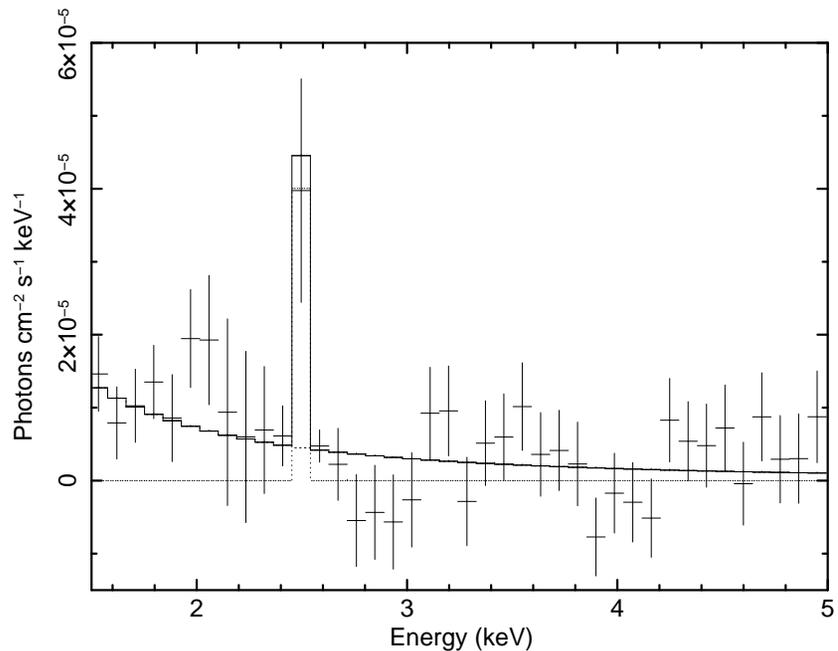
Maccio & Fontanot (2009);

Polisensky & Ricotti (2010)

- **Thermal relic** with exponential cut-off  $\sim 1$  Mpc would erase **too many substructures**. Anything “colder” would produce enough structures to explain observed Milky Way structures

# Checking DM origin of a line

- *Dark Matter Search Using Chandra Observations of Willman 1, and a Spectral Feature Consistent with a Decay Line of a 5 keV Sterile Neutrino* Loewenstein & Kusenko (Dec'2009)

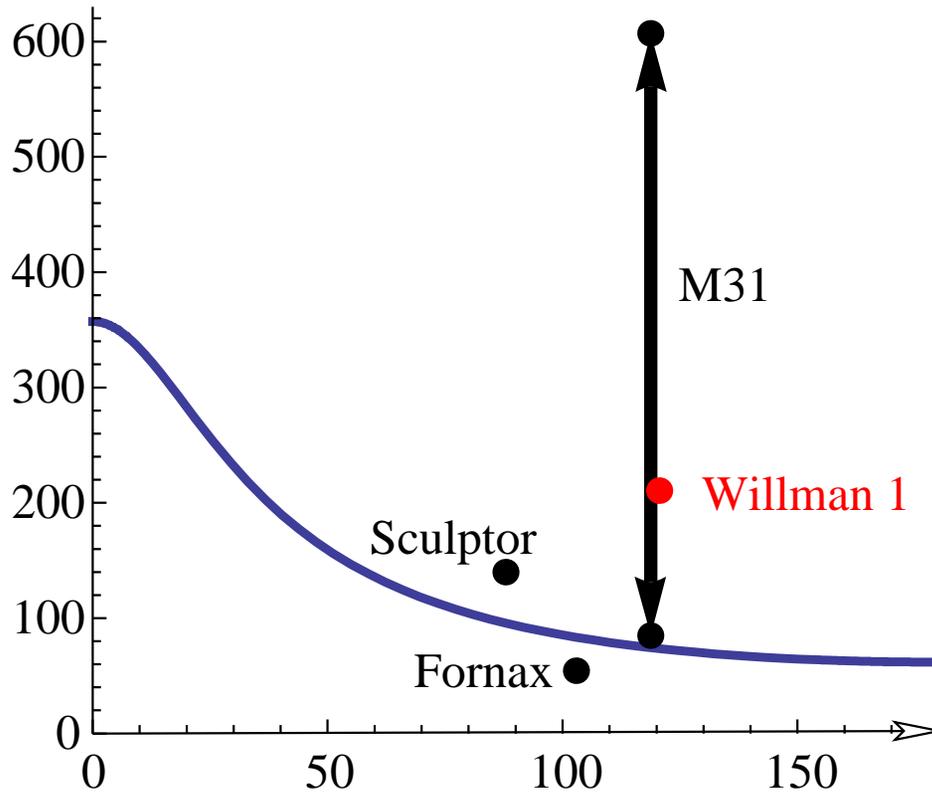


68%, 90% and 99% confidence intervals

- *Can the excess in the FeXXVI Ly gamma line from the Galactic Center provide evidence for 17 keV sterile neutrinos?* Prokhorov & Silk (Jan'2010)

