

Natural and flavorful SUSY at the LHC

Andreas Weiler



Blois
29/5/12

w/ M. Papucci^{1,2}, J. Ruderman^{1,2} (LBL Berkely)
Gilad Perez², Rakhi Mahbubani² (CERN)

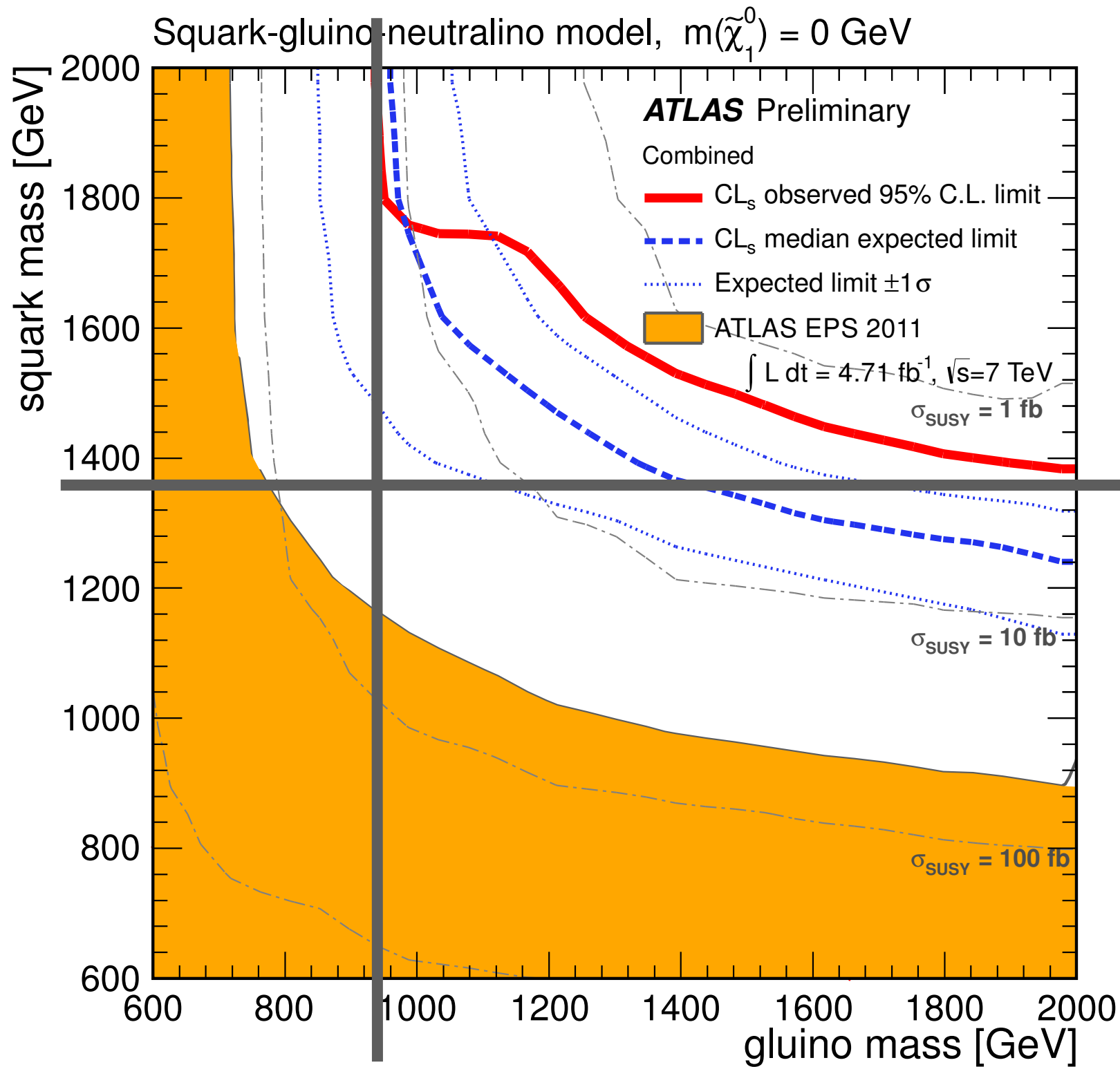
¹ | 10.6926 ² | 206.xxxx

Direct searches

→ Sunil Somalwar's talk

- Implications of a 125 GeV Higgs for the MSSM → ~~Nima's~~ & Csaba Csaki's talk

4.7 /fb Susy, post-Moriond



~930 GeV

~1400 GeV

Susy searches

What have we learned about the susy spectrum after 5 fb⁻¹ ?

- 1st & 2nd generation squarks need to be heavy $> 1.2-1.5 \text{ TeV}$ from jets+MET searches with 5/fb
- gluino limits above $\sim 900 \text{ GeV}$ (also from various other channels)

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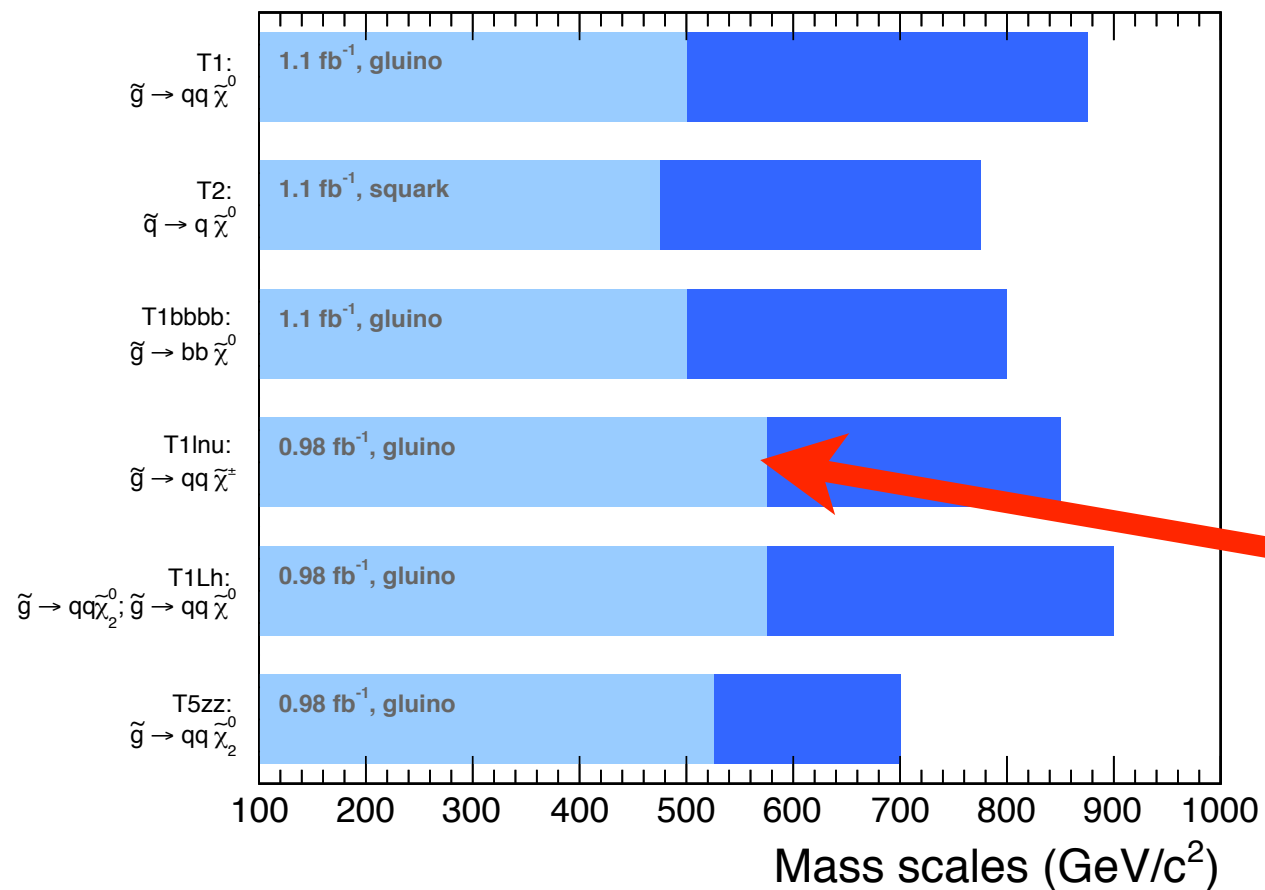
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Caveats?

Impact of LSP mass

Compressed spectrum alleviates bounds

Ranges of exclusion limits for gluinos and squarks, varying $m(\tilde{\chi}^0)$
CMS preliminary



$$m_{\text{gluino}} - m_{\text{LSP}} = 200 \text{ GeV}$$

For limits on $m(\tilde{g}), m(\tilde{q}) \gg m(\tilde{g})$ (and vice versa), $\sigma^{\text{prod}} = \sigma^{\text{NLO-QCD}}$.

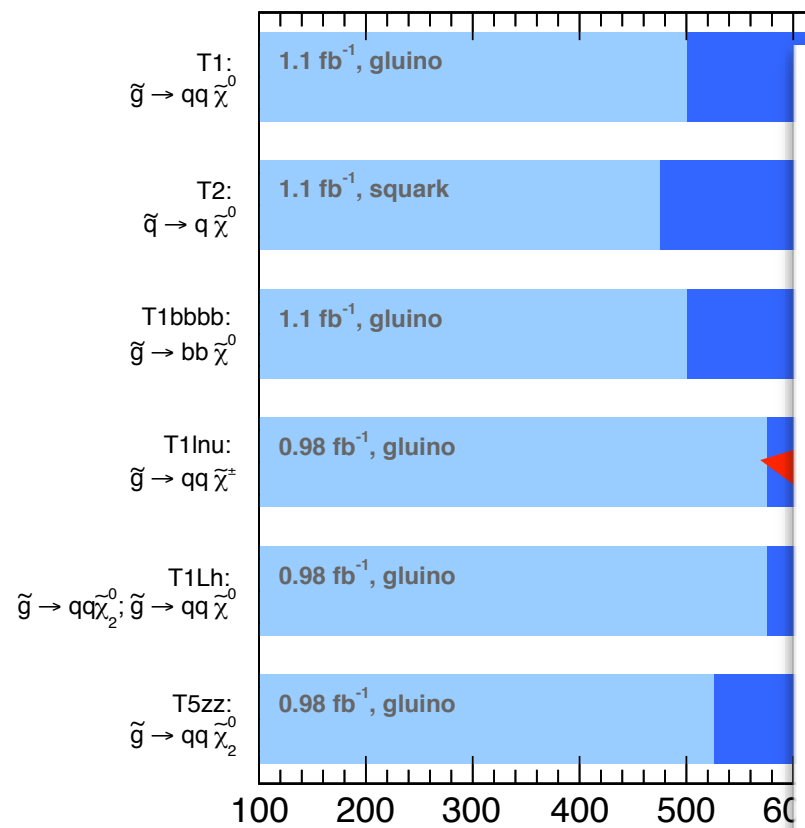
$$m(\tilde{\chi}^\pm), m(\tilde{\chi}_2^0) = \frac{m(\tilde{g}) + m(\tilde{\chi}^0)}{2}$$

$m(\tilde{\chi}^0)$ is varied from 0 GeV/c² (dark blue) to $m(\tilde{g}) - 200$ GeV/c² (light blue).

Impact of LSP mass

Compressed spectrum alleviates bounds

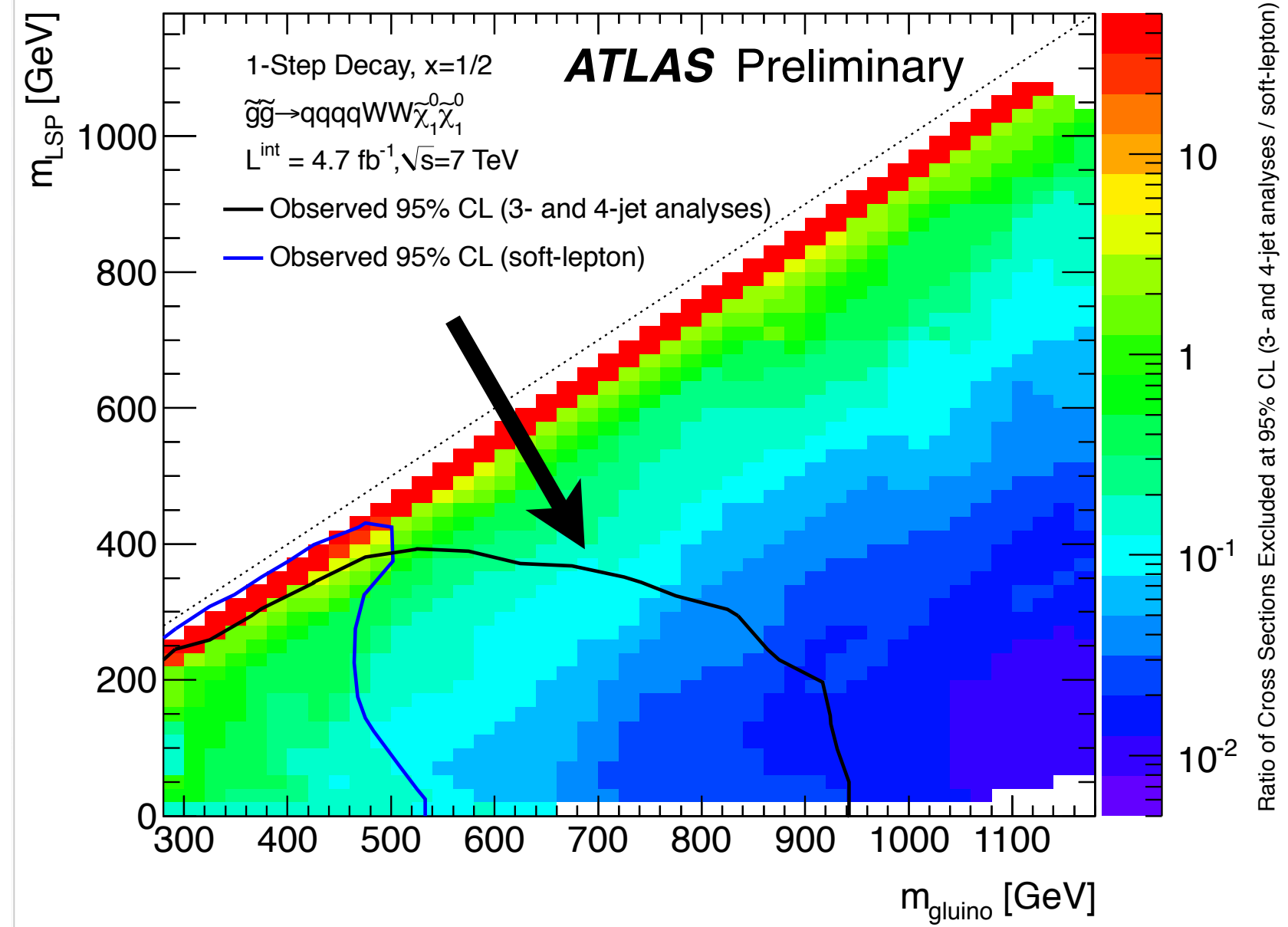
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For limits on $m(\tilde{g}), m(\tilde{q}) \gg m(\tilde{g})$ (and vice versa), σ^{prod}

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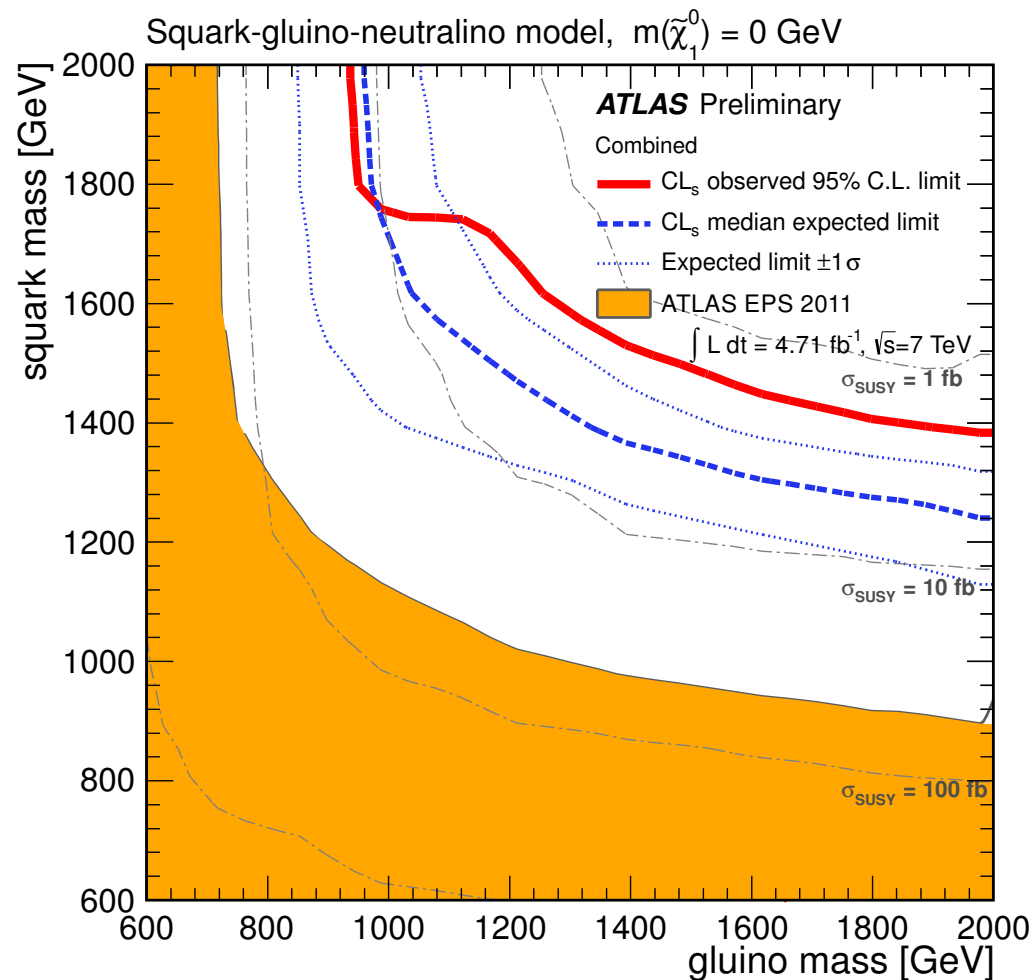
While limits are pushed we need to make sure that no stone is left unturned:

- Globally scan SUSY parameter space (e.g. pMSSM, e.g. Hewett et al., I205.5903)
- Modify CMSSM/mSUGRA, ... talk by H. Rzehak
- Focus the relevant kinematic features (this talk)
- Focus on theoretically-motivated models first (also this talk)

1st and 2nd generation squark limits

*work in progress with
Michele Papucci, Josh Ruderman,
Gilad Perez, Rakhi Mahbubani*

1st & 2nd generation squark limits



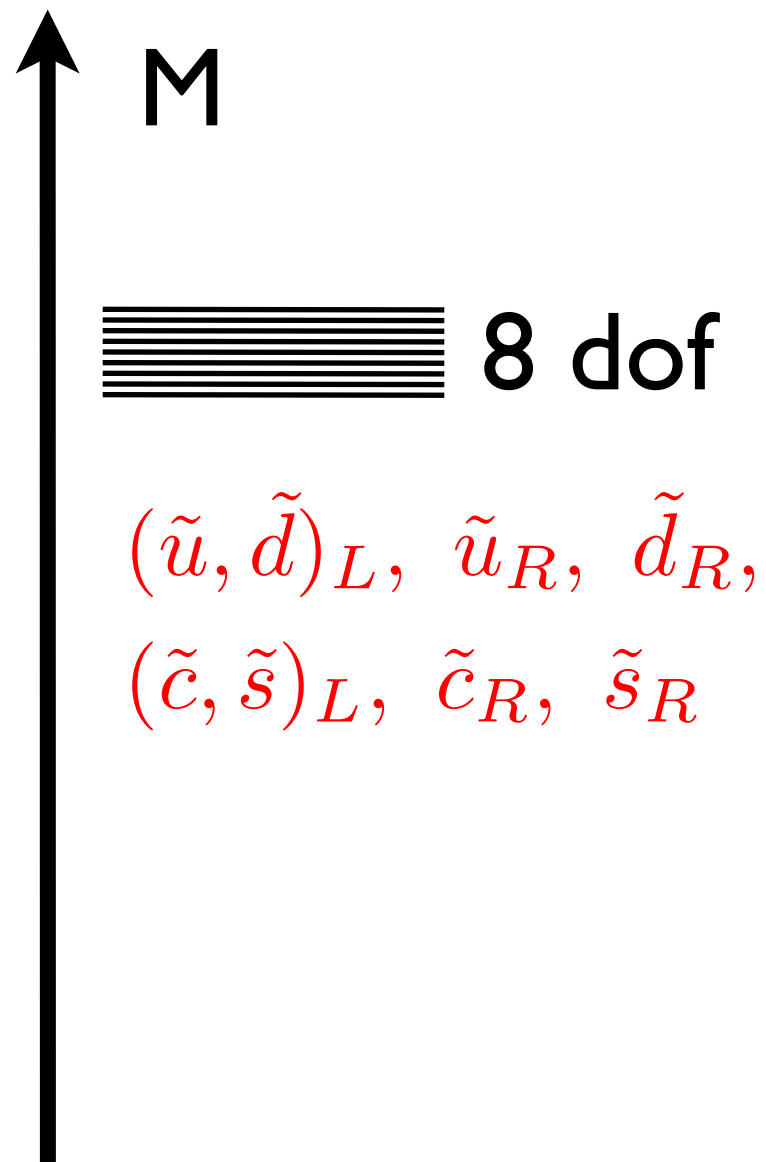
Light squarks **> 1.4 TeV?**

Assumptions? Necessary b/c of flavor?

What is driving the limit?

Holes in the net?

Do 1st & 2nd gen' squarks have to be degenerate?



- Because of flavor constraints?
Not really.

spectrum in ATLAS/CMS plots

→ Gino Isidori's talk

UTfit 08, Isidori, Perez, Nir '10

Operator	Bounds on Λ in TeV ($c_{ij} = 1$)		Bounds on c_{ij} ($\Lambda = 1$ TeV)		Observables
	Re	Im	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^4	9.0×10^{-7}	3.4×10^{-9}	$\Delta m_K; \epsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	1.8×10^4	3.2×10^5	6.9×10^{-9}	2.6×10^{-11}	$\Delta m_K; \epsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^3	2.9×10^3	5.6×10^{-7}	1.0×10^{-7}	$\Delta m_D; q/p , \phi_D$
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Very strong suppression! New flavor violation must either **approximately (exactly?) follow SM structure...**

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Very strong suppression! New flavor violation must either **approximately (exactly?) follow SM structure...**

... or exist only at **very high scales ($10^2 - 10^5$ TeV)**

SUSY & Flavor

Flavor Bounds (K, D, B, Bs mixing, ...) controlled by

$$(\delta_{ij}^q)_{MM} = \frac{1}{\tilde{m}_q^2} \sum_{\alpha} (K_M^q)_{i\alpha} (K_M^q)_{j\alpha}^* \Delta \tilde{m}_{q\alpha}^2$$

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mixing matrices

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mixing matrices mass splitting

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mixing matrices

mass splitting

(m=1TeV)

q	ij	$(\delta_{ij}^q)_{MM}$	$\langle \delta_{ij}^q \rangle$
d	12	0.03	0.002
d	13	0.2	0.07
d	23	0.6	0.2
u	12	0.1	0.008

Isidori et. al '10

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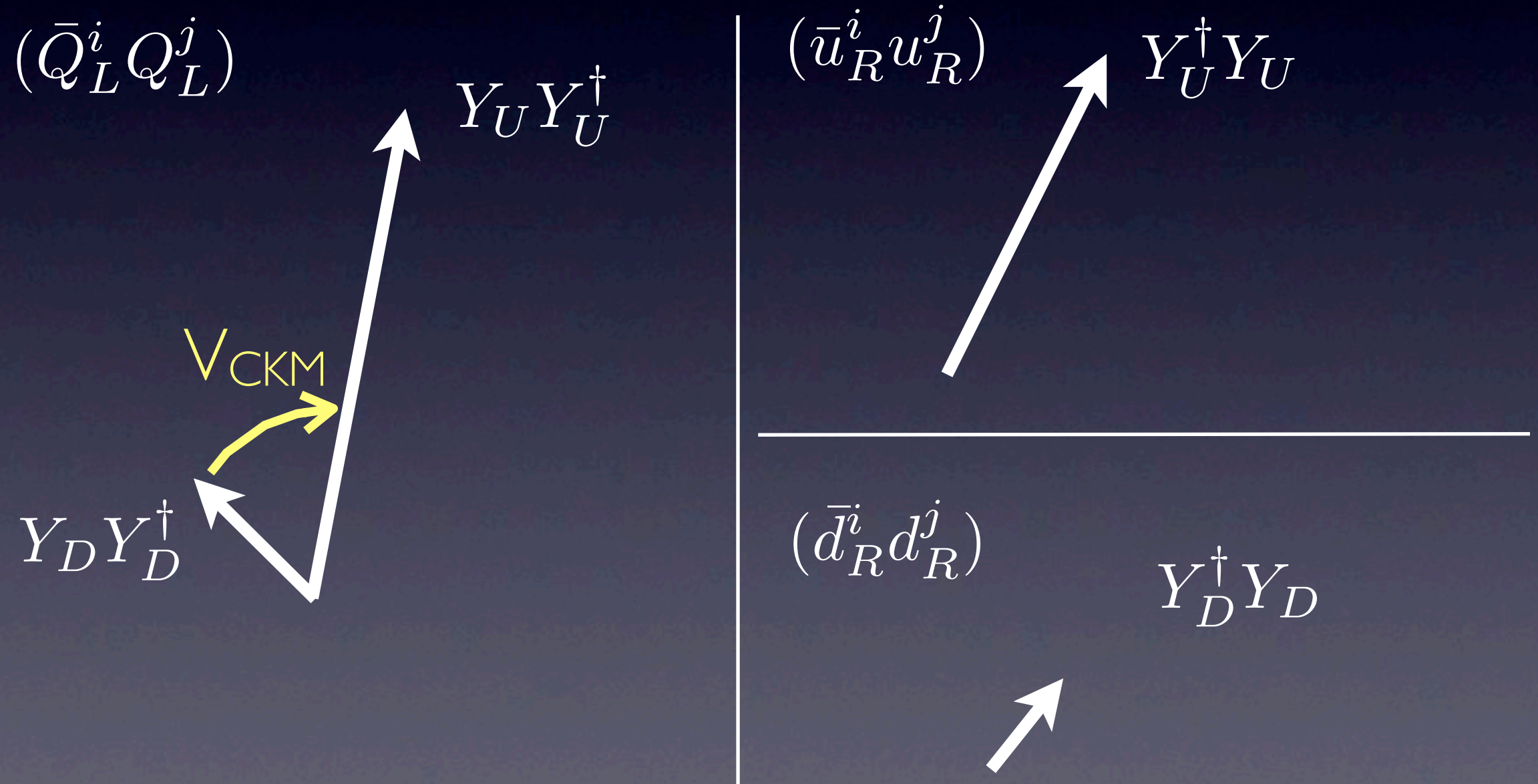
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large mixing
means splitting
must be $\ll 1$

A picture of flavor

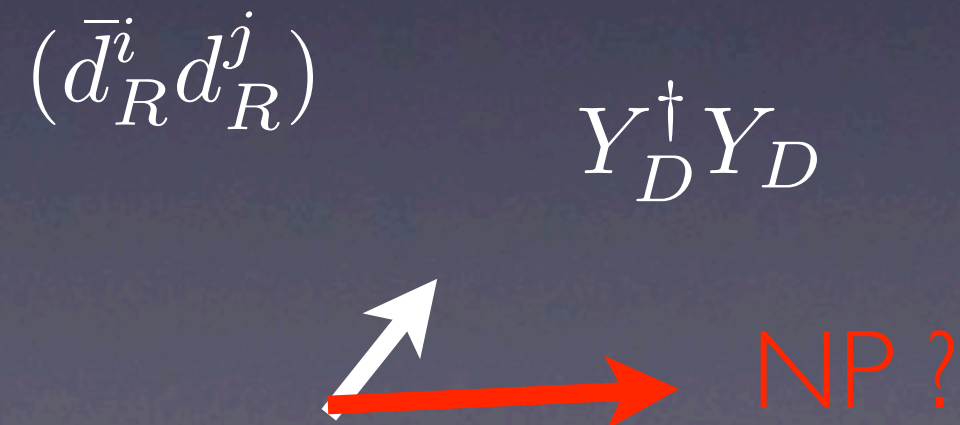
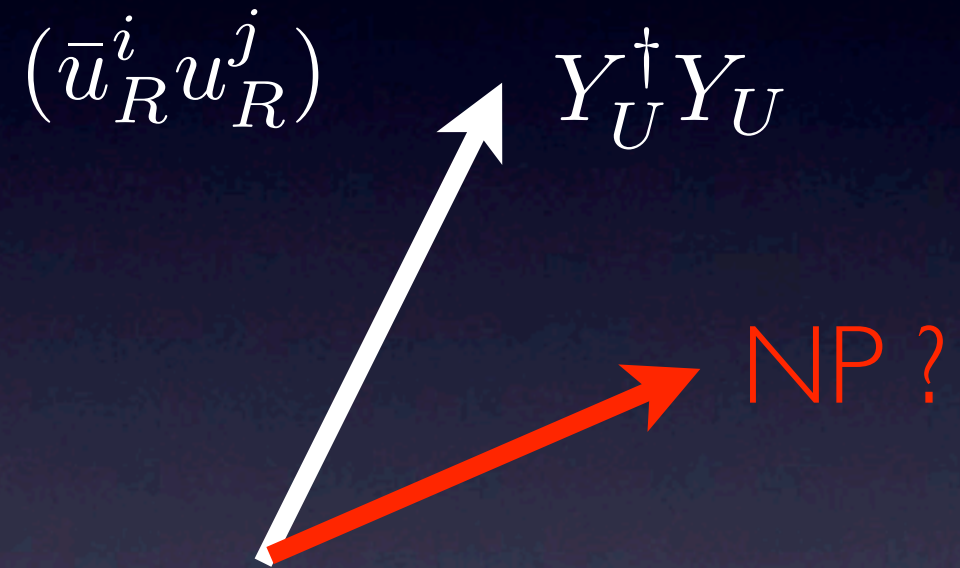
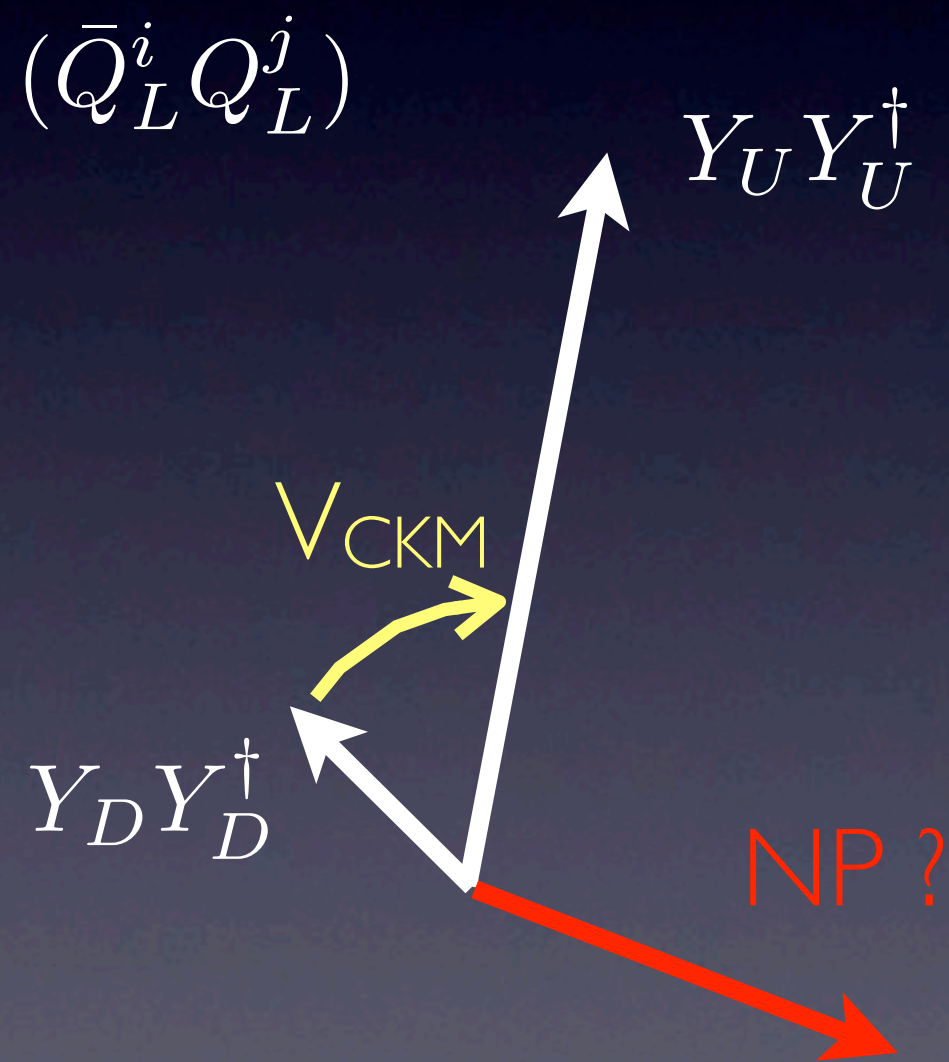
Yukawa matrices Y_U & Y_D encode flavor violation



+ LR, RL

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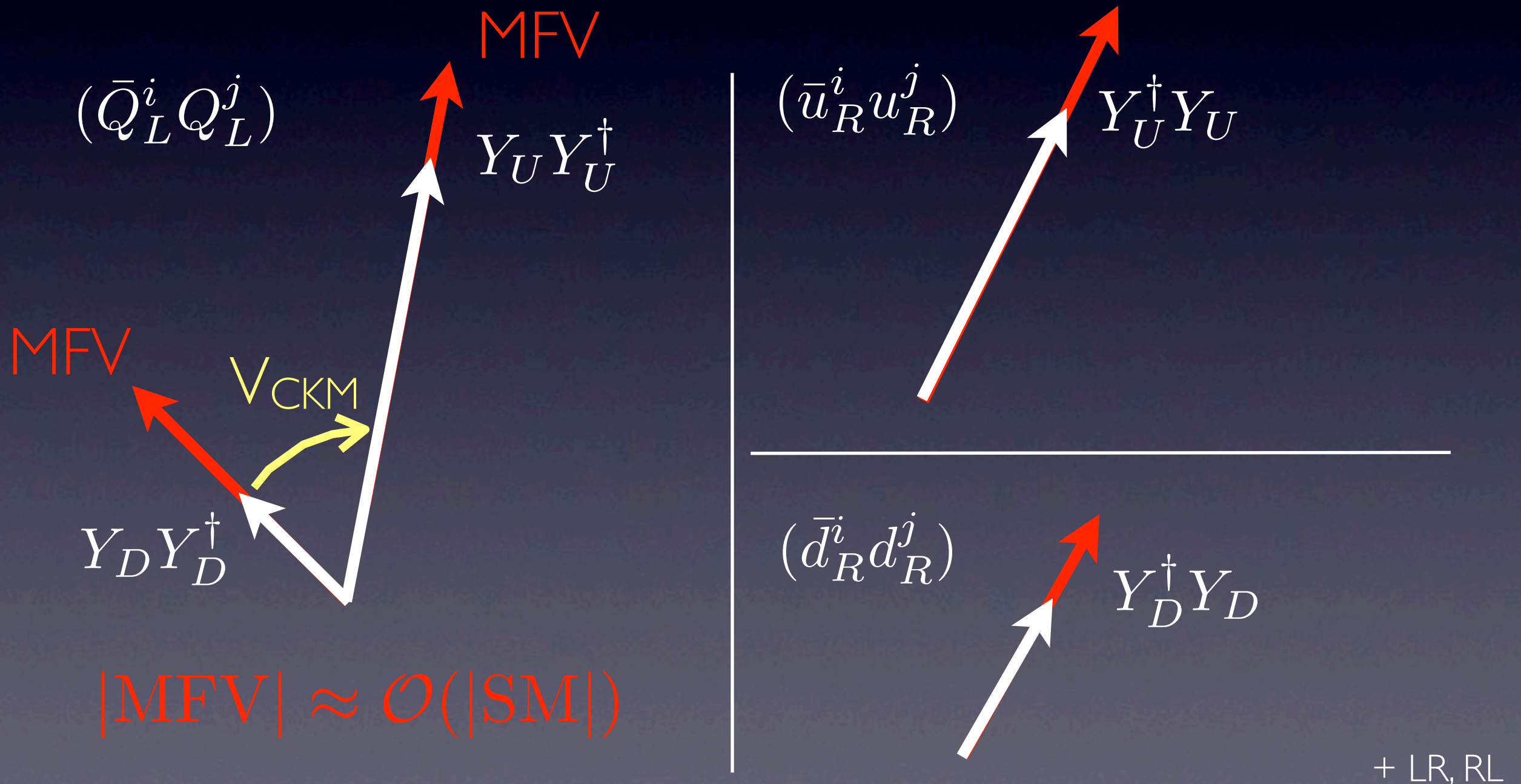
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Minimal flavor violation

Chivukula Georgi; Buras et. al; D'Ambrosio et. al

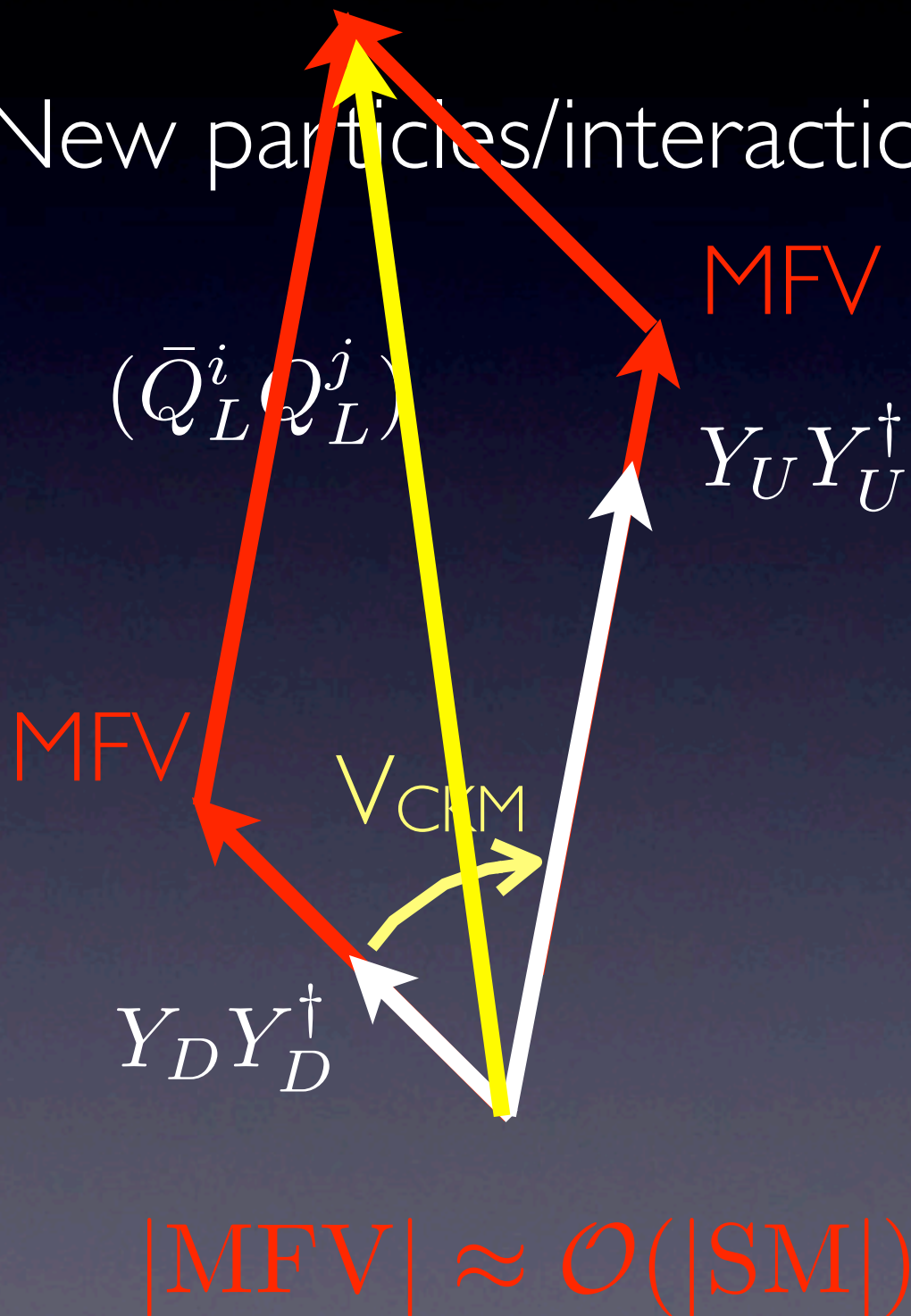
New particles/interactions, but flavor structure $\sim V_{\text{CKM}}$



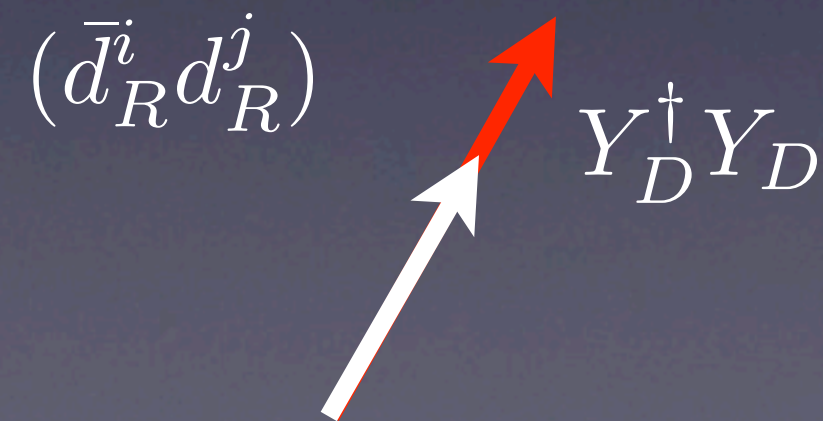
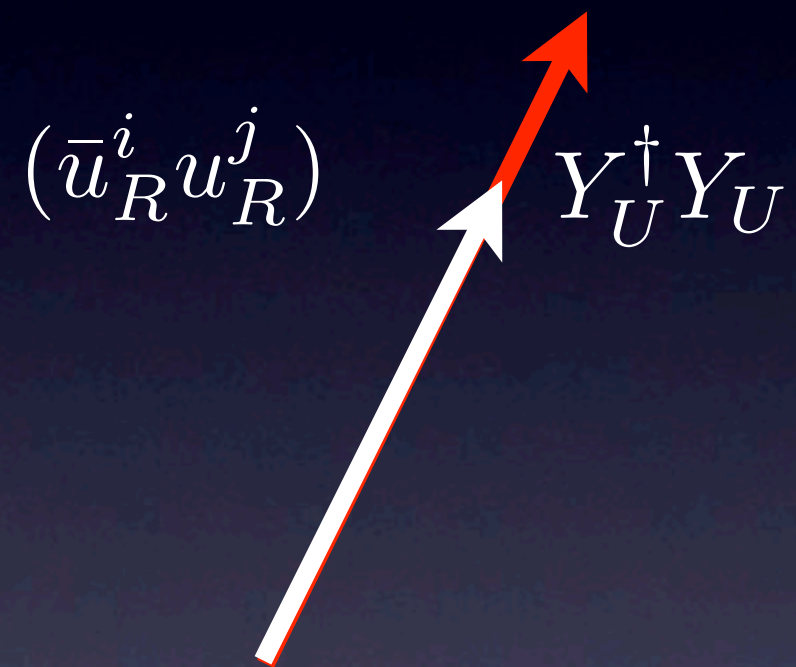
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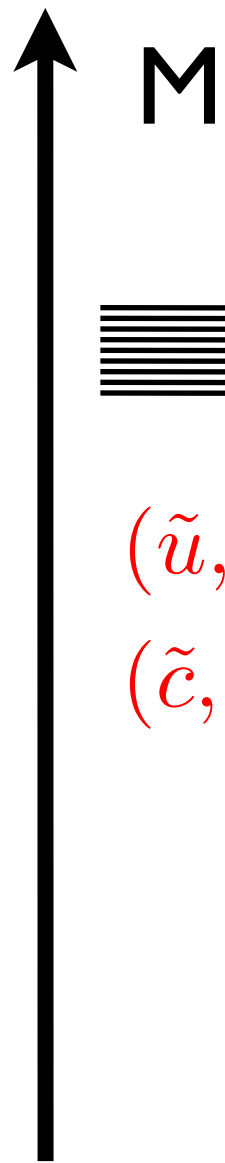
$$|MFV| \approx \mathcal{O}(|SM|)$$



+ LR, RL

Minimal Flavor Violation

- **Trivial:** Squark masses same for all three generations but split between $\tilde{Q}_L, \tilde{u}_R, \tilde{d}_R$
- Split among generations but split like in SM:
mass-differences $\propto Y_{U,D} \sim (0,0,1)$



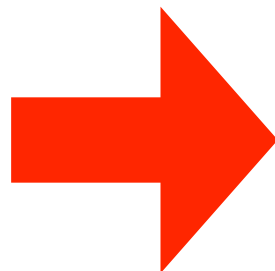
M



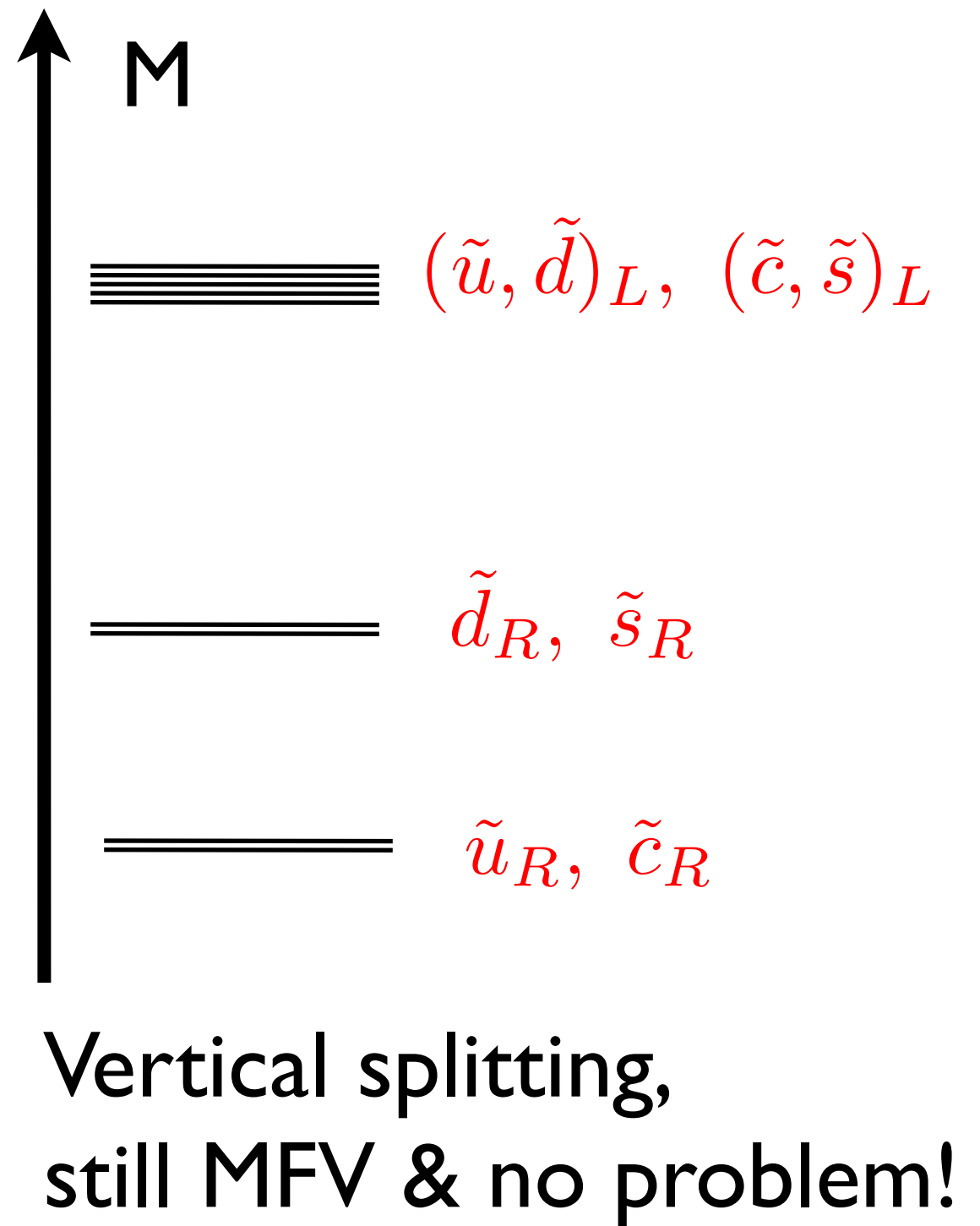
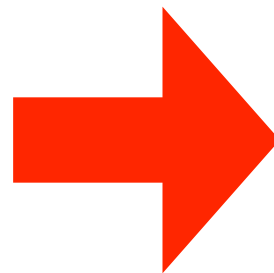
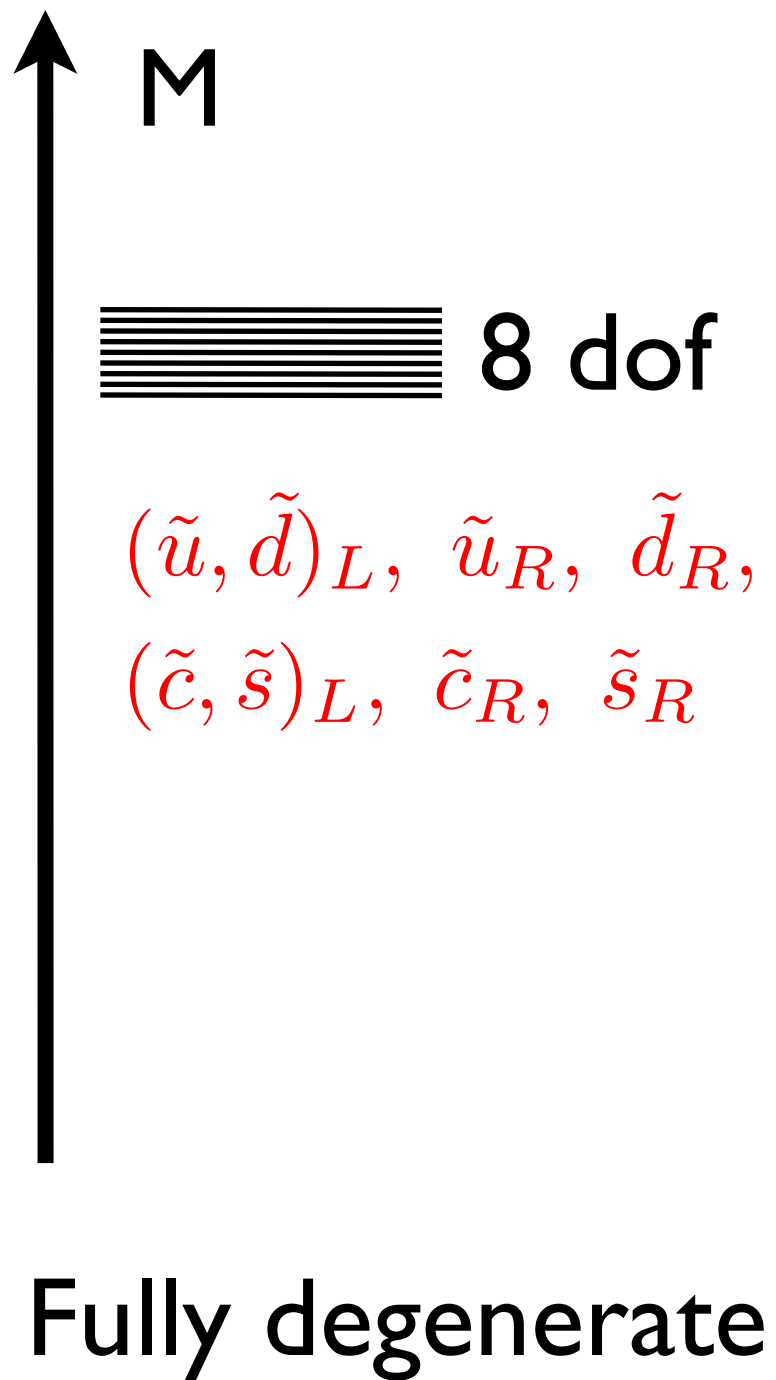
8 dof