## Do we see this line anywhere else?

$S_{MW}$  Msun/pc<sup>2</sup>



Objects with comparable expected signal for which archival data is available

■ **Fornax dSph (XMM)**

$$S_F = 54.4 M_{\odot} \text{pc}^{-2}$$

■ **Sculptor dSph (Chandra)**

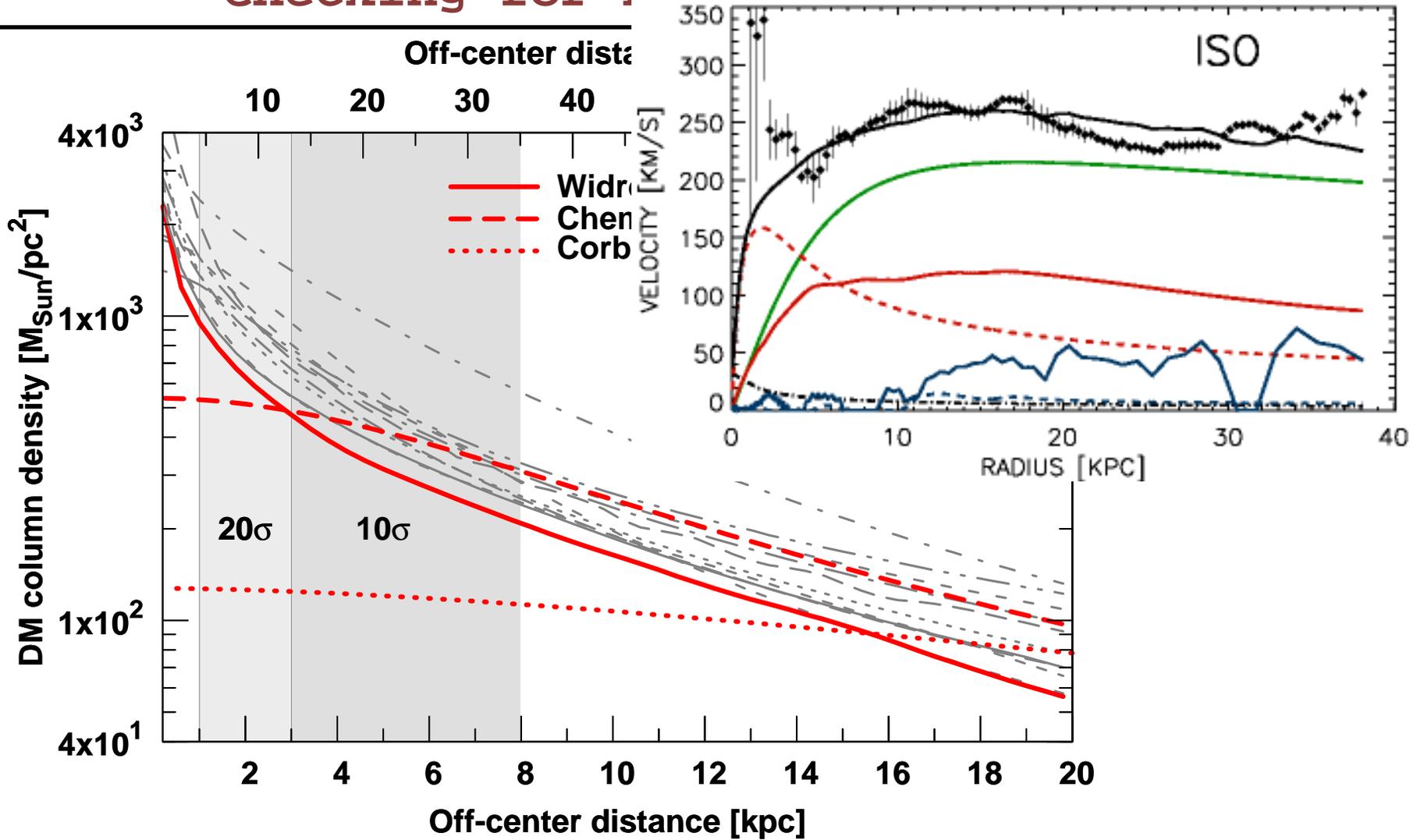
$$S_{Sc} = 140 M_{\odot} \text{pc}^{-2}$$

■ **Andromeda galaxy (M31) :**

$$S_{M31} \sim 100 - 600 M_{\odot} / \text{pc}^2$$

Do we see this 2.5 keV line?

# Checking for $\tilde{\nu}_m$ lines in M31



Willman 1 spectral feature excluded with high significance from archival observations of M31 and Fornax and Sculptor dSphs