$(\tilde{u}, \tilde{d})_L, \tilde{u}_R, \tilde{d}_R,$

$(\tilde{c}, \tilde{s})_L, \tilde{c}_R, \tilde{s}_R$

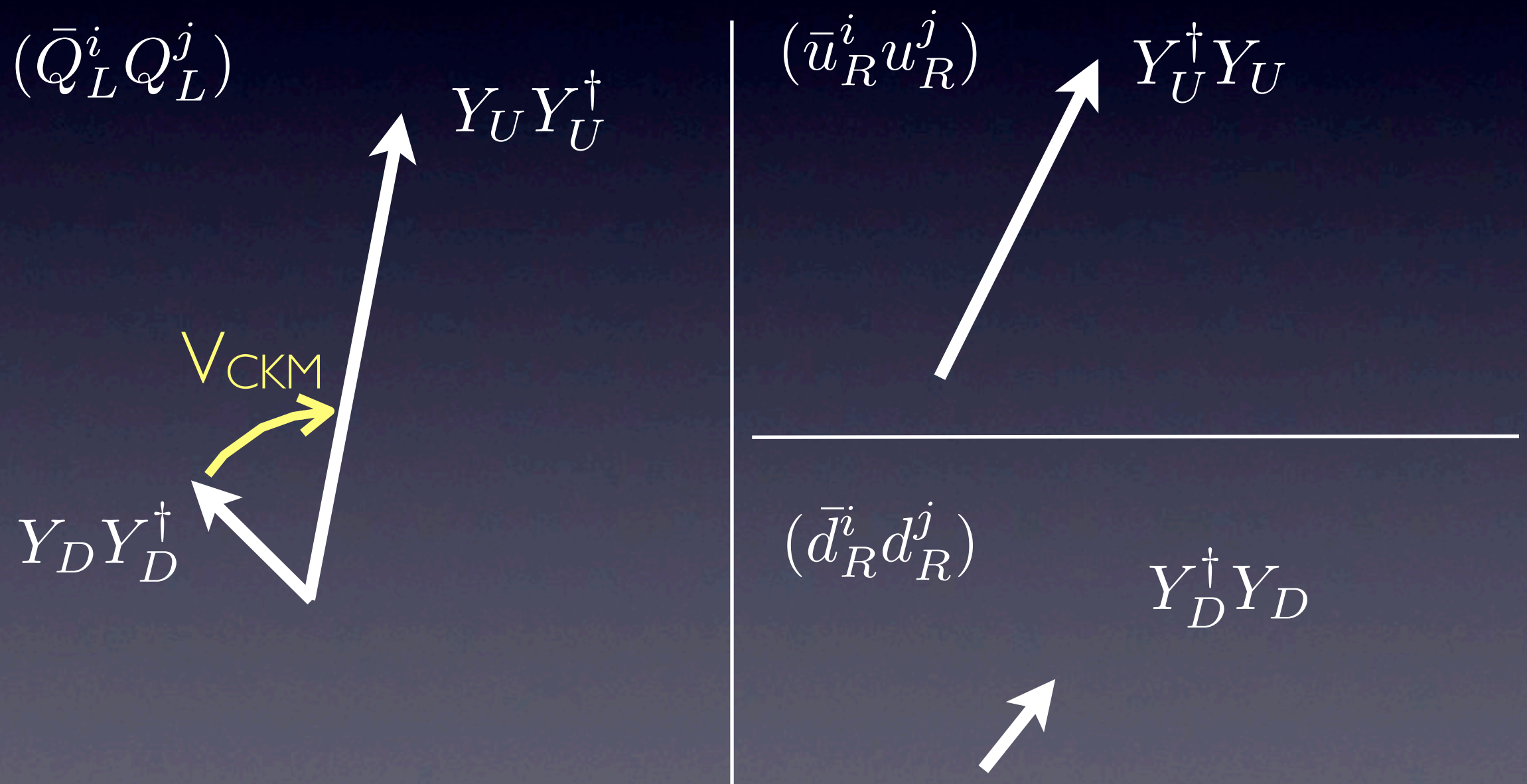


Fully degenerate



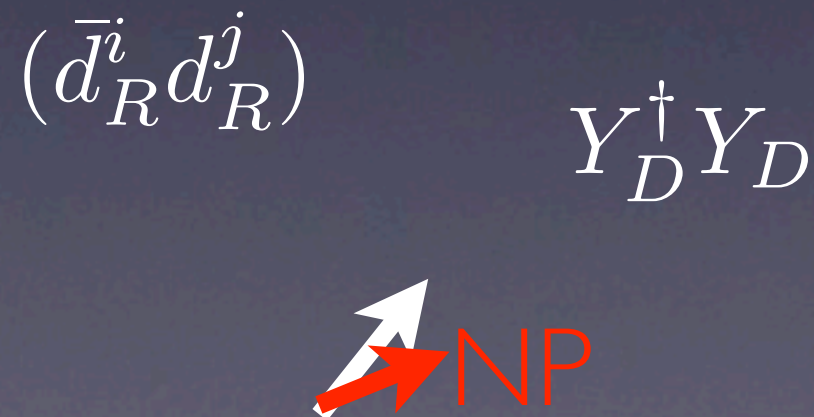
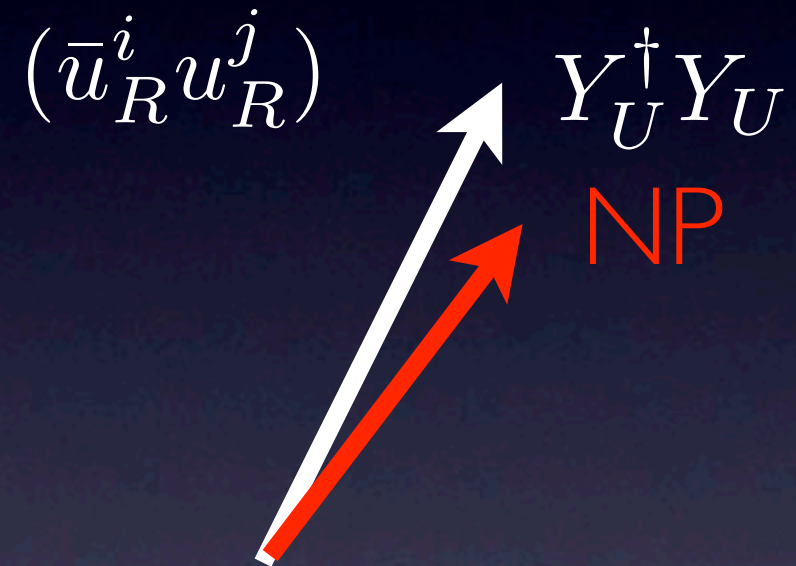
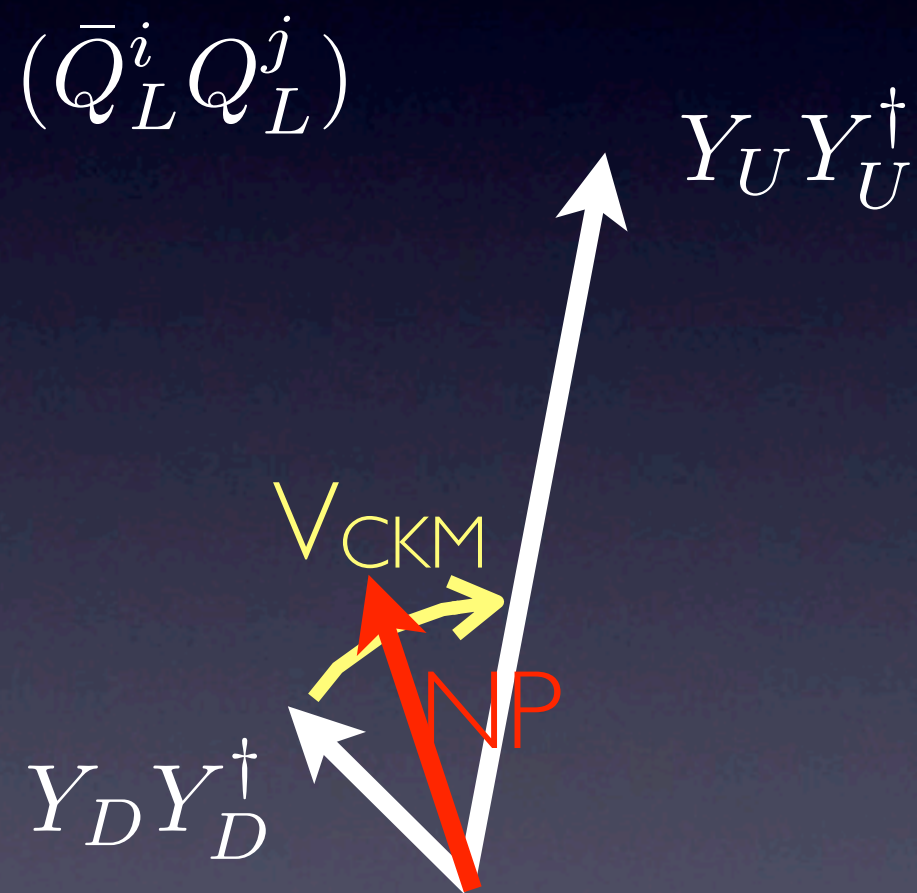
Flavor dynamics: alignment

Dynamics (e.g. $U(1)_{\text{horiz.}}$) generates hierarchies in masses & mixings. Consequence: **partial alignment** with SM

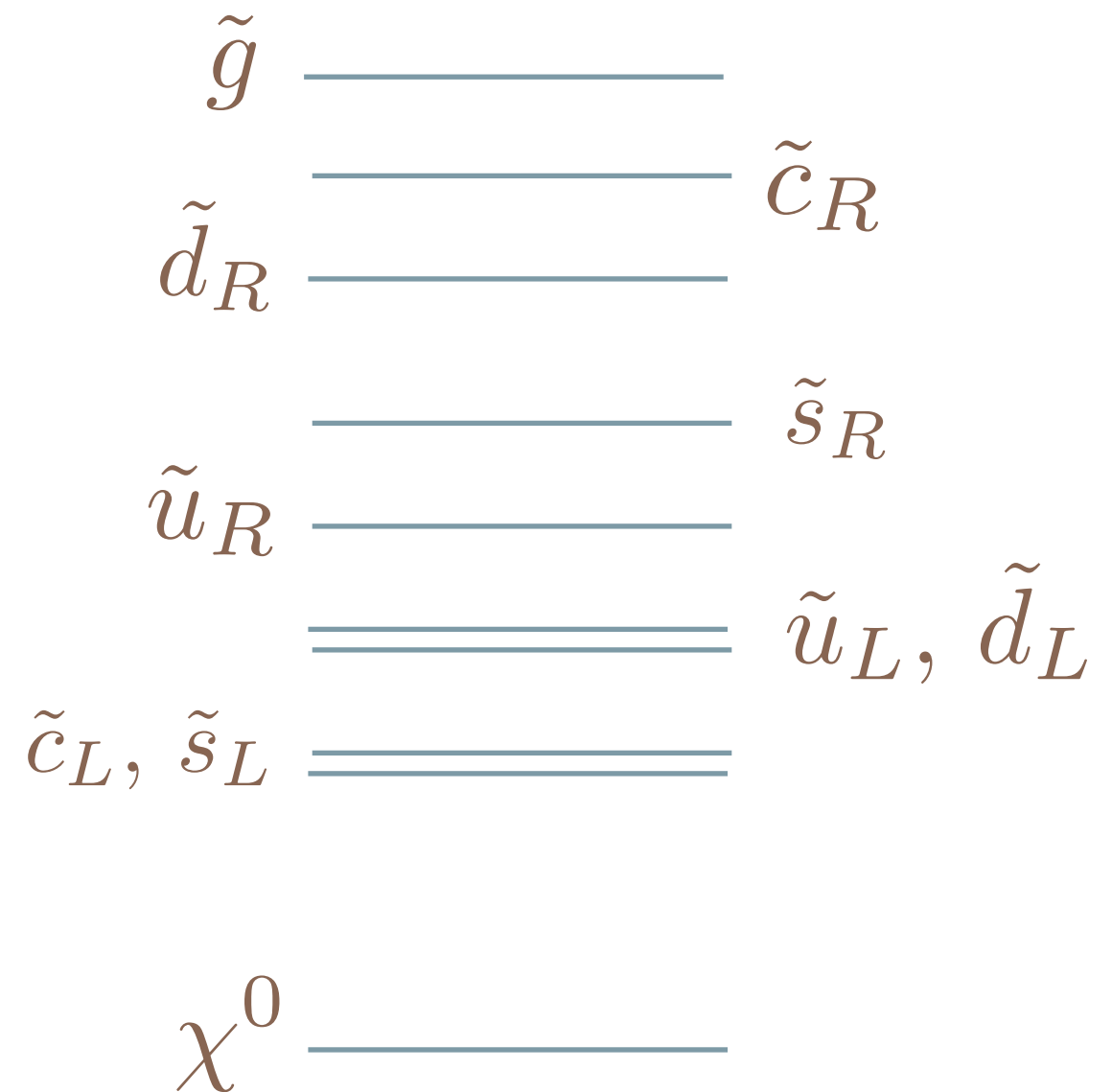


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- Right handed squarks can be strongly aligned $\rightarrow O(1)$ splitting possible
- Left handed squarks can be aligned **either to up or down** quarks
- Existence of V_{CKM} \rightarrow Left handed squarks contribute to K or D mixing
- Left handed squark splitting mildly constrained (CP phases generically small)



Alignment

$$(\delta_{ij}^q)_{MM} = \frac{1}{\tilde{m}_q^2} \sum_{\alpha} (K_M^q)_{i\alpha} (K_M^q)_{j\alpha}^* \Delta \tilde{m}_{q\alpha}^2$$

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mixing / misalignment between
SM Yukawas and squark mass matrices

Alignment

Seiberg & Nir

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If by symmetry: $K_{ij} \sim$ diagonal \Rightarrow $O(1)$ mass splitting
allowed!

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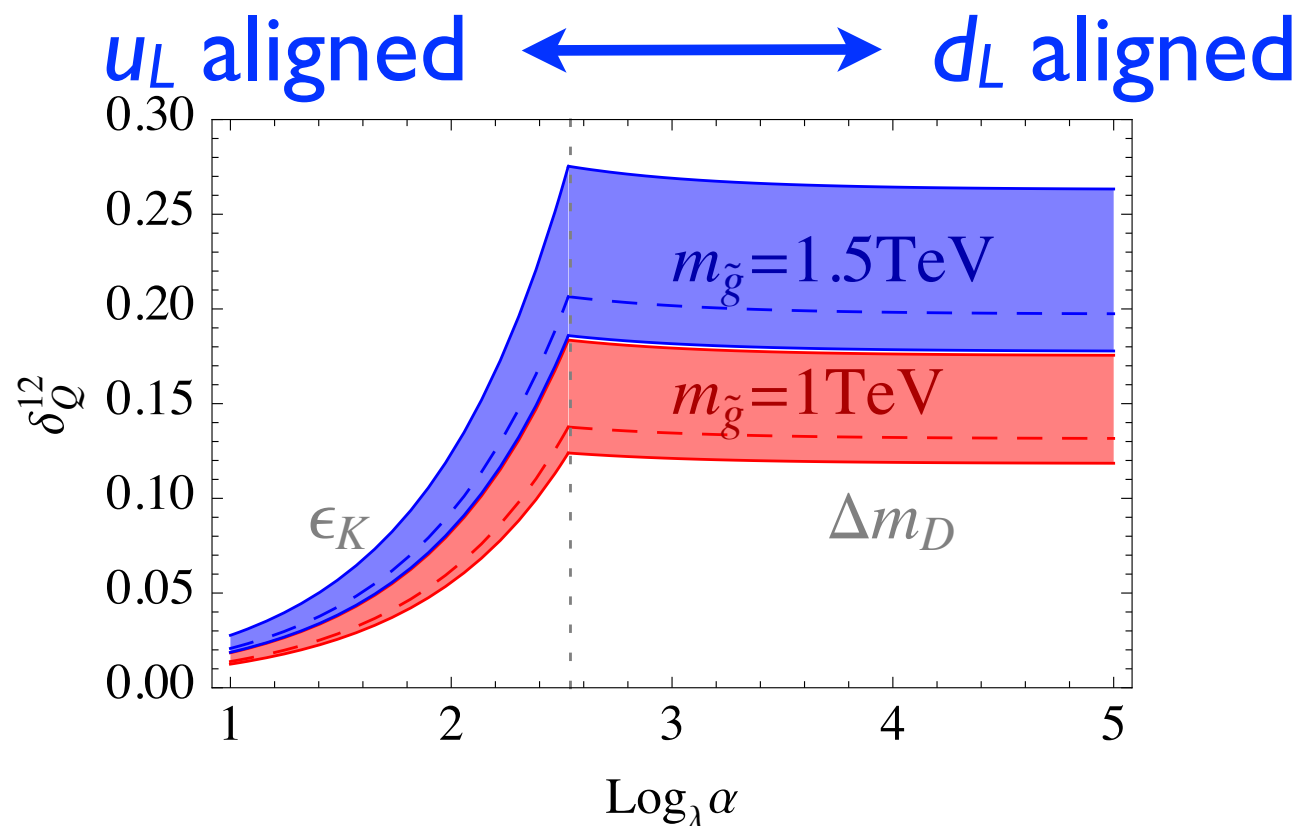
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Gedalia et. al

squark splitting



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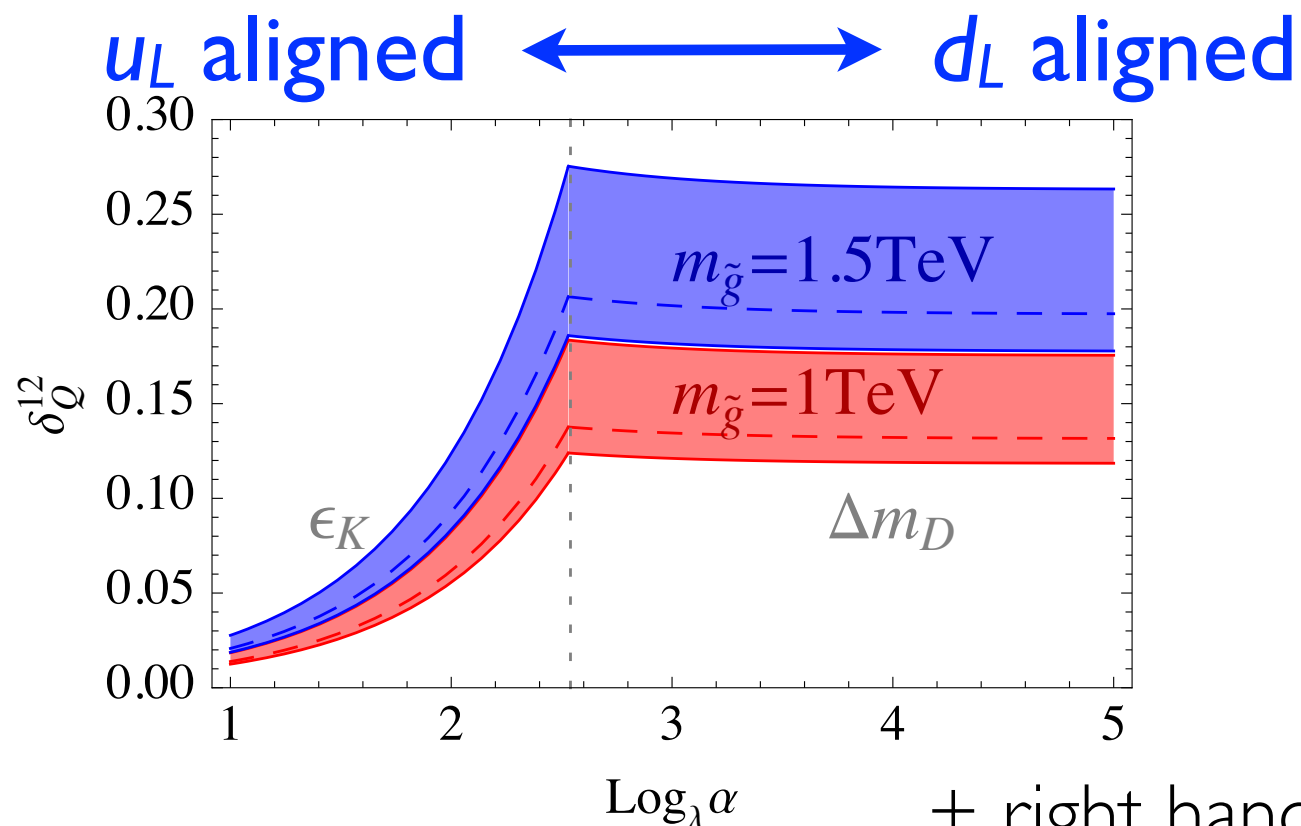
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Gedalia et. al

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Example:

$$m_{\text{gluino}} = 1.3 \text{ TeV}$$

$$m_{Q1} = 550 \text{ GeV}$$

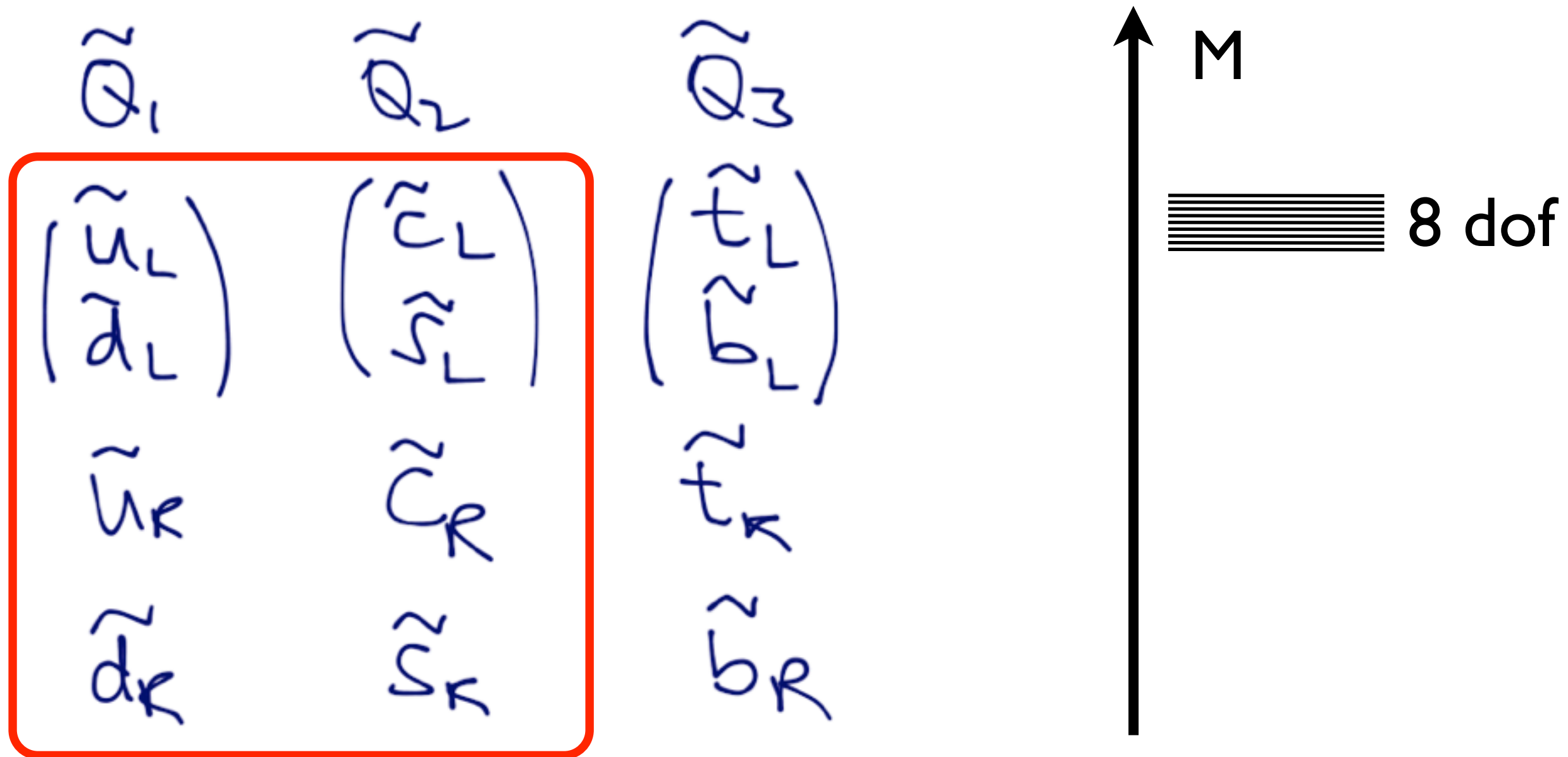
$$m_{Q2} = 950 \text{ GeV}$$

+ right handed squarks split by arbitrary amount

Flavor vs. squark masses: summary

- Generic 1-2 splitting has to be small, **BUT:**
- Can split **vertically**: split irrep's, MFV
- Can split **horizontally**, if squark flavor aligned

1st & 2nd gen' squarks degenerate



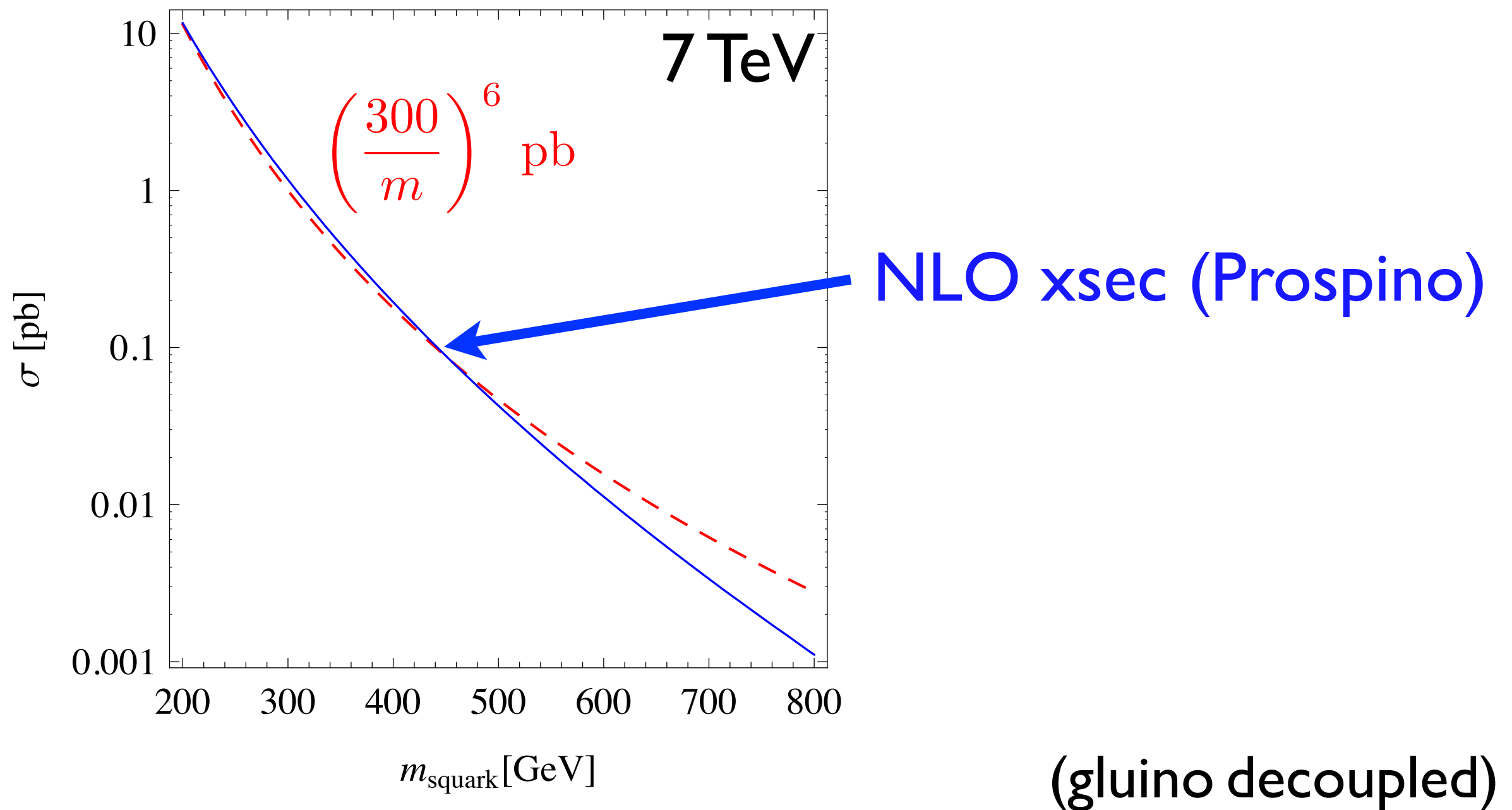
assumption in ATLAS/CMS plots

Does it matter if
we relax the degeneracy
assumption?

Naive answer: **not so much.**

Cross-sections vs. mass

$$\sigma(pp \rightarrow \tilde{u}_R \tilde{u}_R^*) \propto \frac{1}{m^6} \quad (\text{roughly})$$



Back of the envelope estimate



Cross-sections roughly scale like $\sim 1/m^6$.

Example: 8 light squarks \rightarrow 2 light squarks

Shift limit only by $\sim 4^{1/6} - 1 \approx 25\%$

\rightarrow **too naive!**

Dedicated study needed

- Production cross-section can be **flavor dependent** if gluino is not fully decoupled through p.d.f's (*u vs. d, sea vs. valence*)
- Experimental **efficiencies** for light squarks are not very good : efficiencies have hard thresholds and current limits are on the thresholds

How can we extract limits on non-degenerate 1st and 2nd gen' squarks from experimental searches?

“The experiments haven’t covered my favorite model”

Relax & Wait?



vs.

* not his real attitude.

“The experiments haven’t covered my favorite model”

Relax & Wait?



vs.



Let’s check!

* not his real attitude.

DYI limits

CERN-PH-EP-2011-145

Search for squarks and gluinos using final states with jets and missing transverse momentum with the ATLAS detector in $\sqrt{s} = 7$ TeV proton-proton collisions

The ATLAS Collaboration

Example:
jets+ MET 1.041/fb

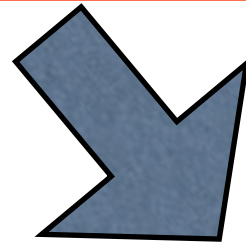
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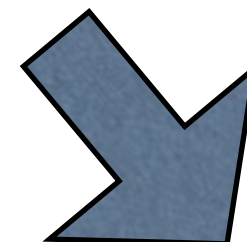
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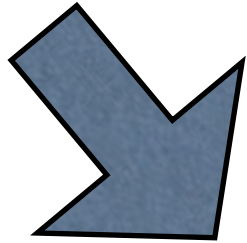


Signal Region	≥ 2 -jet	≥ 3 -jet	≥ 4 -jet	High mass
E_T^{miss}	> 130	> 130	> 130	> 130
Leading jet p_T	> 130	> 130	> 130	> 130
Second jet p_T	> 40	> 40	> 40	> 80
Third jet p_T	–	> 40	> 40	> 80
Fourth jet p_T	–	–	> 40	> 80
$\Delta\phi(\text{jet}, \vec{P}_T^{\text{miss}})_{\text{min}}$	> 0.4	> 0.4	> 0.4	> 0.4
$E_T^{\text{miss}}/m_{\text{eff}}$	> 0.3	> 0.25	> 0.25	> 0.2
m_{eff}	> 1000	> 1000	$> 500/1000$	> 1100

signal bins



Bgd's are left to the experimentalists...
 stay out of **control regions!**



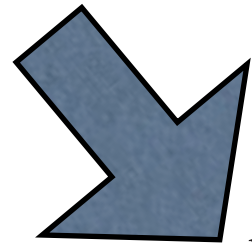
Process	Signal Region				
	$\geq 2\text{-jet}$	$\geq 3\text{-jet}$	$\geq 4\text{-jet},$ $m_{\text{eff}} > 500 \text{ GeV}$	$\geq 4\text{-jet},$ $m_{\text{eff}} > 1000 \text{ GeV}$	High mass
$Z/\gamma\text{+jets}$	$32.3 \pm 2.6 \pm 6.9$	$25.5 \pm 2.6 \pm 4.9$	$209 \pm 9 \pm 38$	$16.2 \pm 2.2 \pm 3.7$	$3.3 \pm 1.0 \pm 1.3$
$W\text{+jets}$	$26.4 \pm 4.0 \pm 6.7$	$22.6 \pm 3.5 \pm 5.6$	$349 \pm 30 \pm 122$	$13.0 \pm 2.2 \pm 4.7$	$2.1 \pm 0.8 \pm 1.1$
$t\bar{t}\text{+ single top}$	$3.4 \pm 1.6 \pm 1.6$	$5.9 \pm 2.0 \pm 2.2$	$425 \pm 39 \pm 84$	$4.0 \pm 1.3 \pm 2.0$	$5.7 \pm 1.8 \pm 1.9$
QCD multi-jet	$0.22 \pm 0.06 \pm 0.24$	$0.92 \pm 0.12 \pm 0.46$	$34 \pm 2 \pm 29$	$0.73 \pm 0.14 \pm 0.50$	$2.10 \pm 0.37 \pm 0.82$
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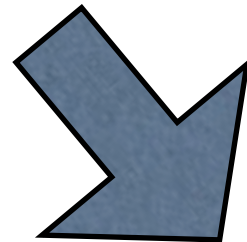
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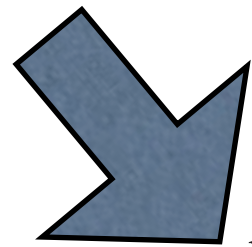


upper
 bound on
 signal xsec

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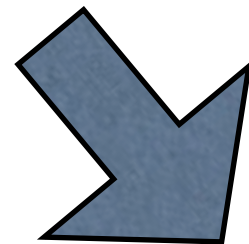
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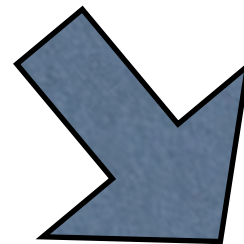


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upper
 bound on
 signal xsec



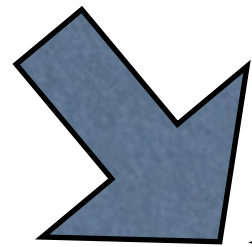
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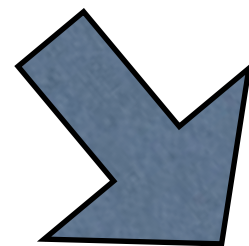
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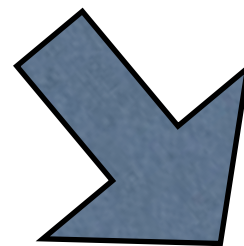


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upper
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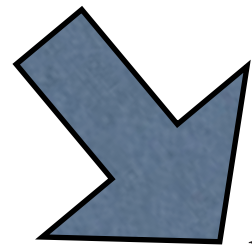
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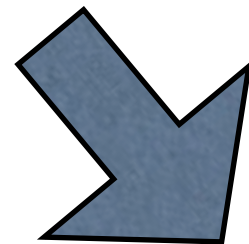
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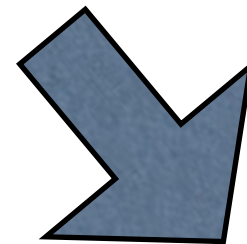


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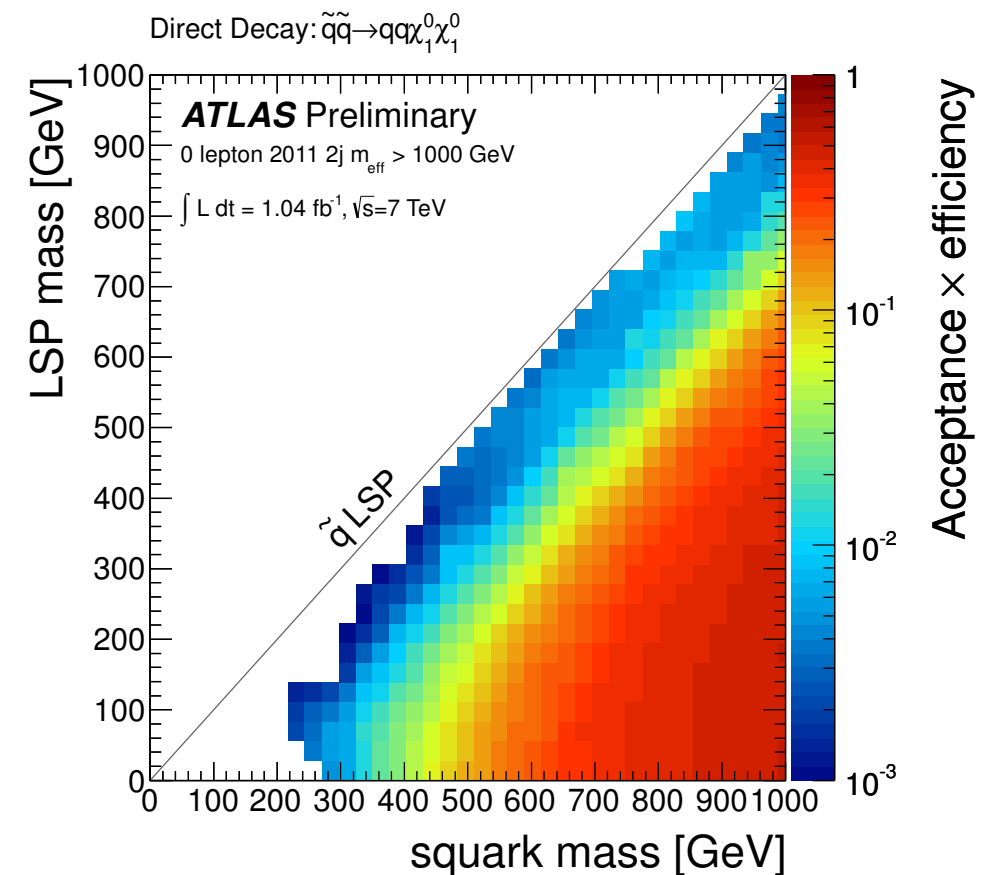
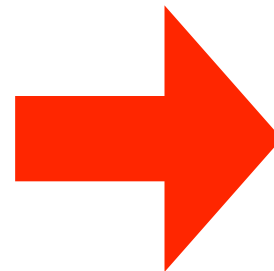
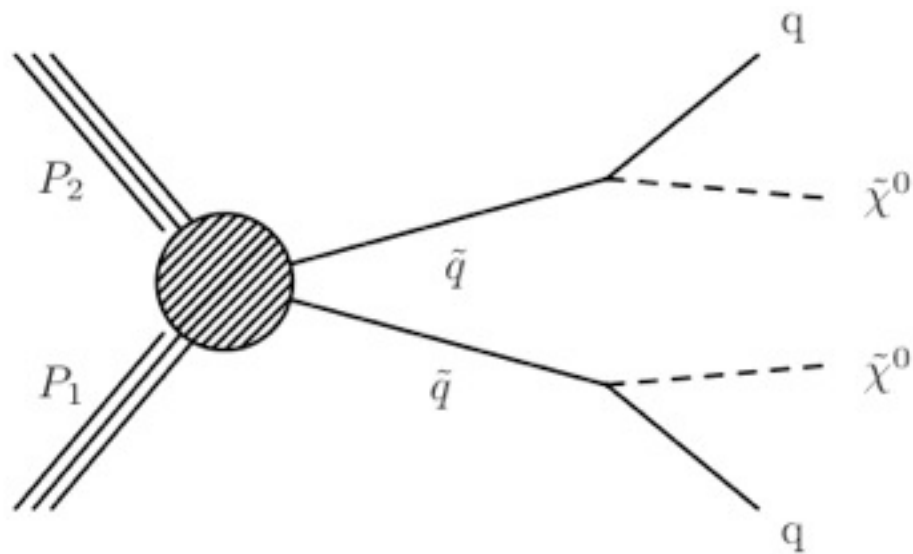


LIMIT!



Simplified Models to the rescue!

- Luckily ATLAS and CMS provide **efficiencies** for **simplified models** (so far only for 1/fb)

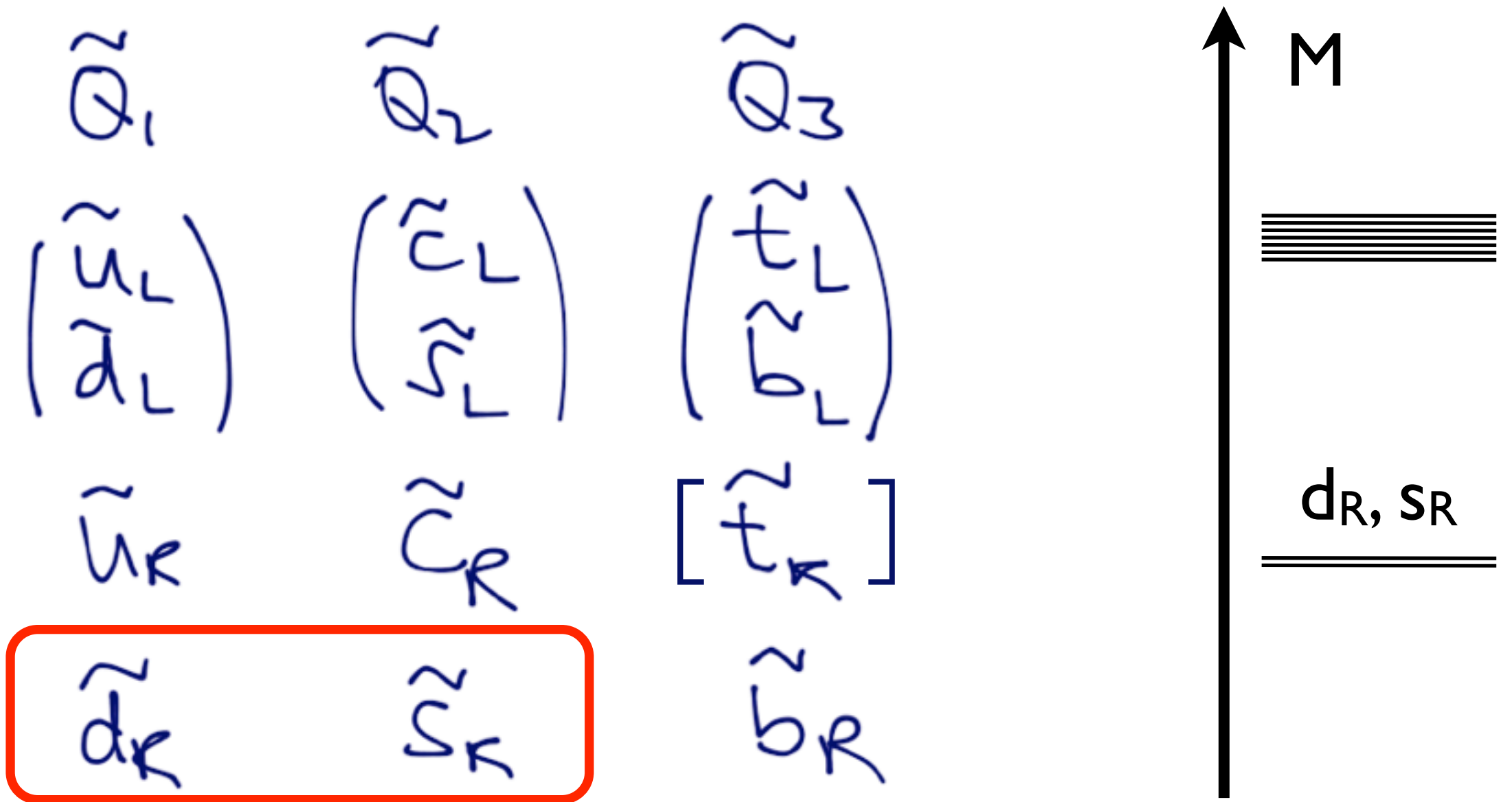


simplified topologies

e.g. [CMSPublic/PhysicsResultsSUSI 1004](https://arxiv.org/abs/1004.1004) or [ATLAS-CONF-2011-155/](https://arxiv.org/abs/1011.1551)

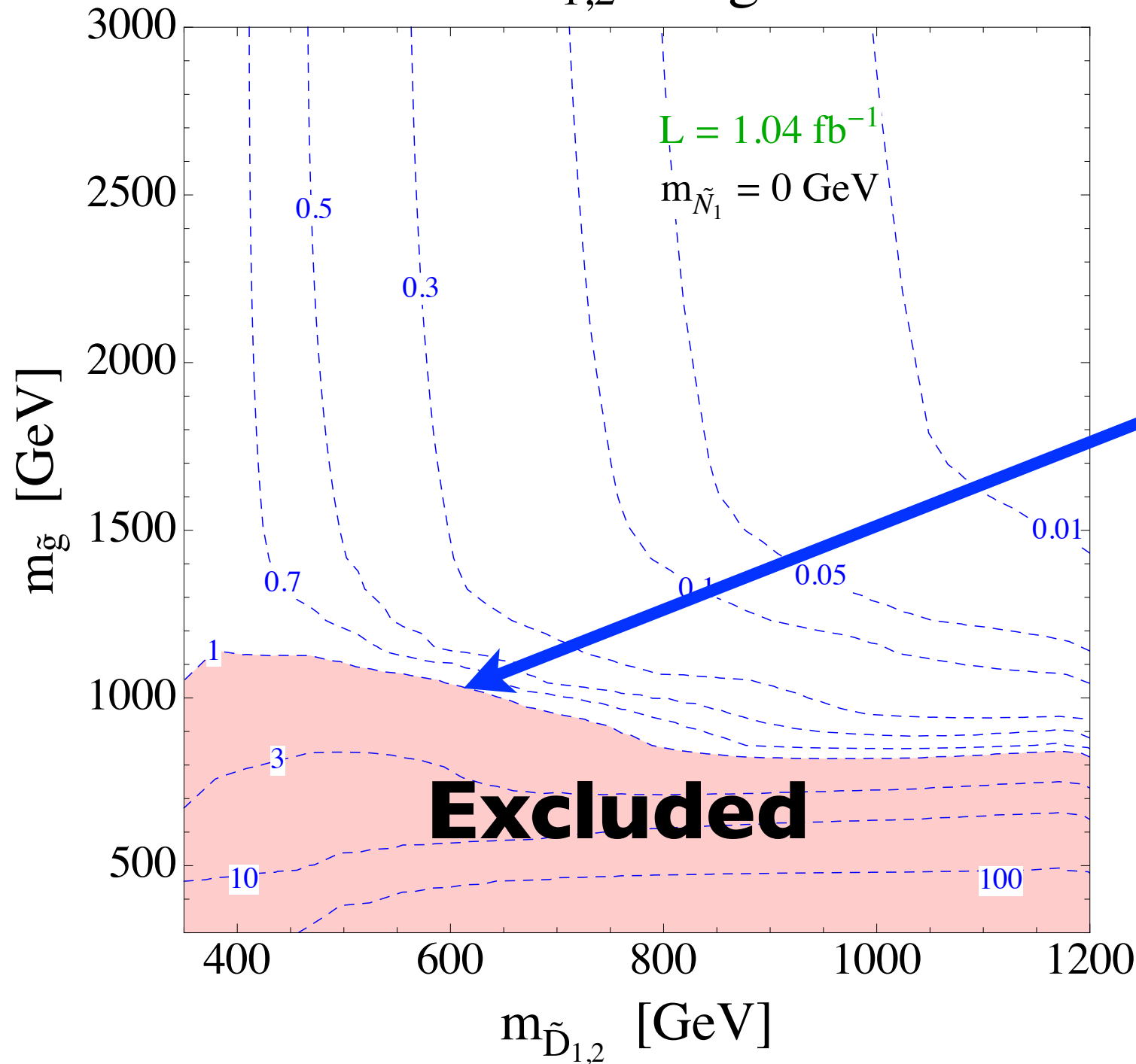
Thanks for providing root files, HEPDATA,...

Split squarks w/ different quantum numbers (vertical)



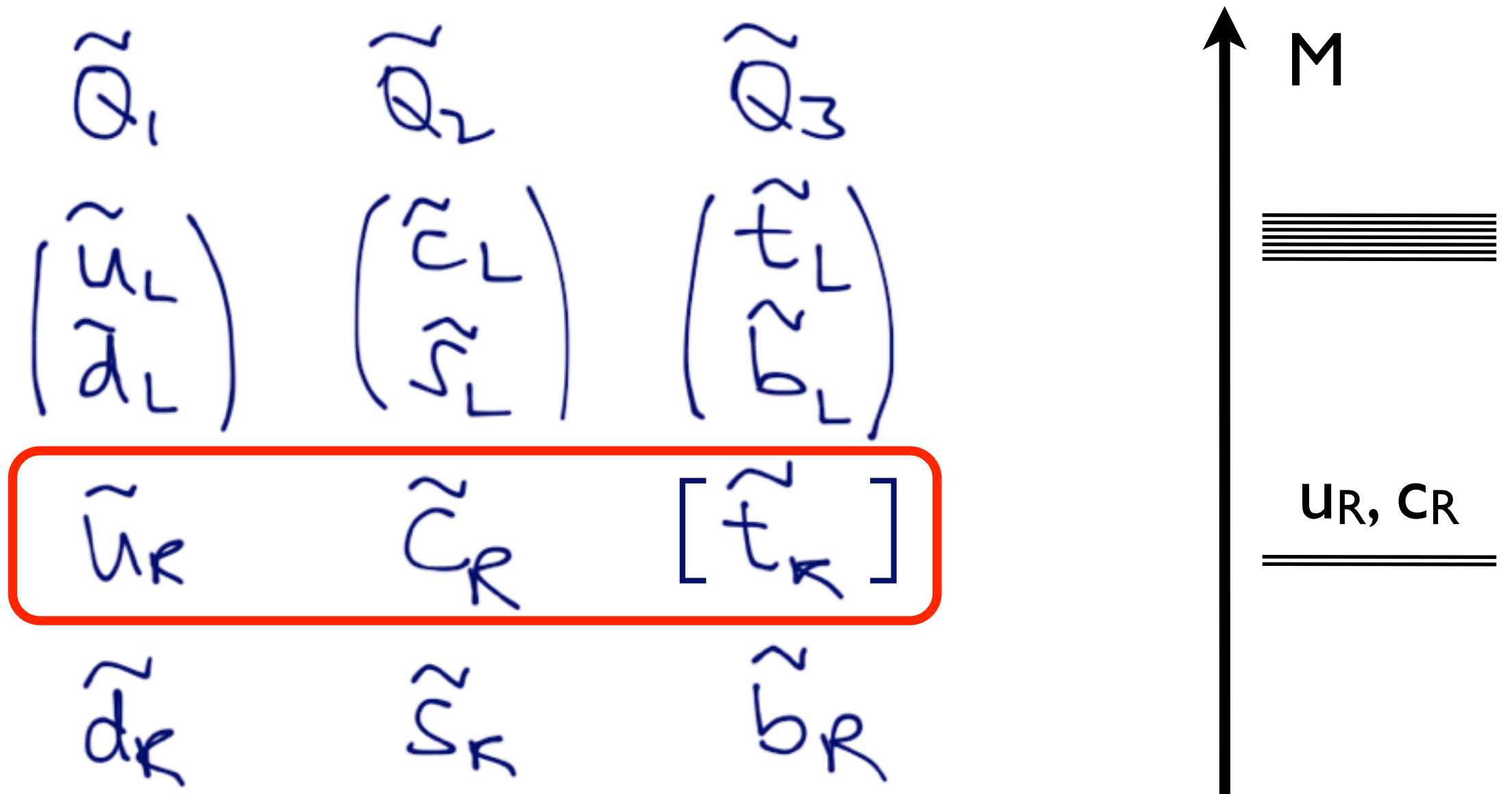
Vertical non-degeneracy

$\tilde{D}_{1,2}$ v. \tilde{g}

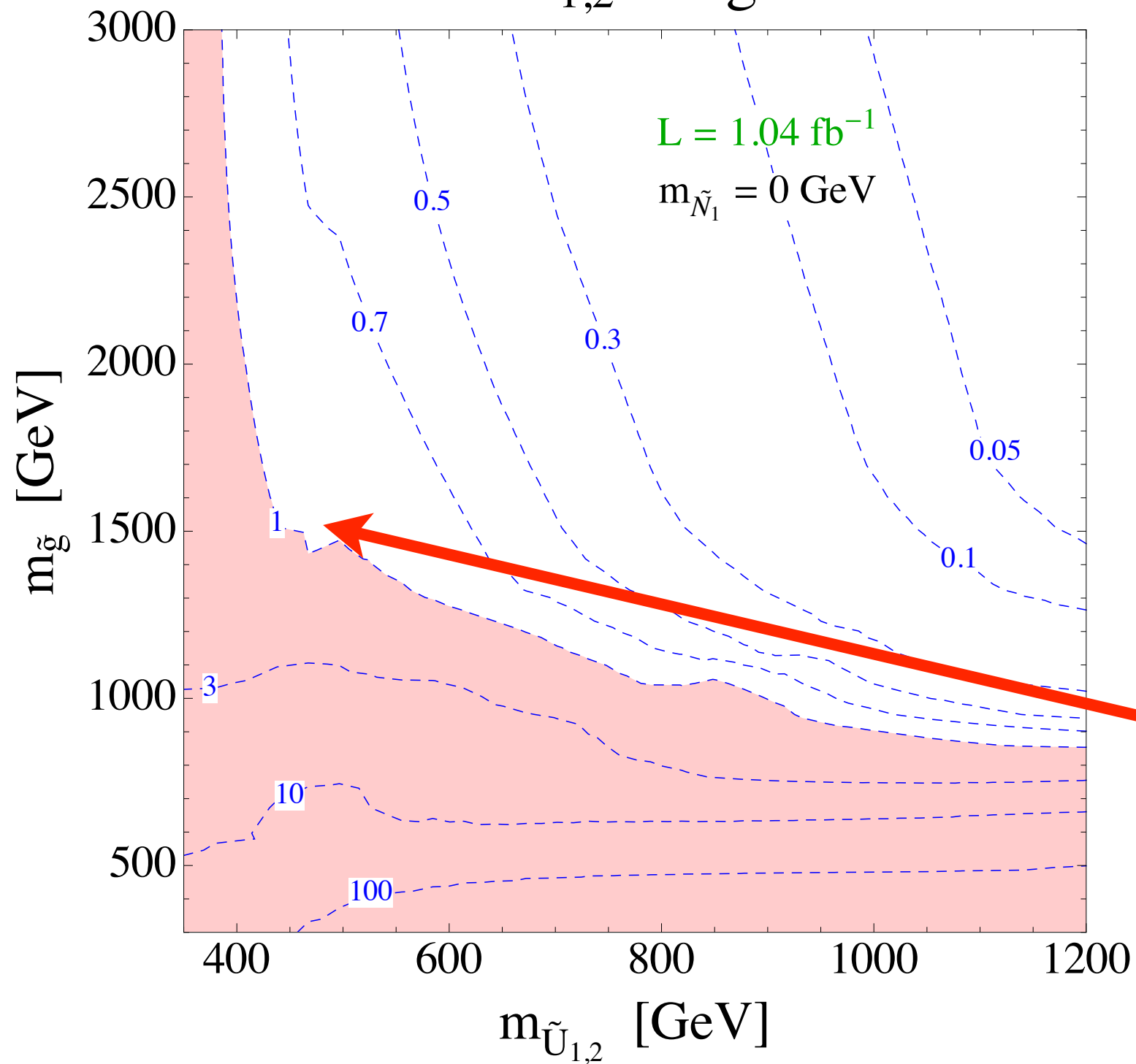


For $m_{\text{gluino}} > 1.2 \text{ TeV}$, limit can even disappear ($m_{\text{LSP}} = 0 \text{ GeV}$)

Split squarks w/ different quantum number (vertical)



$\tilde{U}_{1,2}$ v. \tilde{g}

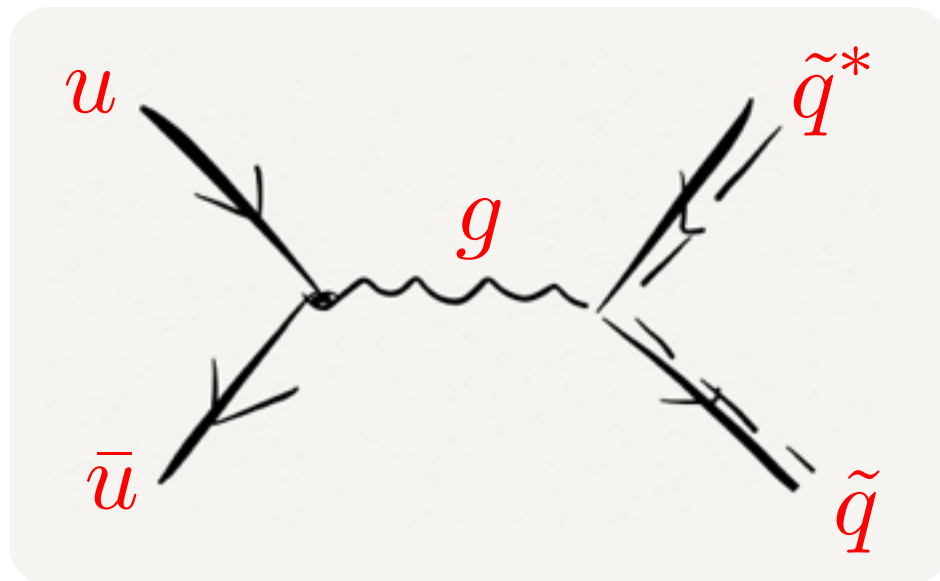


Stronger limits,
later decoupling
of gluino

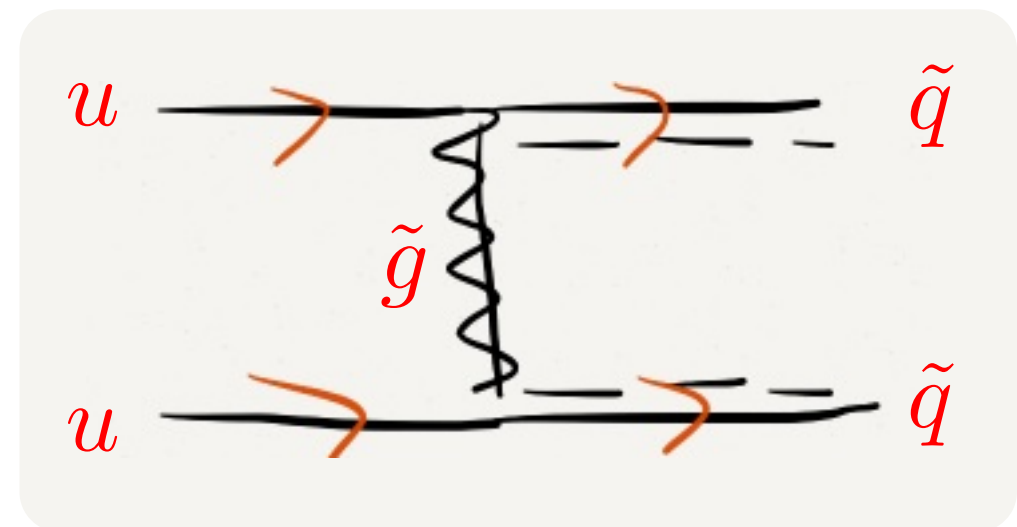
Why much more
sensitive to
gluino mass? →

What is driving the strong ATLAS/CMS limit?

Squark - Squark production:

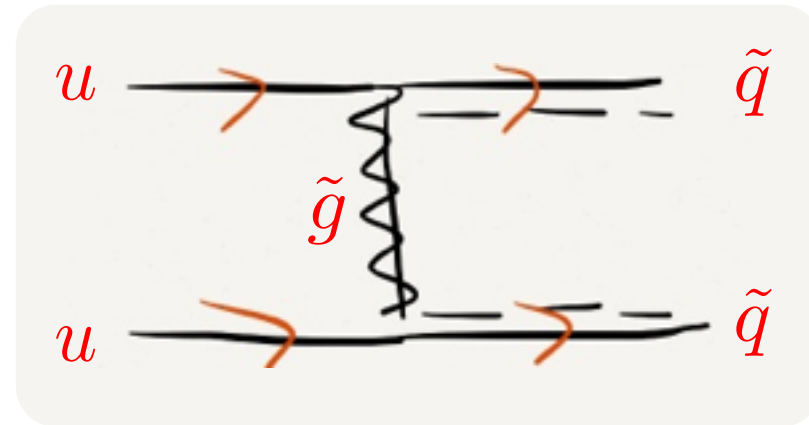
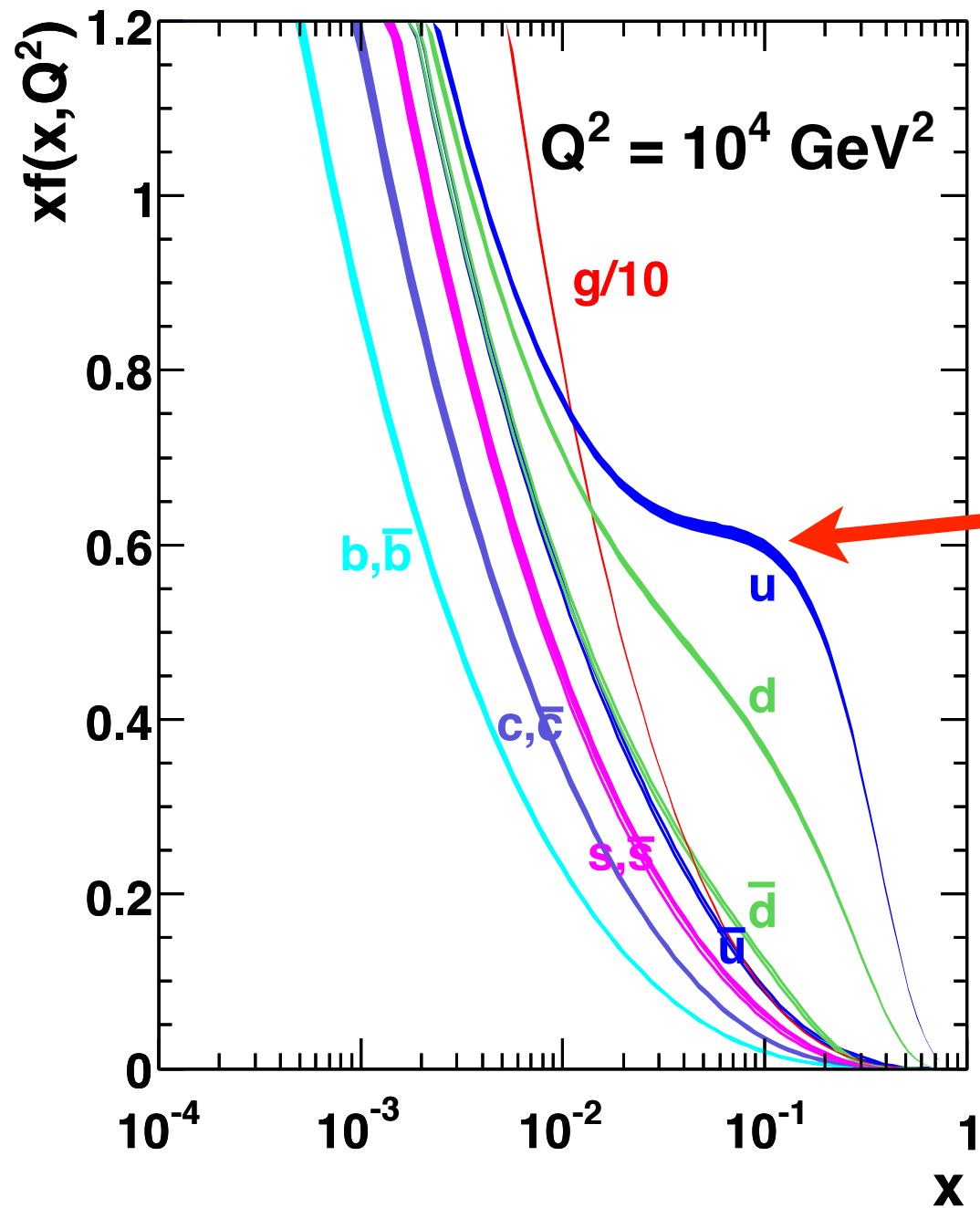


Independent of
squark flavor
(and gluino mass)



Majorana nature of
gluino allows **u u**
initial state!

Simple d.o.f rescaling



access to large **up**
quark pdf

$$E \approx x \cdot 7 \text{ TeV}$$

$$\frac{1}{m_{\tilde{g}}} \tilde{q}\tilde{q} u_R u_R \quad \text{dim5 op.}$$

$$\rightarrow \sigma \sim 1/m_{\tilde{g}}^2$$

slow decoupling

Aligned LH squarks

Squarks

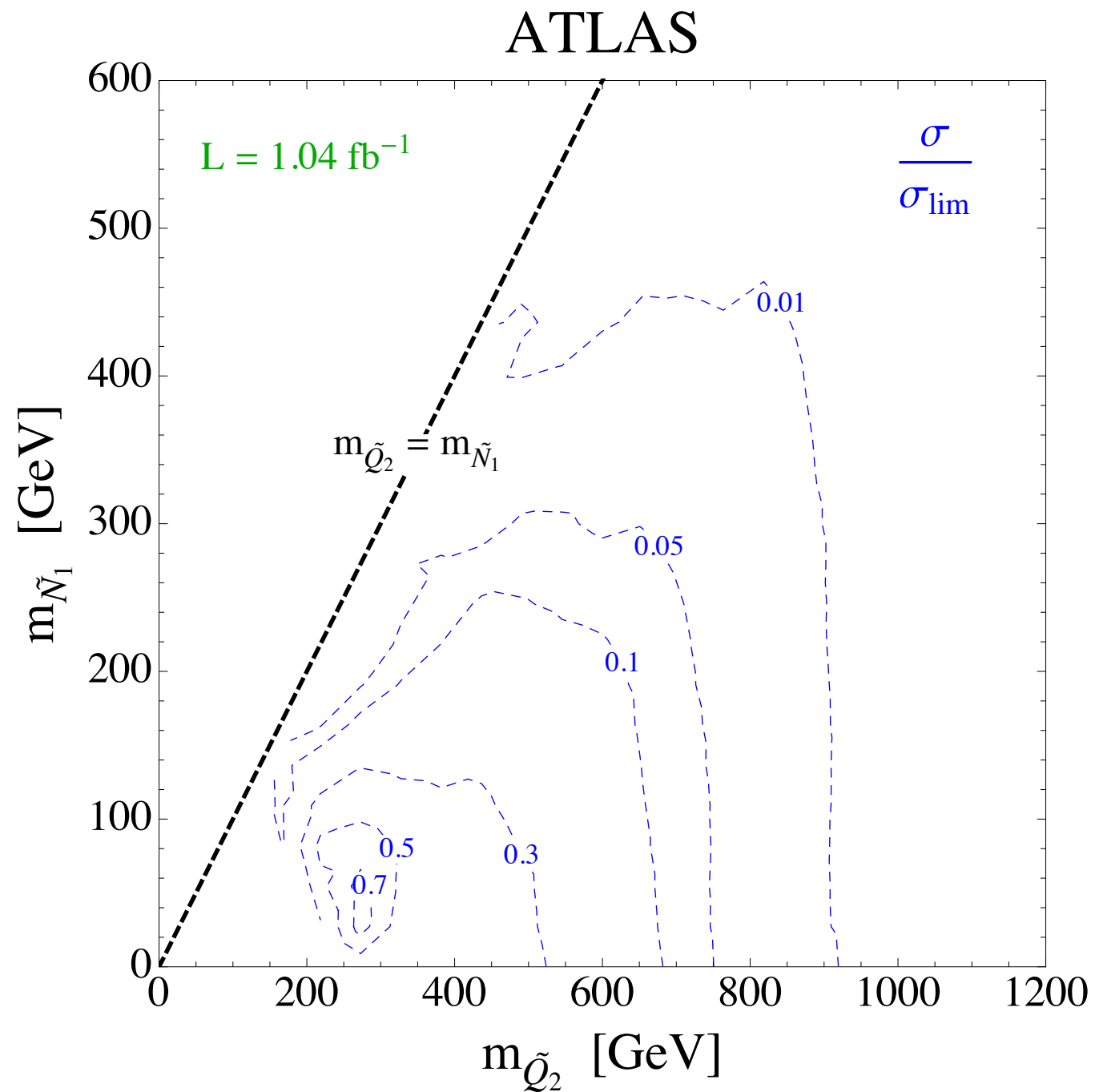
$$Q_1 \begin{pmatrix} u_L \\ d_L \end{pmatrix} \quad \begin{matrix} u_L \\ d_L \end{matrix}$$

$$Q_2 \begin{pmatrix} u_L \\ d_L \end{pmatrix} \quad \begin{matrix} u_L \\ d_L \end{matrix}$$

$$Q_3 \begin{pmatrix} u_L \\ d_L \end{pmatrix} \quad \begin{matrix} u_L \\ d_L \end{matrix}$$

2 d.o.f

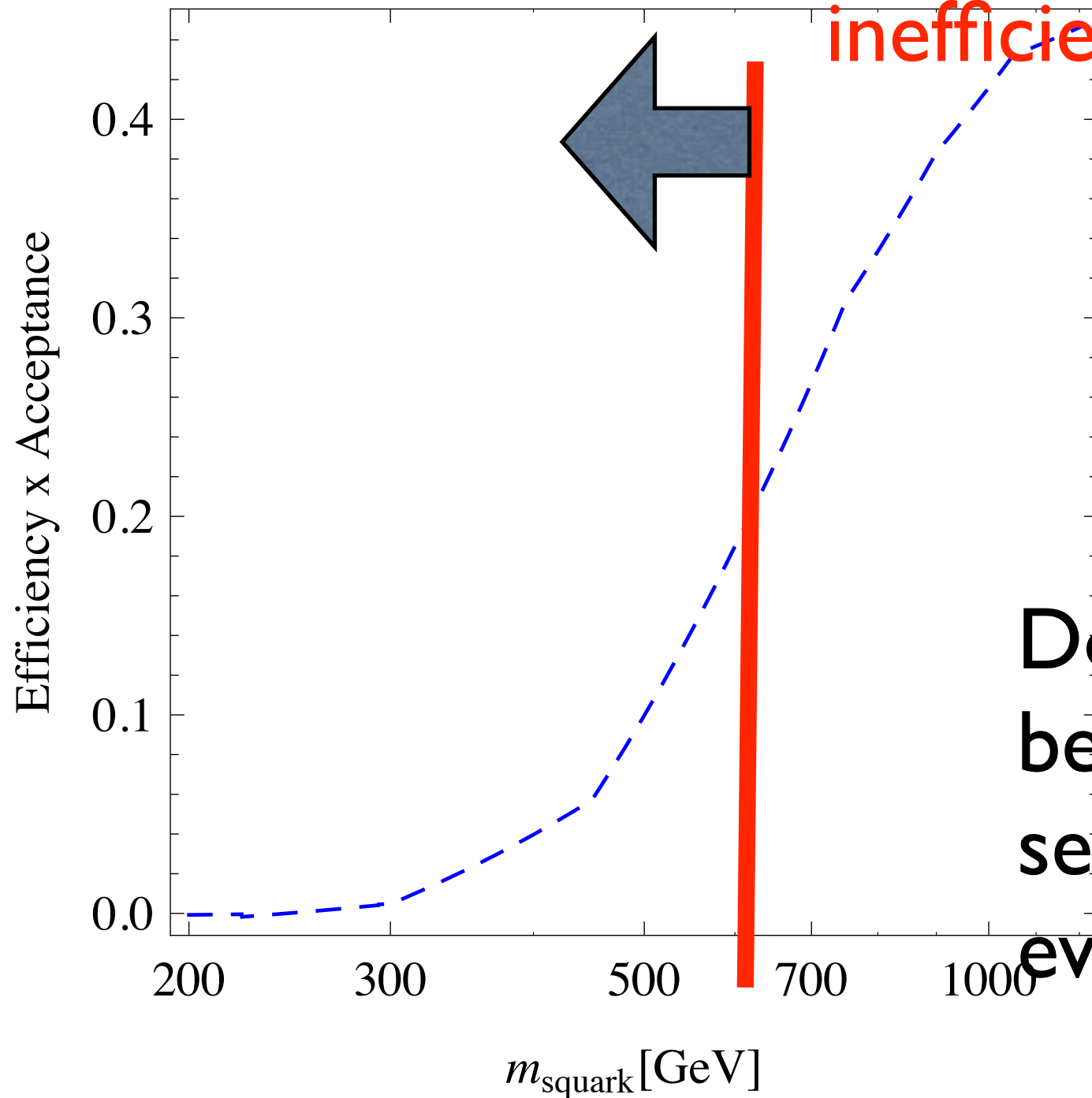
No limit on LH 2nd gen' squark



No bounds for \tilde{Q}_2 in 1/fb of data
(decoupled gluinos at 3TeV)

Efficiencies

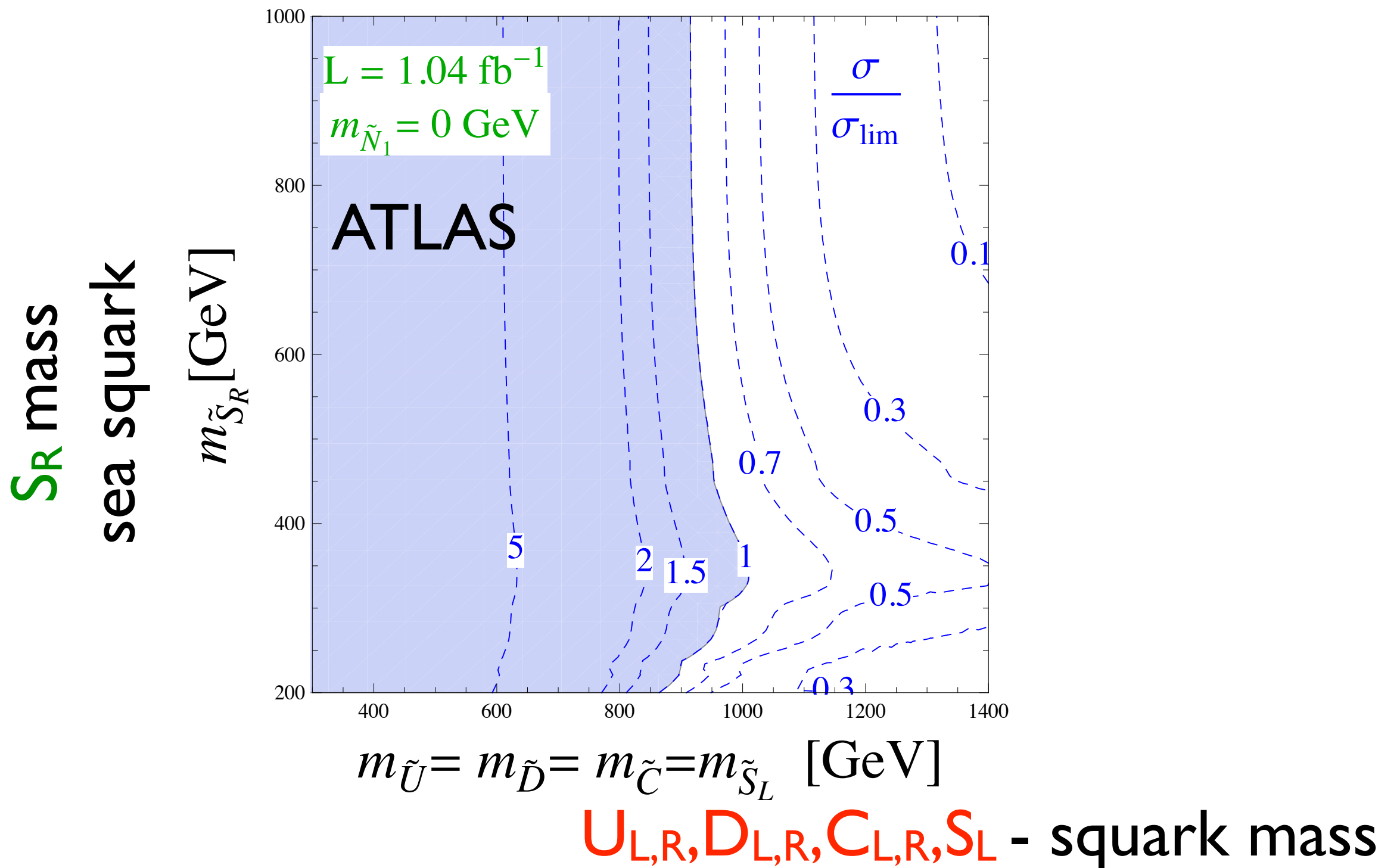
Searches have hard thresholds,
inefficient for light squarks



Example: ATLAS 1/fb
2jet, $M_{\text{eff}} > 1 \text{ TeV}$,
 $m_{\text{LSP}} = 0 \text{ GeV}$

Do not expect it to be much
better for higher luminosity
searches ($> 5 \text{ /fb}$) b/c
even harder cuts used

Decoupling of the remaining squarks?



Todo

- No efficiency maps available yet for 5 /fb searches. We **simulate and validate** using Monte-Carlo mockups (see part 2 on how to do that) to check limits
- Associated prod' of squarks of different mass important in various cases (b/c of efficiency behavior) → need to simulate also for 1/fb, cannot use Prospino NLO out-of-the-box for xsecs
- BUT: similar conclusions expected due to hard cuts on M_{eff} / H_T in 5 1/fb searches

work in progress with

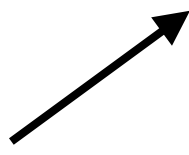
Michele Papucci, Josh Ruderman (LBL Berkely)

Gilad Perez, Rakhi Mahbubani (CERN)

Summary part I

- Squarks spectra can be **vertically** and **horizontally** split.
- Limits for 1st gen' squarks very dependent on gluino mass, for heavy gluino almost no limit
- Are there light squarks hiding in the data?
- **Need dedicated light squark searches!**

low mass



Natural SUSY

- Bottom-up naturalness reminder → Csaba's talk
- Current limits?
→ Weber's talk on 3rd gen squarks

h = linear combination of fields whose
vev breaks EW symmetry

$$V = m_H^2 |h|^2 + \frac{\lambda}{4} |h|^4 \quad m_h^2 = \lambda v^2 = -2m_H^2$$

$$\Delta = \frac{2|\delta m_H^2|}{m_h^2}$$

measures fine-tuning

Natural EWSB & SUSY

MSSM, NMSSM, DMSSM, ...

Fine-tuning of (Higgs mass)²

→ Bertuzzo's talk

$$\frac{m_{Higgs}^2}{2} = -|\mu|^2 + \dots + \delta m_H^2$$

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Higgsinos

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Higgsinos

1 loop

$$\delta m_H^2|_{stop} = -\frac{3}{8\pi^2} y_t^2 \left(m_{U_3}^2 + m_{Q_3}^2 + |A_t|^2 \right) \log \left(\frac{\Lambda}{\text{TeV}} \right)$$

stops, sbottom_L

2 loop

$$\delta m_H^2|_{gluino} = -\frac{2}{\pi^2} y_t^2 \left(\frac{\alpha_s}{\pi} \right) |M_3|^2 \log^2 \left(\frac{\Lambda}{\text{TeV}} \right)$$

gluino

- Amount of cancelation has not been directly probed yet! (experimental question)
- Interesting to look first for those cases where this cancelation is not strong (naturalness)

- Minimal requirements for a "natural" weak-scale SUSY?
 - 2 light stops
 - 1 light "left-handed" sbottom (near the stops by weak isospin)
 - light higgsinos, i.e. 2 neutralinos and 1 chargino
 - a not-too-heavy gluino

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Model dependence:

if low scale mediation, a light gravitino
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Model dependence:

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Rest could be decoupled...

What about numbers?

Difficult to make sharp quantitative statements (just a guidance):
What is “natural”? $10^{-9}=1$? $100^{-99}=1$? $1000^{-999}=1$? 1 part in
 10^4 ? ...

Stops:

$$\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim 600 \text{ GeV} \frac{\sin \beta}{(1 + x_t^2)^{1/2}} \left(\frac{\log(\Lambda / \text{TeV})}{3} \right)^{-1/2} \left(\frac{m_h}{120 \text{ GeV}} \right) \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

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$$\mu \lesssim 190 \text{ GeV} \left(\frac{m_h}{120 \text{ GeV}} \right) \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

Stops:

$$\sqrt{m_{\tilde{t}_1}^2 + m_{\tilde{t}_2}^2} \lesssim 600 \text{ GeV} \frac{\sin \beta}{(1 + x_t^2)^{1/2}} \left(\frac{\log(\Lambda / \text{TeV})}{3} \right)^{-1/2} \left(\frac{m_h}{120 \text{ GeV}} \right) \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

(e.g. Kitano & Nomura 2006)

Less problems w/ low
scale mediation

bound ameliorated if physics
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Gluginos:

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(Bino, Wino)

$$(M_1, M_2) \lesssim (2.7 \text{ TeV}, 870 \text{ GeV}) \left(\frac{\log(\Lambda / \text{TeV})}{3} \right)^{-1/2} \left(\frac{m_h}{120 \text{ GeV}} \right) \left(\frac{\Delta^{-1}}{20\%} \right)^{-1/2}$$

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1st-2nd gen' squarks not very constrained

3rd gen spectrum:

Two Stops: $\begin{pmatrix} m_{Q_3}^2 + (165\text{GeV})^2 & m_t(A_t - \mu \cot \beta) \\ m_t(A_t - \mu \cot \beta) & m_{U_3}^2 + (170\text{GeV})^2 \end{pmatrix}$

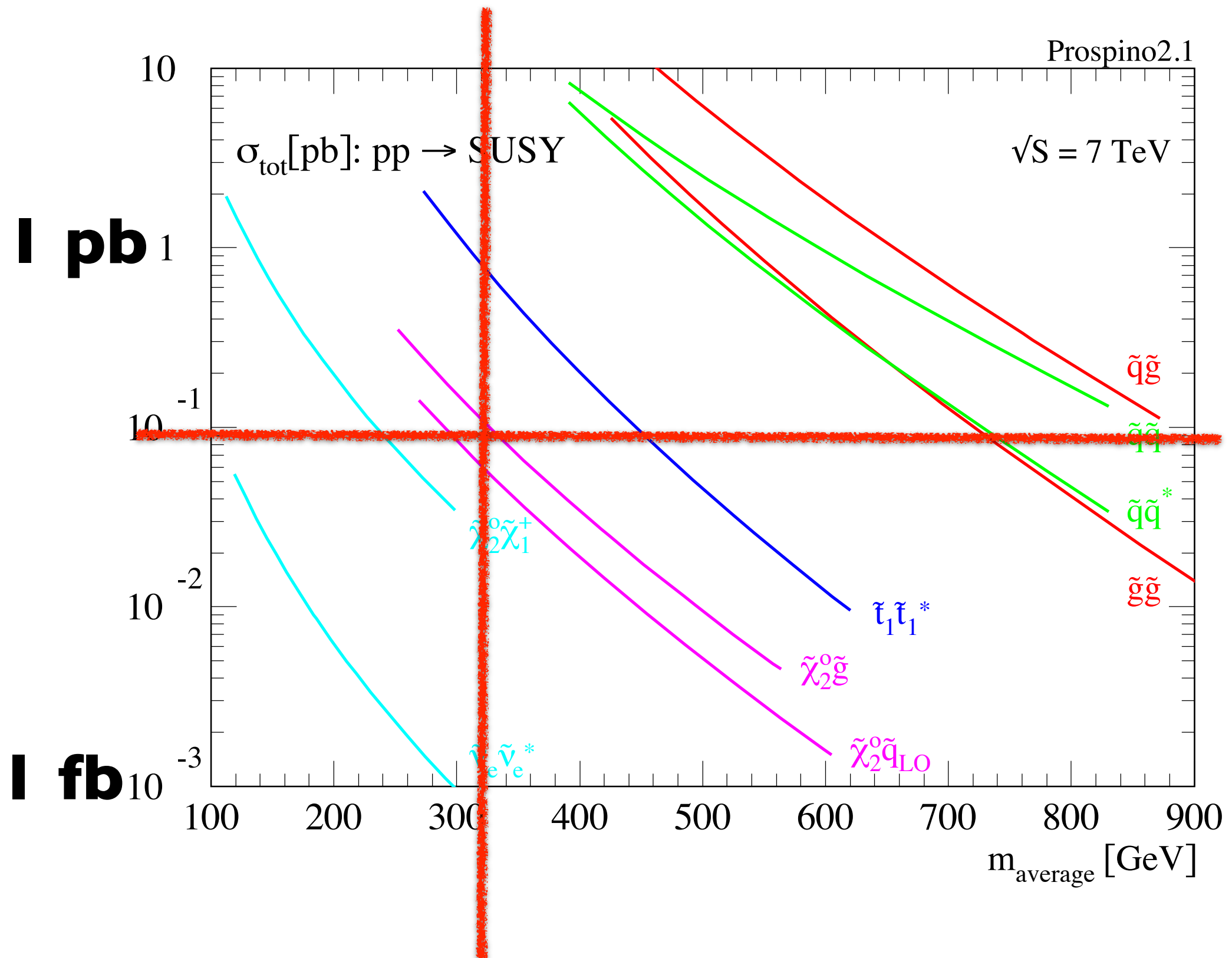
One Sbottom: $m_{Q_3}^2 + (70\text{GeV})^2$

$X_t = A_t - \mu \cot \beta$

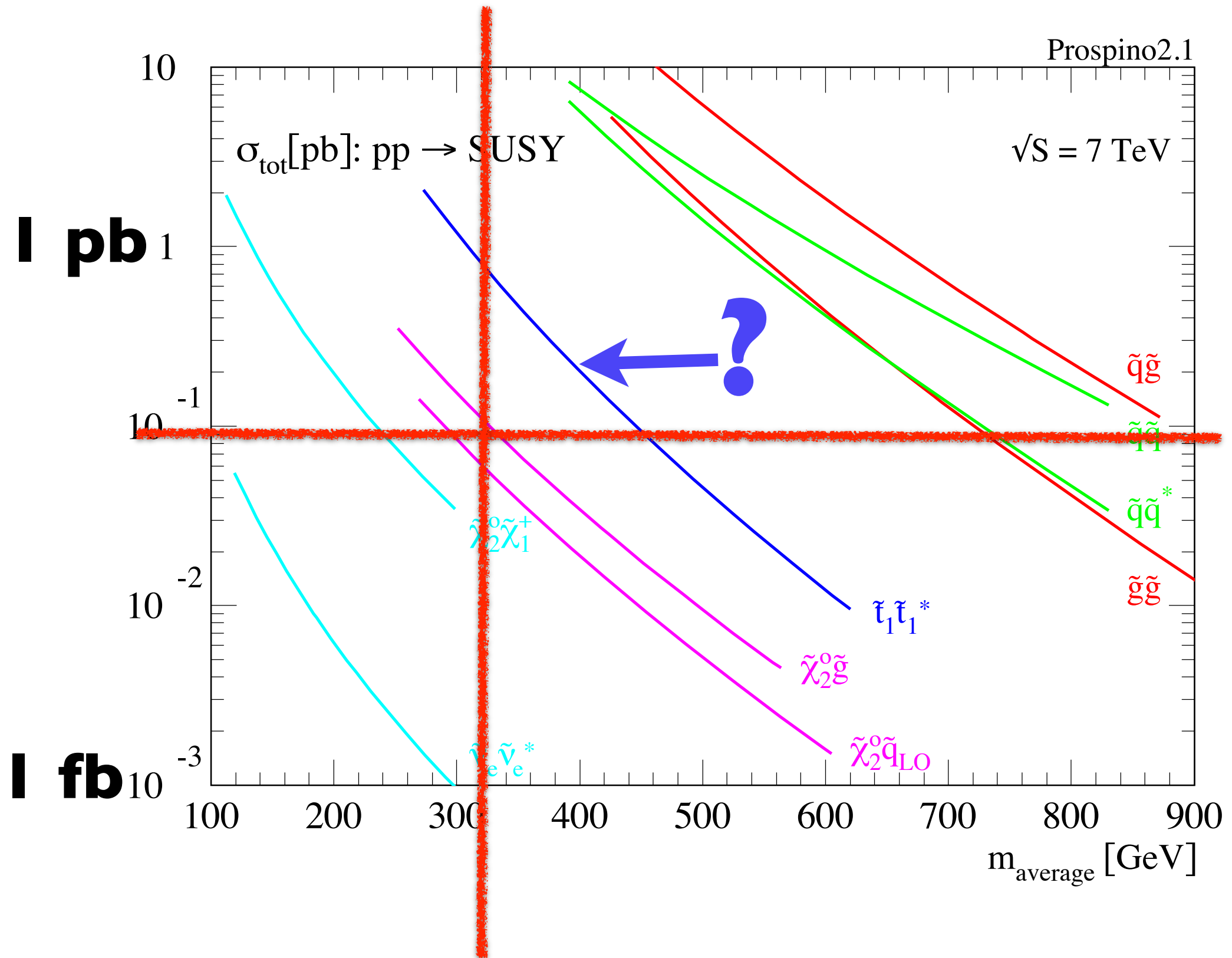
Fixed by Isospin

The other sbottom mass is basically free

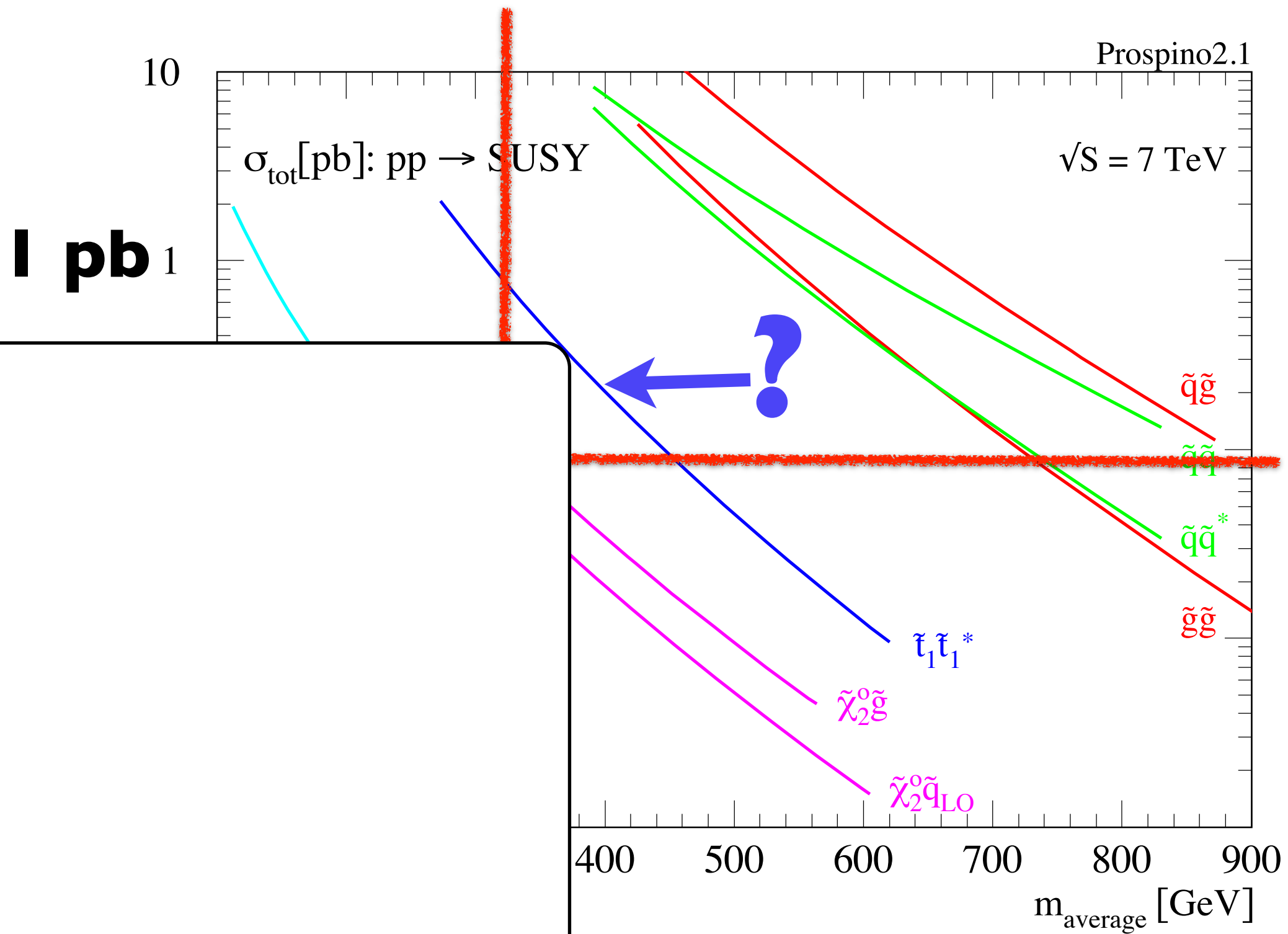
Reach?



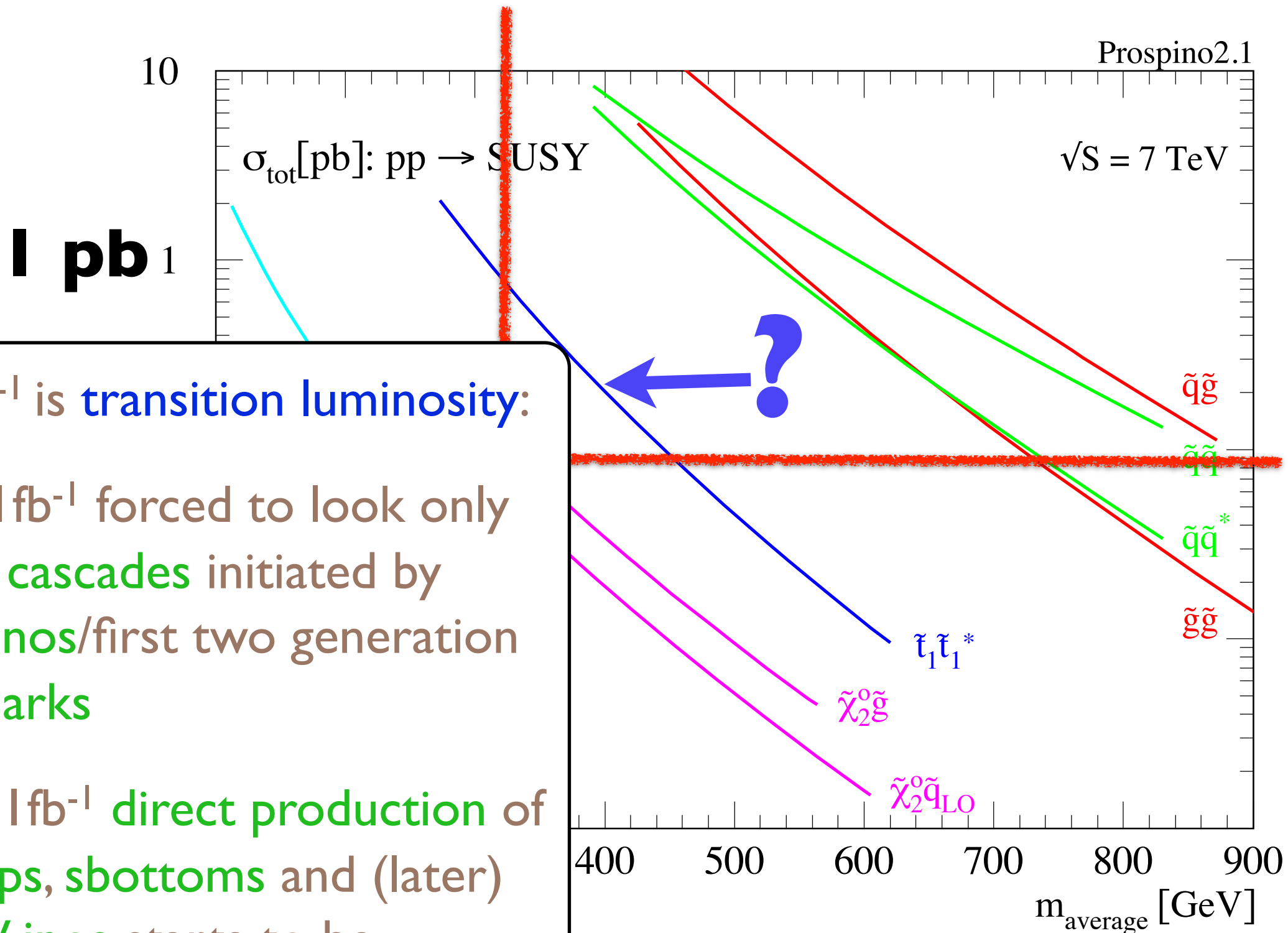
Reach?



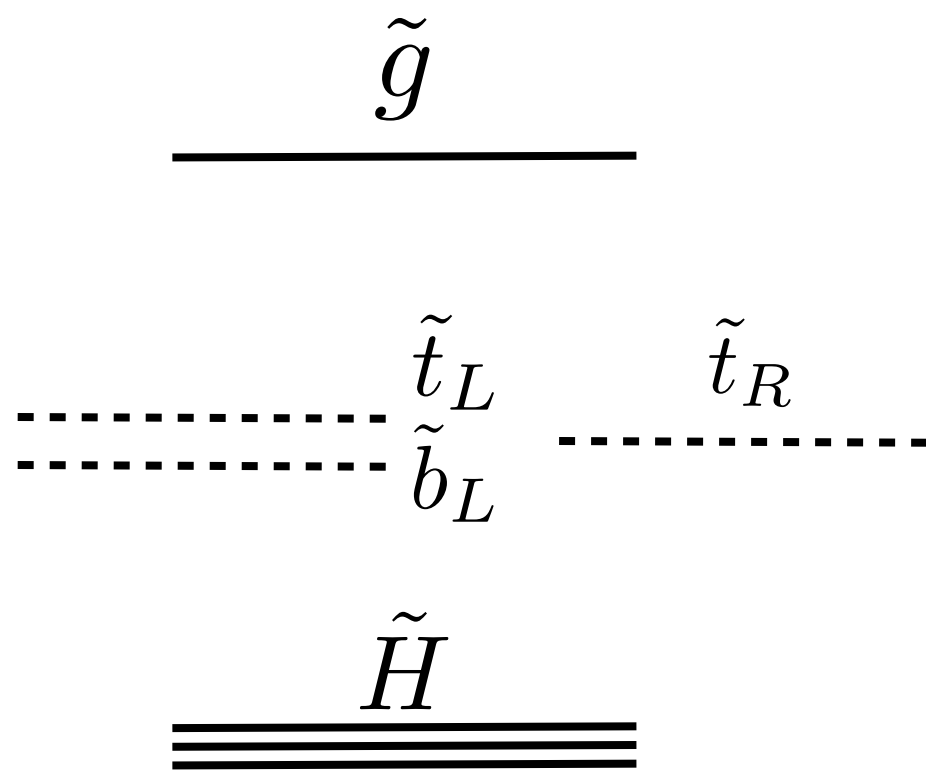
Reach?



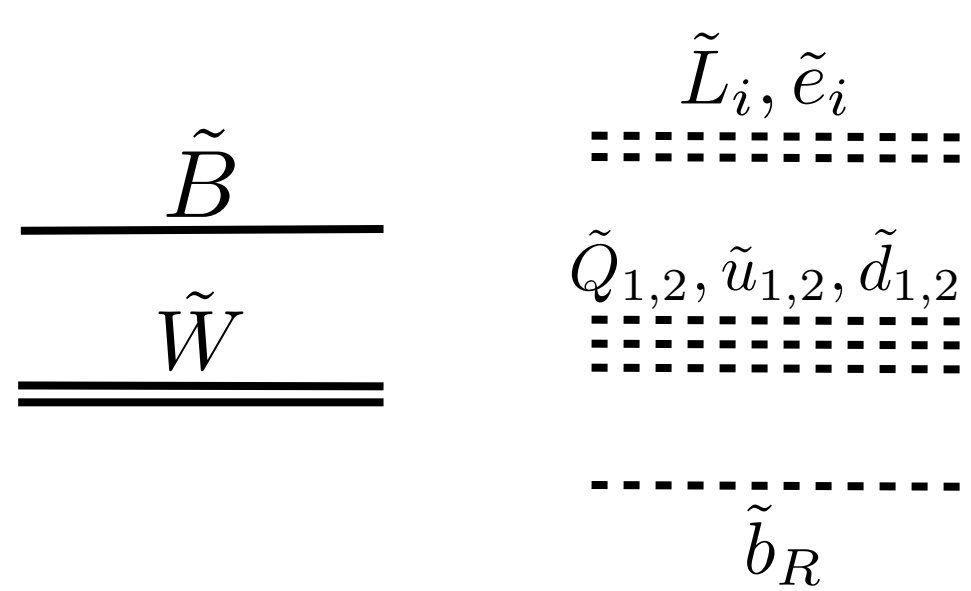
Reach?



- 1 fb^{-1} is transition luminosity:
- $L < 1 \text{ fb}^{-1}$ forced to look only for **cascades** initiated by **gluinos**/first two generation **squarks**
- $L > 1 \text{ fb}^{-1}$ **direct production** of **stops, sbottoms** and (later) **EW-inos** starts to be accessible



natural SUSY



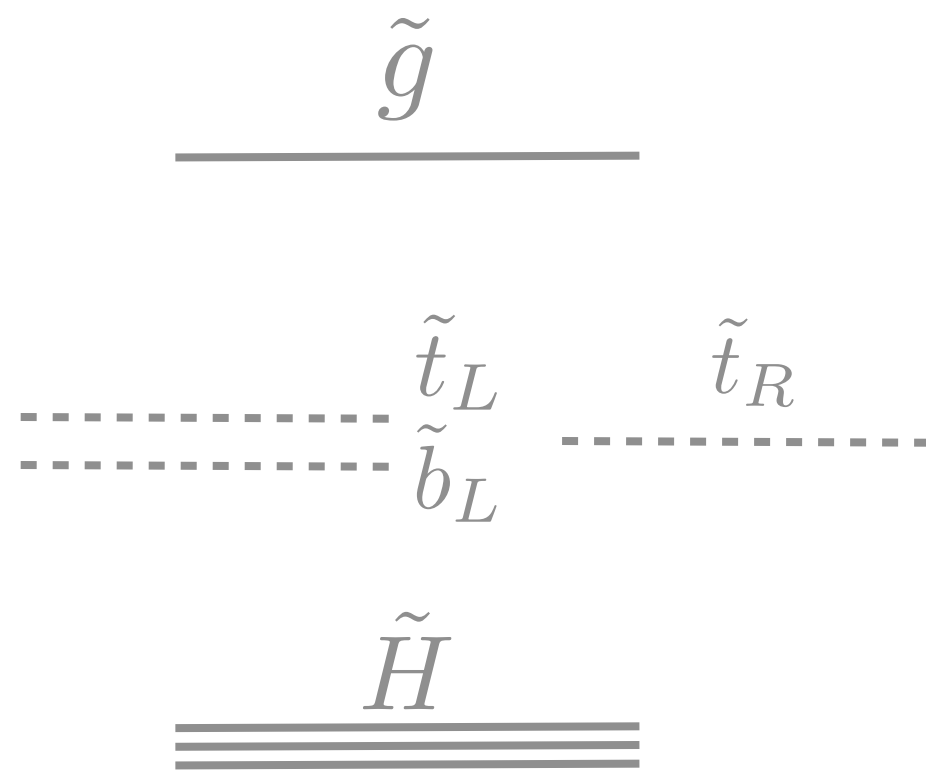
decoupled SUSY

Related:

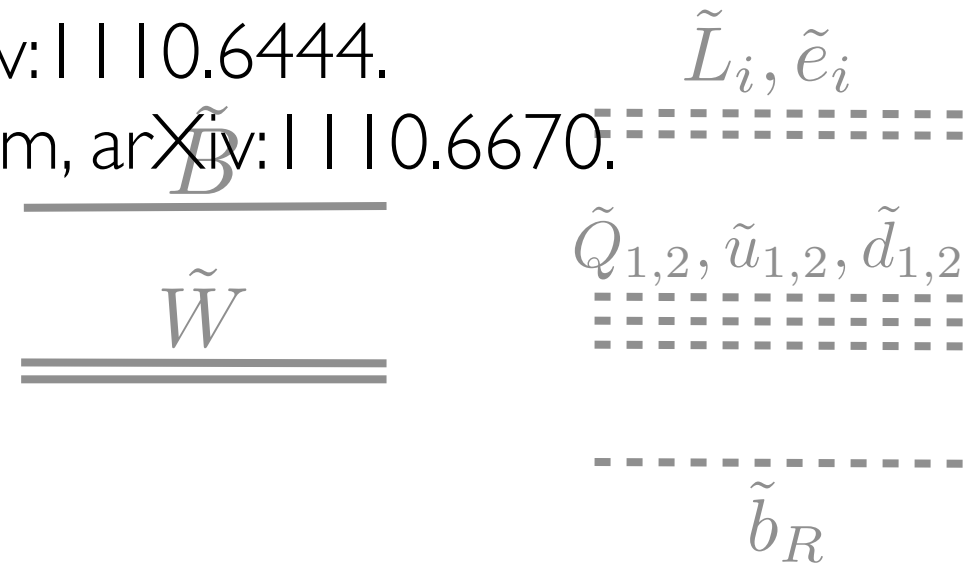
R. Essig, E. Izaguirre, J. Kaplan, J. G. Wacker, arXiv:1110.6443.

Y. Kats, P. Meade, M. Reece, D. Shih, arXiv:1110.6444.

C. Brust, A. Katz, S. Lawrence, R. Sundrum, arXiv:1110.6670.



natural SUSY



Natural Susy endures

[arXiv:1110.6926](https://arxiv.org/abs/1110.6926)

M. Papucci, J. Ruderman, AW

decoupled SUSY

Large signature space

arXiv:1110.6926

	ATLAS			CMS		
	channel	\mathcal{L} [fb ⁻¹]	ref.	channel	\mathcal{L} [fb ⁻¹]	ref.
jets + \cancel{E}_T	2-4 jets	1.04	[1]	α_T	1.14	[11]
	6-8 jets	1.34	[2]	H_T, \cancel{H}_T	1.1	[12]
b-jets (+ l's + \cancel{E}_T)	1b, 2b	0.83	[3]	m_{T2} (+ b)	1.1	[13]
	b + 1l	1.03	[4]	1b, 2b	1.1	[14]
				$b'b' \rightarrow b + l^\pm l^\pm, 3l$	1.14	[15]
				$t't' \rightarrow 2b + l^+ l^-$	1.14	[16]
multilepton (+ \cancel{E}_T)	1l	1.04	[5]	1l	1.1	[17]
	$\mu^\pm \mu^\pm$	1.6	[6]	SS dilepton	0.98	[18]
	$t\bar{t} \rightarrow 2l$	1.04	[7]	OS dilepton	0.98	[19]
	$t\bar{t} \rightarrow 1l$	1.04	[8]	$Z \rightarrow l^+ l^-$	0.98	[20]
	4l	1.02	[9]	3l, 4l + \cancel{E}_T	2.1	[21]
	2l	1.04	[10]	3l, 4l	2.1	[22]

non susy
analyses

Large signature space

arXiv:1110.6926

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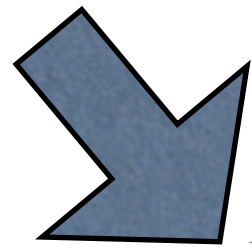
non susy
analyses

too
recent

Bgd's are left to the experimentalists... stay out of control regions!

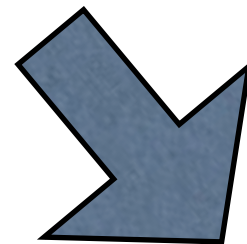
Process	Signal Region				
	$\geq 2\text{-jet}$	$\geq 3\text{-jet}$	$\geq 4\text{-jet}, m_{\text{eff}} > 500 \text{ GeV}$	$\geq 4\text{-jet}, m_{\text{eff}} > 1000 \text{ GeV}$	High mass
$Z/\gamma\text{+jets}$	$32.3 \pm 2.6 \pm 6.9$	$25.5 \pm 2.6 \pm 4.9$	$209 \pm 9 \pm 38$	$16.2 \pm 2.2 \pm 3.7$	$3.3 \pm 1.0 \pm 1.3$
$W\text{+jets}$	$26.4 \pm 4.0 \pm 6.7$	$22.6 \pm 3.5 \pm 5.6$	$349 \pm 30 \pm 122$	$13.0 \pm 2.2 \pm 4.7$	$2.1 \pm 0.8 \pm 1.1$
$t\bar{t}\text{+ single top}$	$3.4 \pm 1.6 \pm 1.6$	$5.9 \pm 2.0 \pm 2.2$	$425 \pm 39 \pm 84$	$4.0 \pm 1.3 \pm 2.0$	$5.7 \pm 1.8 \pm 1.9$
QCD multi-jet	$0.22 \pm 0.06 \pm 0.24$	$0.92 \pm 0.12 \pm 0.46$	$34 \pm 2 \pm 29$	$0.73 \pm 0.14 \pm 0.50$	$2.10 \pm 0.37 \pm 0.82$
Total	$62.4 \pm 4.4 \pm 9.3$	$54.9 \pm 3.9 \pm 7.1$	$1015 \pm 41 \pm 144$	$33.9 \pm 2.9 \pm 6.2$	$13.1 \pm 1.9 \pm 2.5$
Data	58	59	1118	40	18

Table 2: Fitted background components in each SR, compared with the number of events observed in data. The $Z/\gamma\text{+jets}$ background is constrained with control regions CR1a and CR1b, the QCD multi-jet, W and top quark backgrounds by control regions CR2, CR3 and CR4, respectively. In each case the first (second) quoted uncertainty is statistical (systematic). Background components are partially correlated and hence the uncertainties (statistical and systematic) on the background estimates do not equal the quadrature sums of the uncertainties on the components.



[5] has improved the ATLAS reach at large m_0 . The five signal regions are used to set limits on $\sigma_{\text{new}} = \sigma A \epsilon$, for non-SM cross-sections (σ) for which ATLAS has an acceptance A and a detection efficiency of ϵ [44]. The excluded values of σ_{new} are 22 fb, 25 fb, 429 fb, 27 fb and 17 fb, respectively, at the 95% confidence level.

upper bound on signal xsec



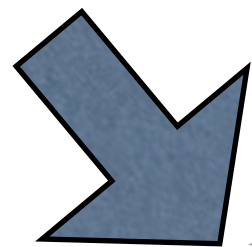
LIMIT!



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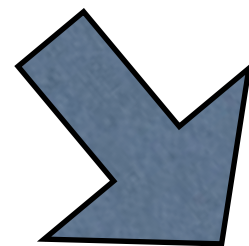
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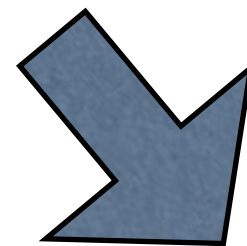


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upper bound on signal xsec



“Only” need **efficiency x Acceptance** of the signal bins for your model...



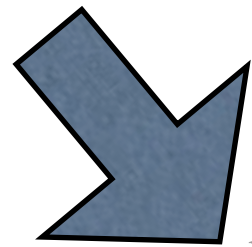
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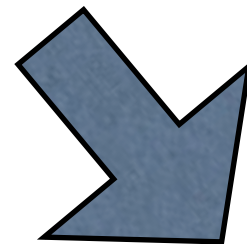
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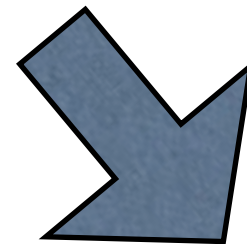


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upper bound on signal xsec



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LIMIT!



our pipelines

ATOM

public code soon

pythia / herwig / etc

+prospino

fastjet

truth leptons / photons /b's

- l/gamma iso
- parameterized efficiencies

checks sensitivity of cut & leakage in control region

pgs

pythia

+prospino

crude detector sim

cone jets

truth
muons/b's

- parameterized efficiencies

crude
simulated e/
gamma

ATOM

an **A**utomated **T**ester **O**f **M**odels

(soon to be) Public Tool developed by

QCD/Jets: C. Bauer (Berkeley), C. Vermillion (Berkeley)

BSM: M. Papucci (Berkeley),
T. Volansky (Tel Aviv), **A.W.** (DESY)

Calibration

“theorist limits”

To calibrate compare:

- 1) key kinematical distributions
- 2) limits

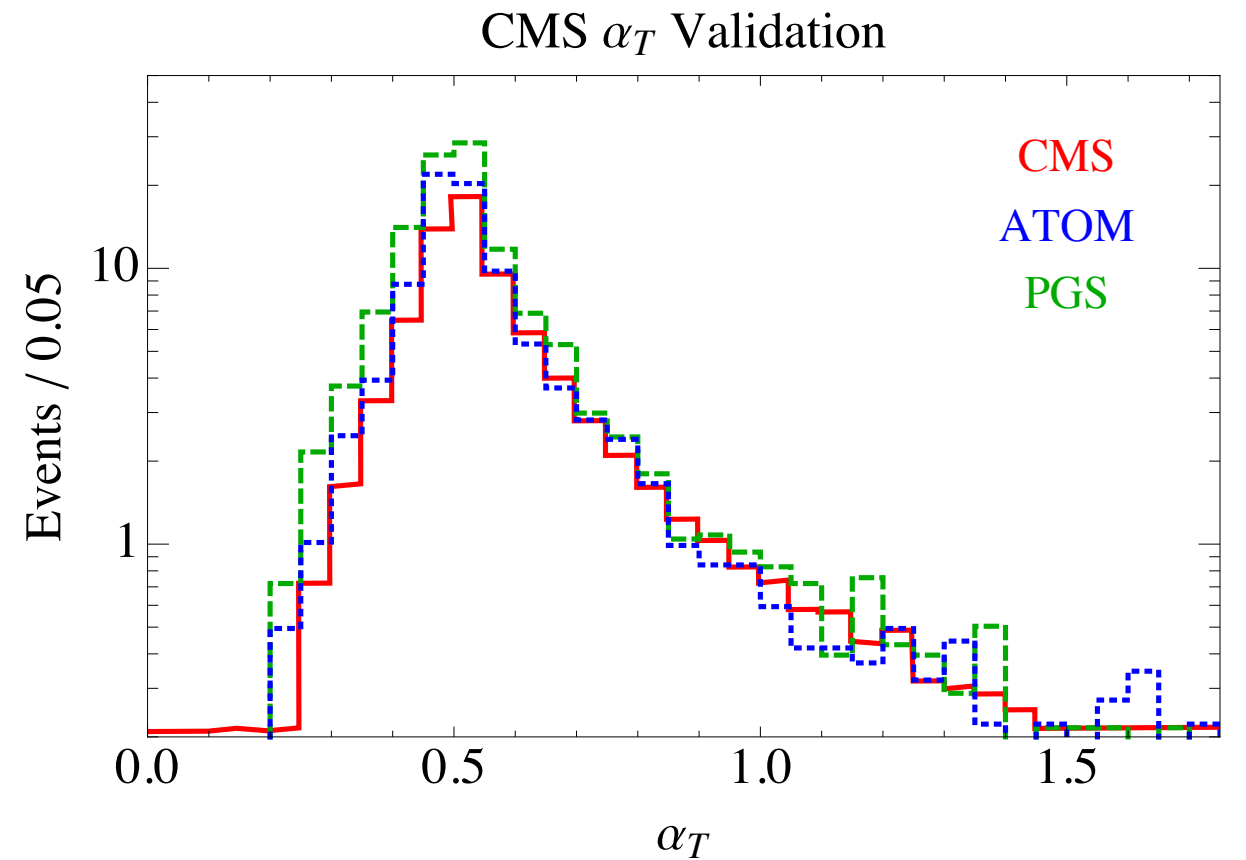
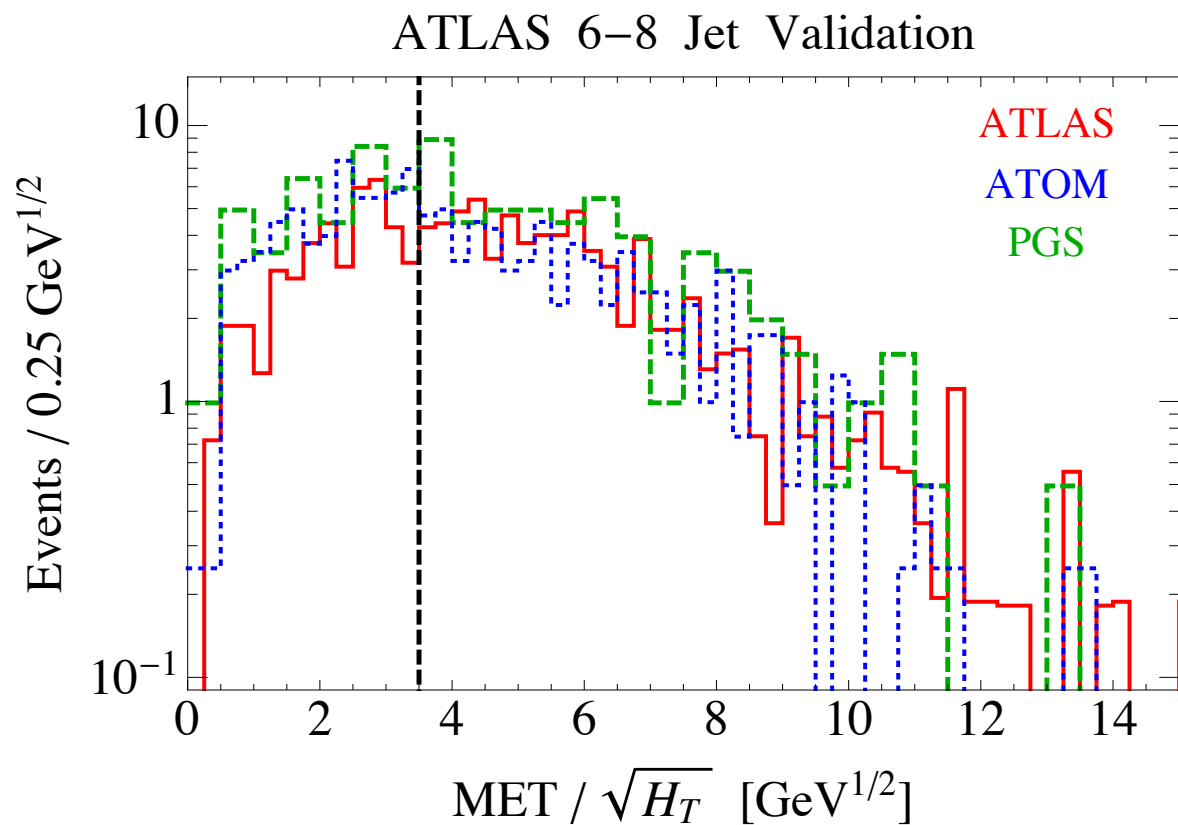
Calibration

“theorist limits”

To calibrate compare:

- 1) key kinematical distributions
- 2) limits

simplified models work best!



Check:

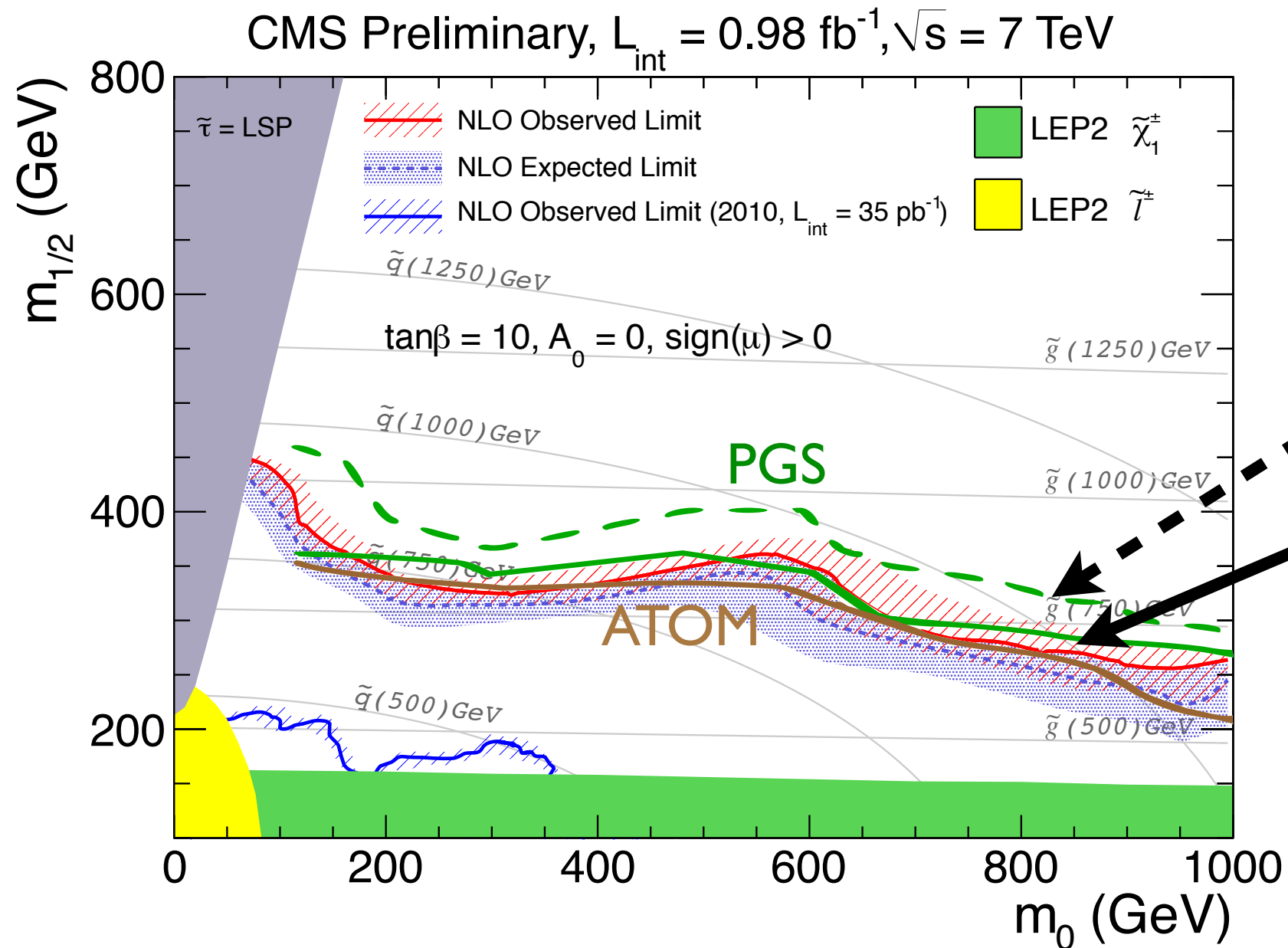
- kinematic distortions (**shape**)
- signal $\epsilon \times \mathcal{A}$ (**normalization**)

+ compare to all available limit plots...

~ 50 GeV accuracy (usually better)

Compare limits

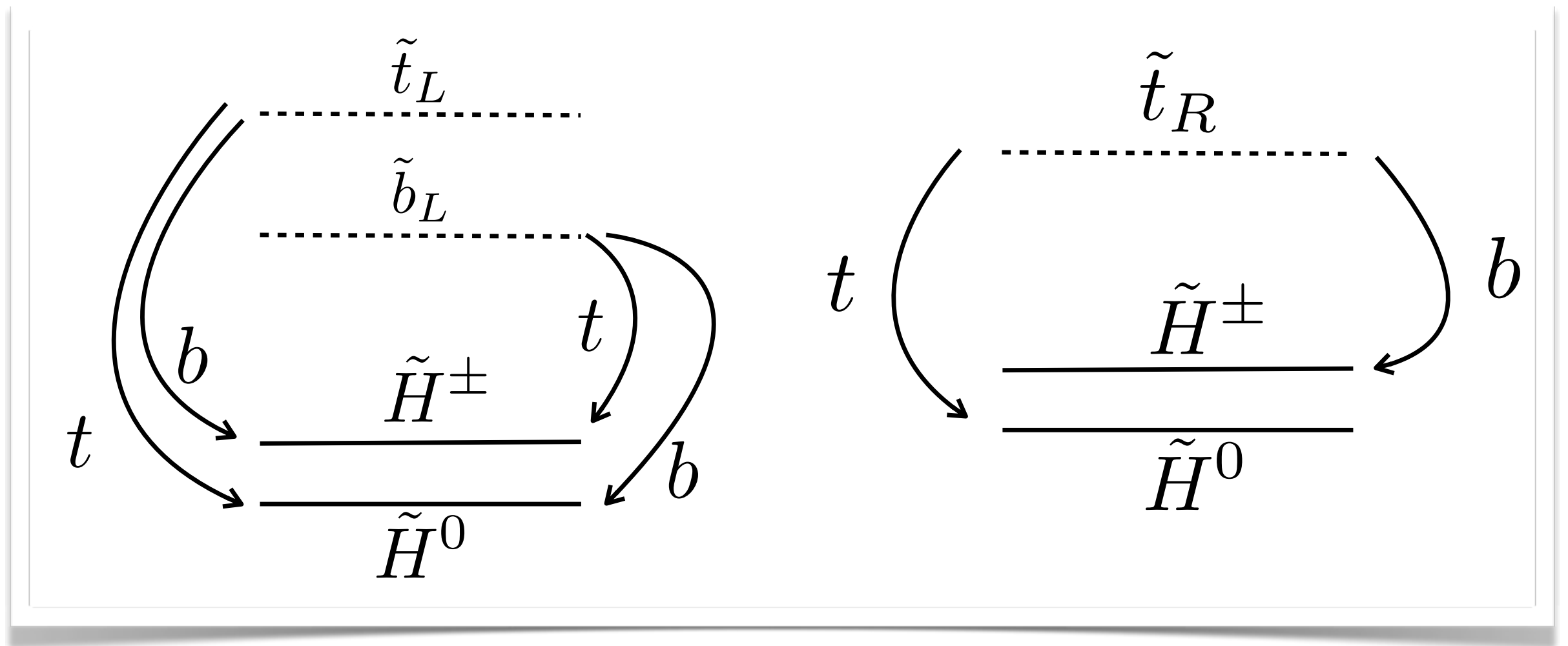
Example: Same-Sign dilepton by CMS



What we actually do

- Simplified topologies to slice the natural SUSY parameter space
- Highlight relevant kinematic features
- Which analyses give the most powerful constraints? Not necessarily those that are “designed” for the “signal”.
 - More recent results ($> 1/\text{fb}$) not included (update in progress)

Stops (sbottom) + Higgsinos

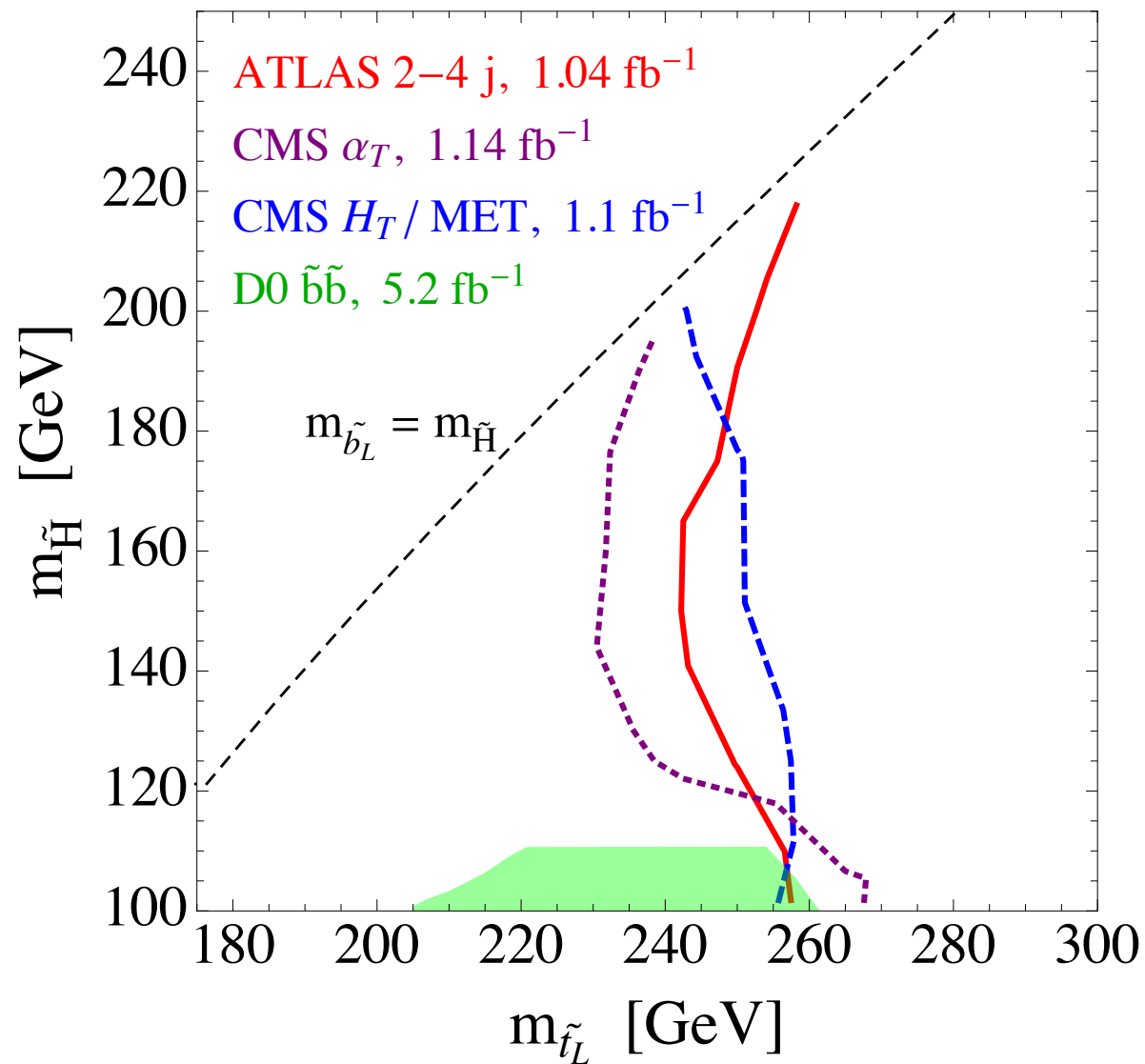


Stops can act as “sbottom” (bjet+ χ) !

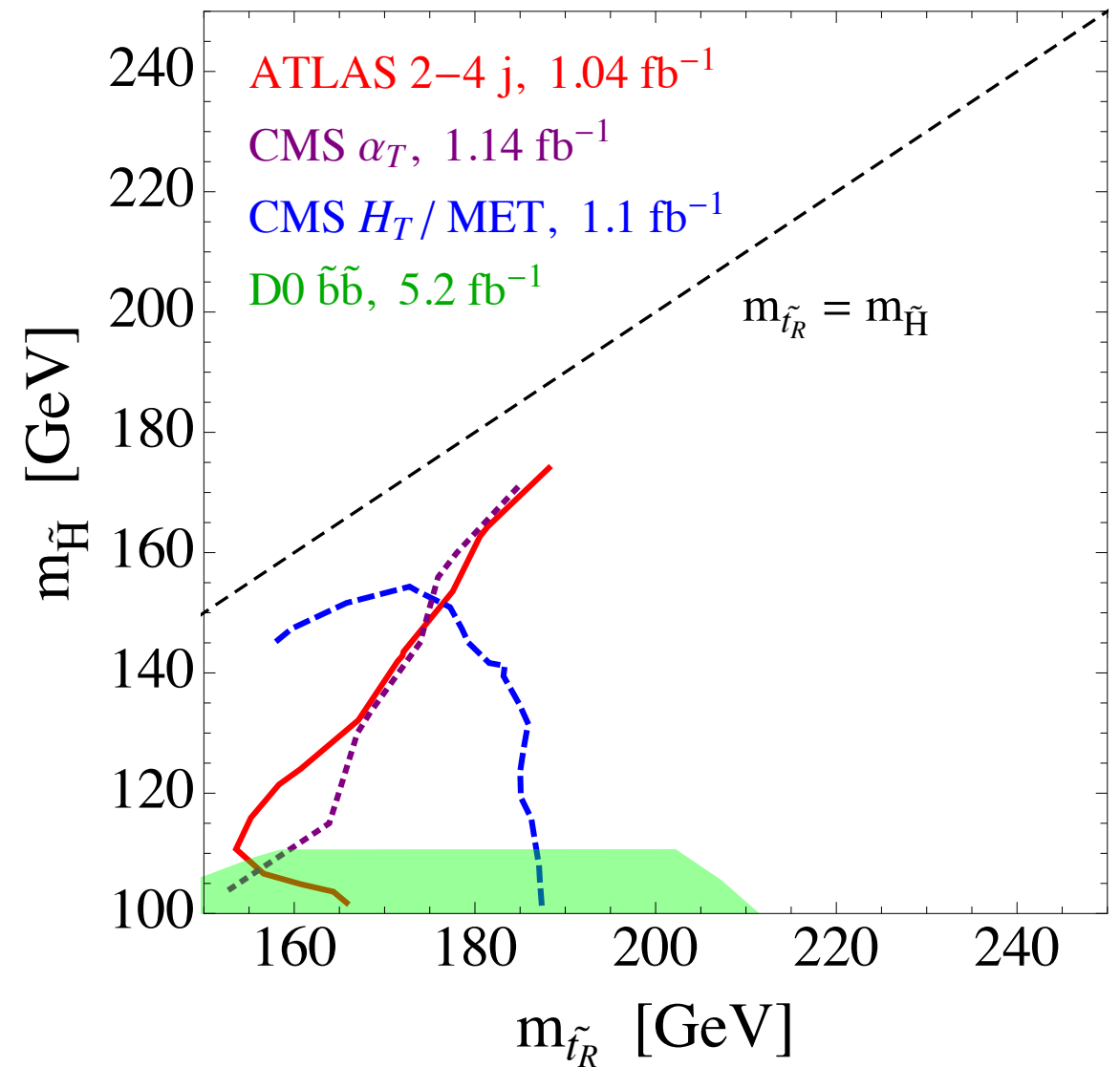
Chargino-neutralino splitting irrelevant for present searches

Stops (sbottom) + Higgsinos

Left-Handed Stop / Sbottom



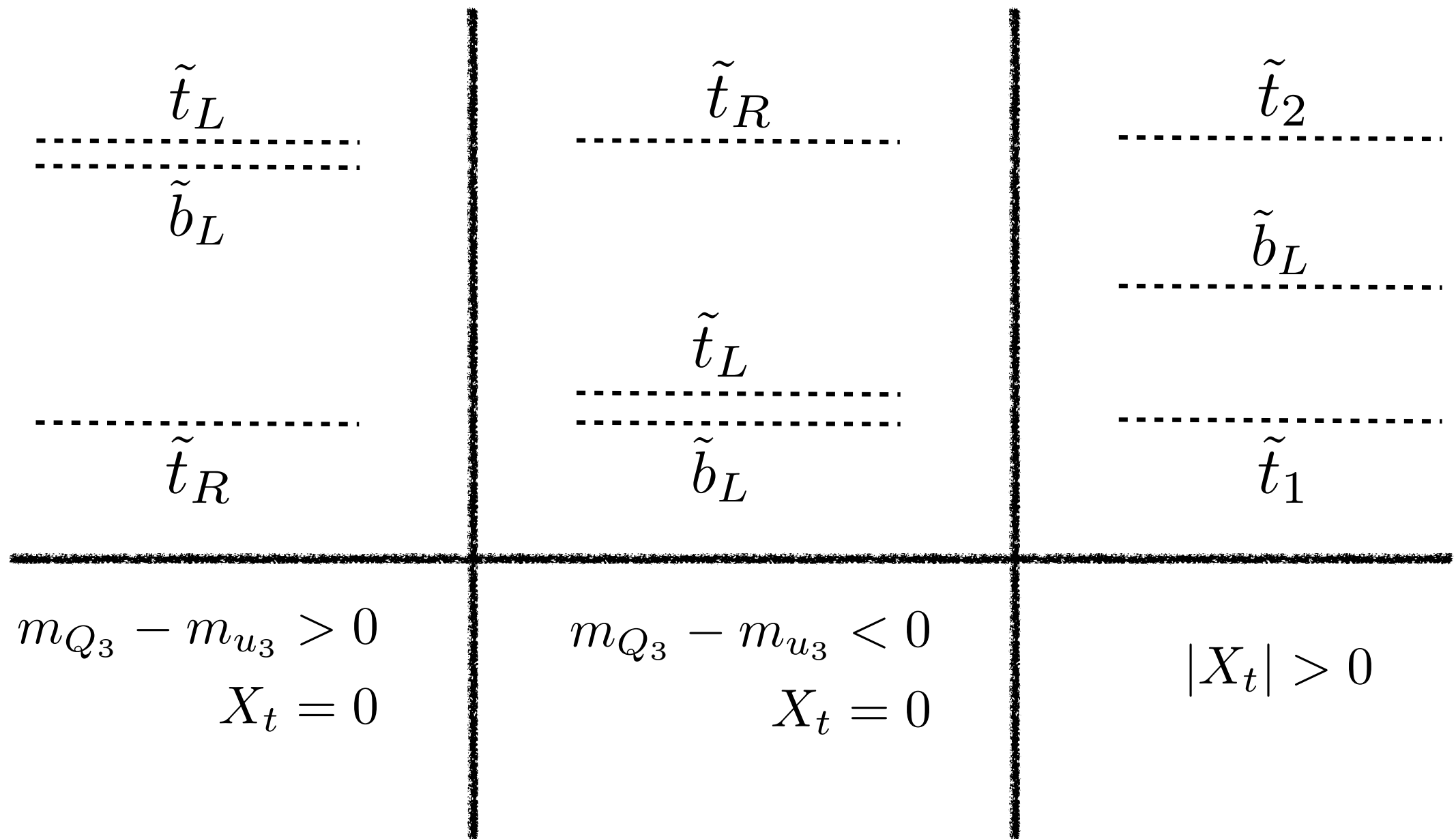
Right-Handed Stop



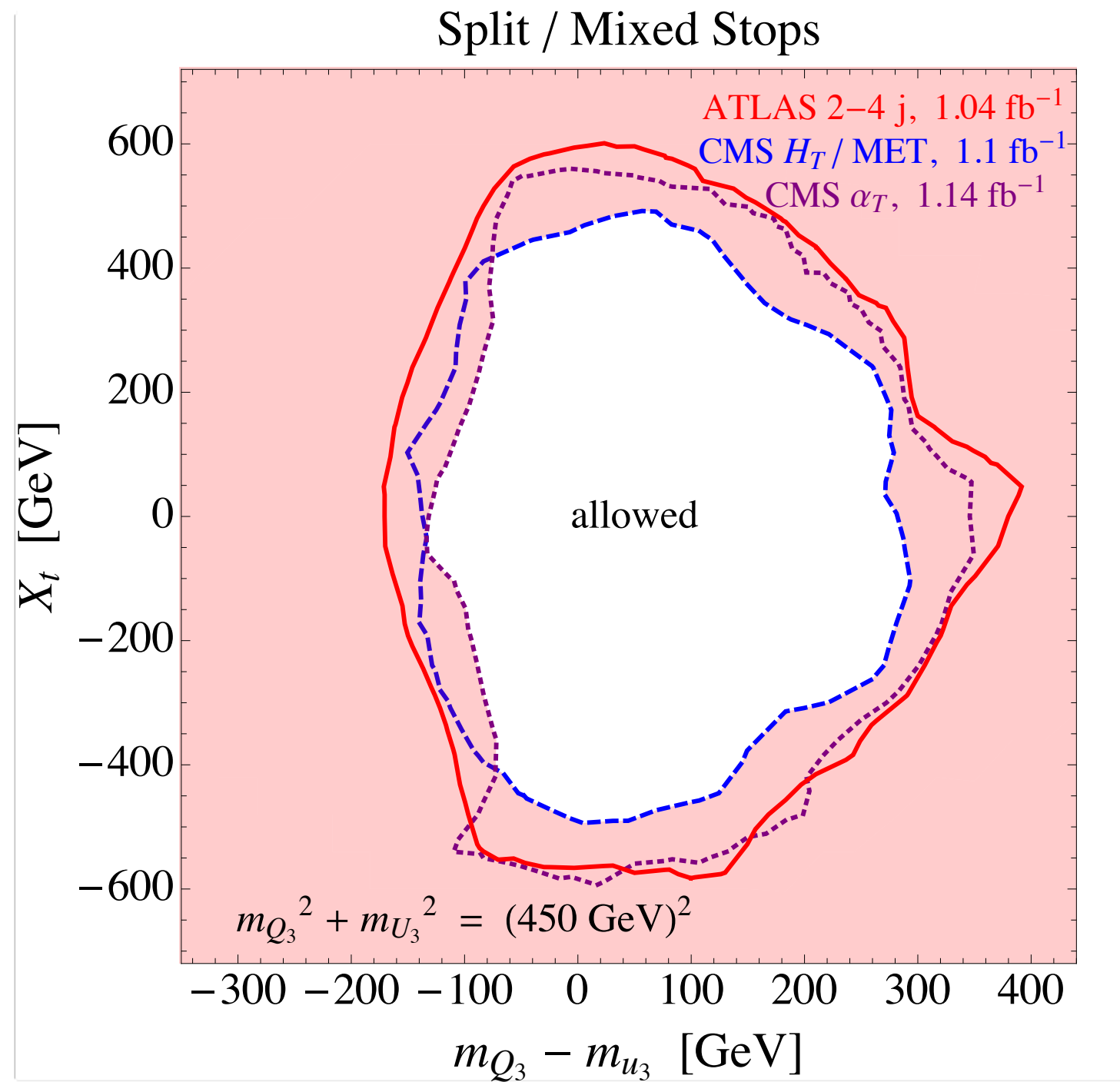
LHC surpasses Tevatron:

Strongest bounds from **jets + MET**

Un-Splitting the spectrum



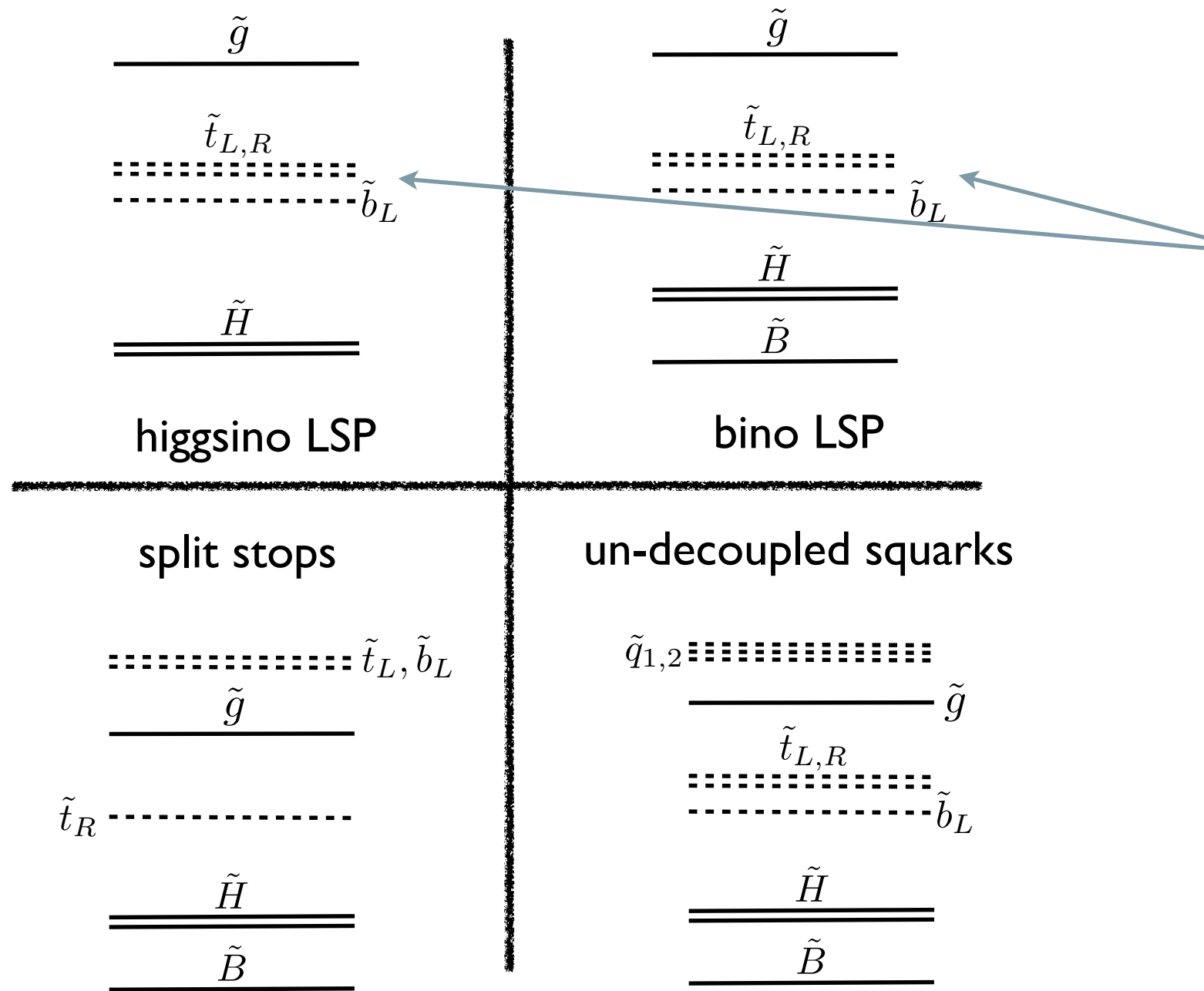
Un-Splitting the spectrum



stronger bound on the left due to light sbottom

TeVatron bounds not shown b/c they have no sensitivity for $m_{\text{LSP}} > 110 \text{ GeV}$

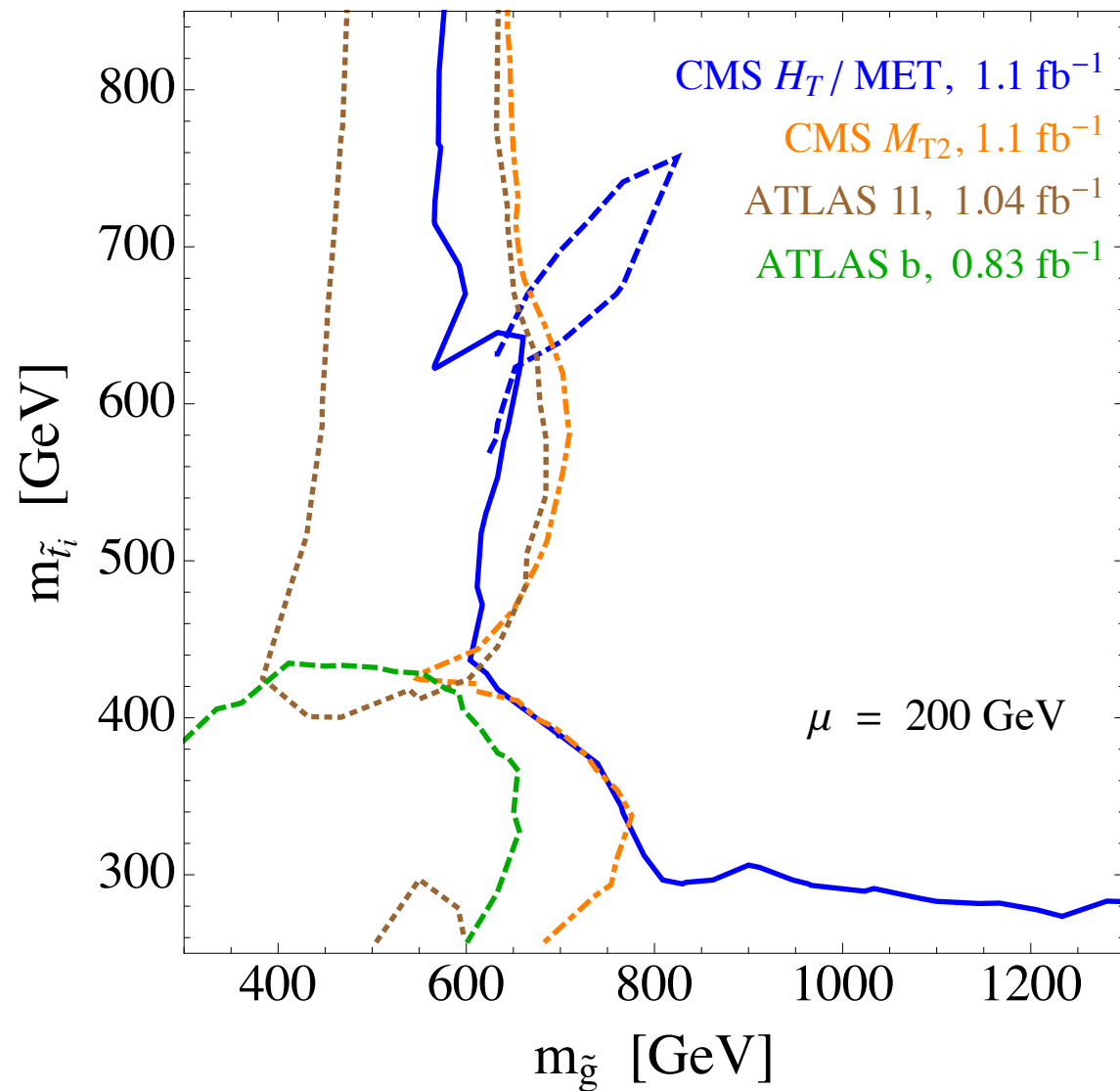
Adding gluinos



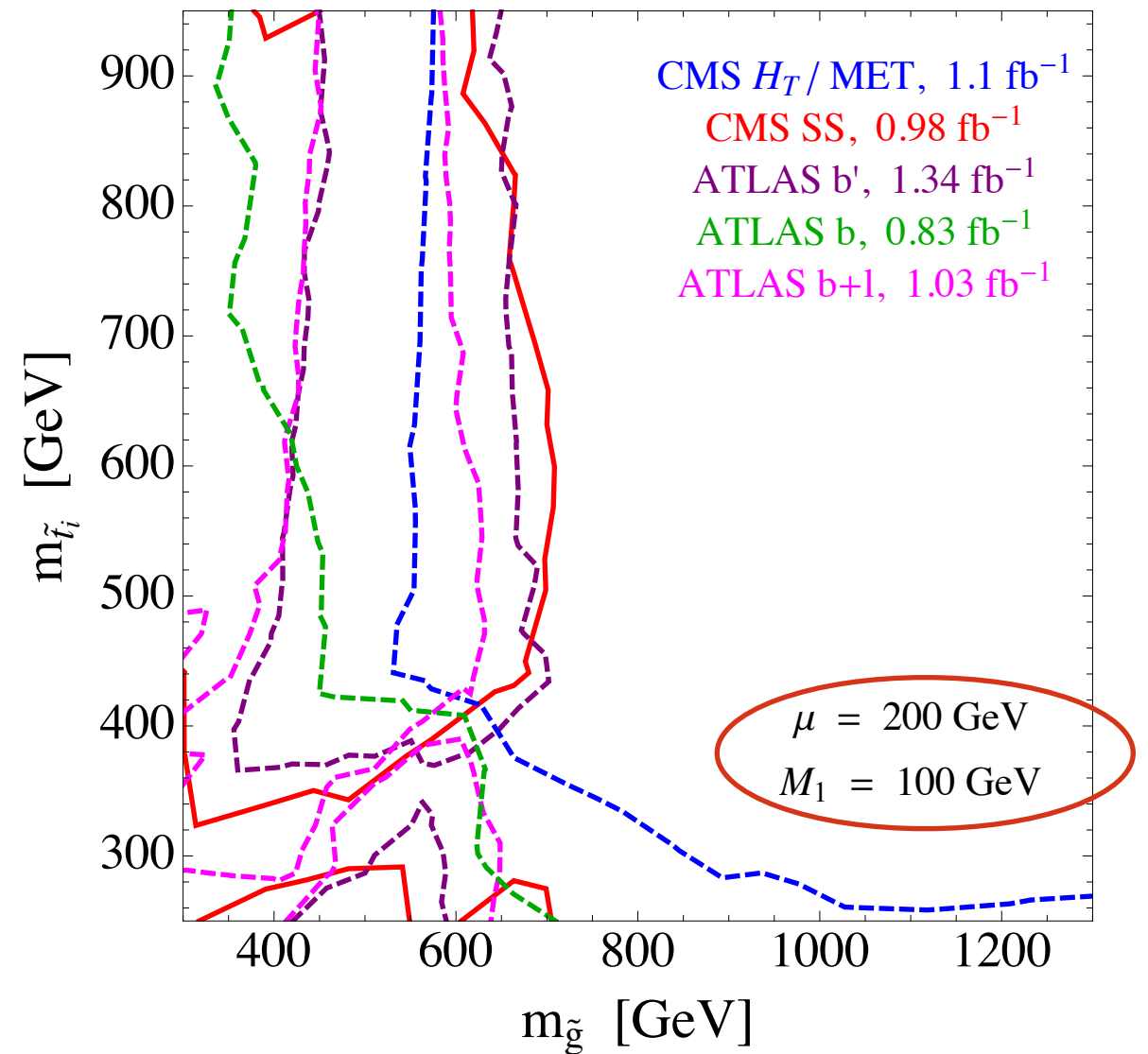
quasi-degenerate
3-rd gen'

Adding the gluinos

Higgsino LSP



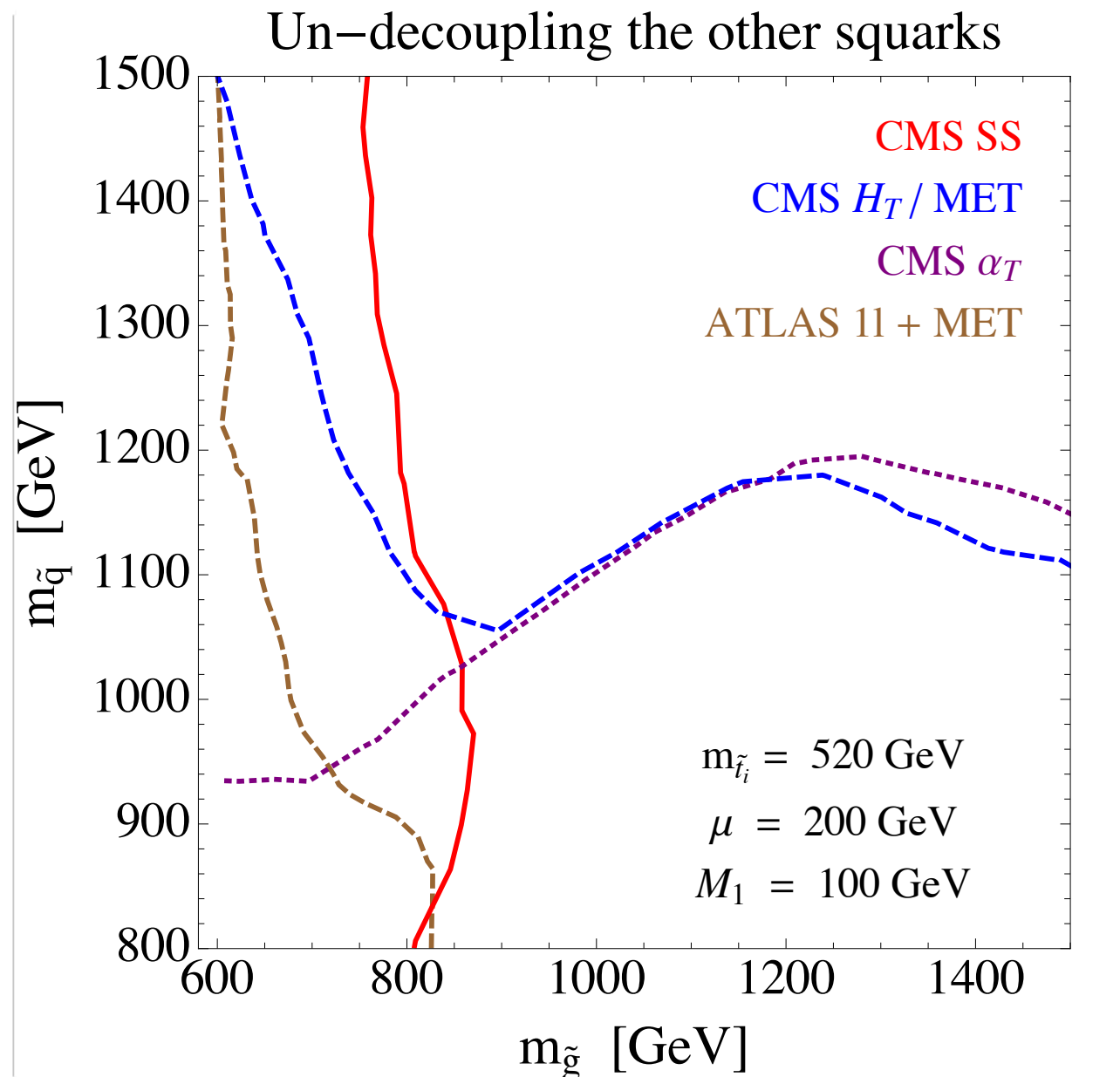
Bino LSP



Gluino bounded (again) by **jets+MET**, and **ll ν** searches

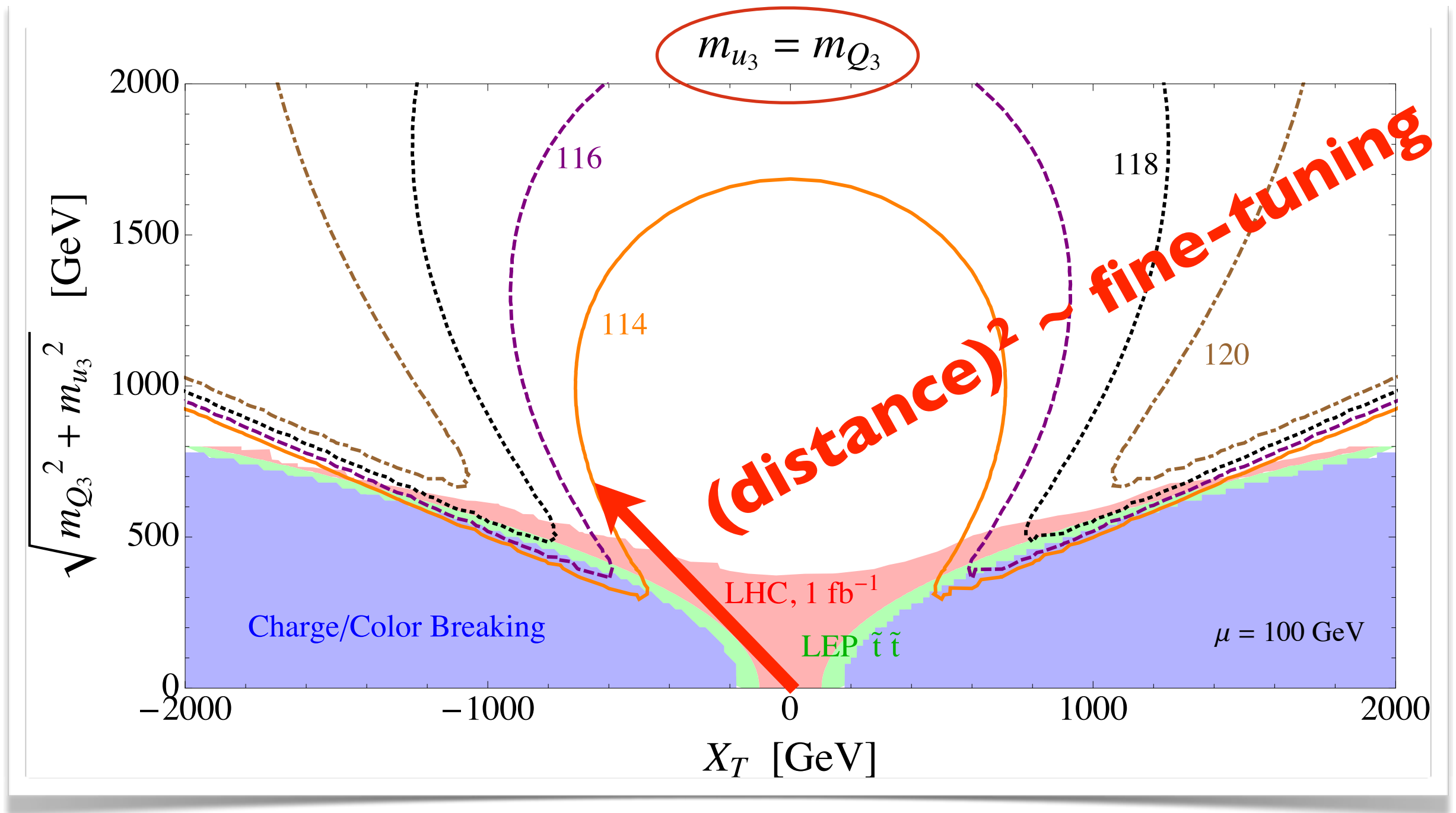
Gluino mostly bounded by **Same Sign** searches

Adding the squarks, too



- Bounds similar to the ATLAS/CMS plots (800GeV-1 TeV)
- Decoupling not effective until 1.2-1.4 TeV

MSSM higgs: LEP2 tuning vs. direct stop

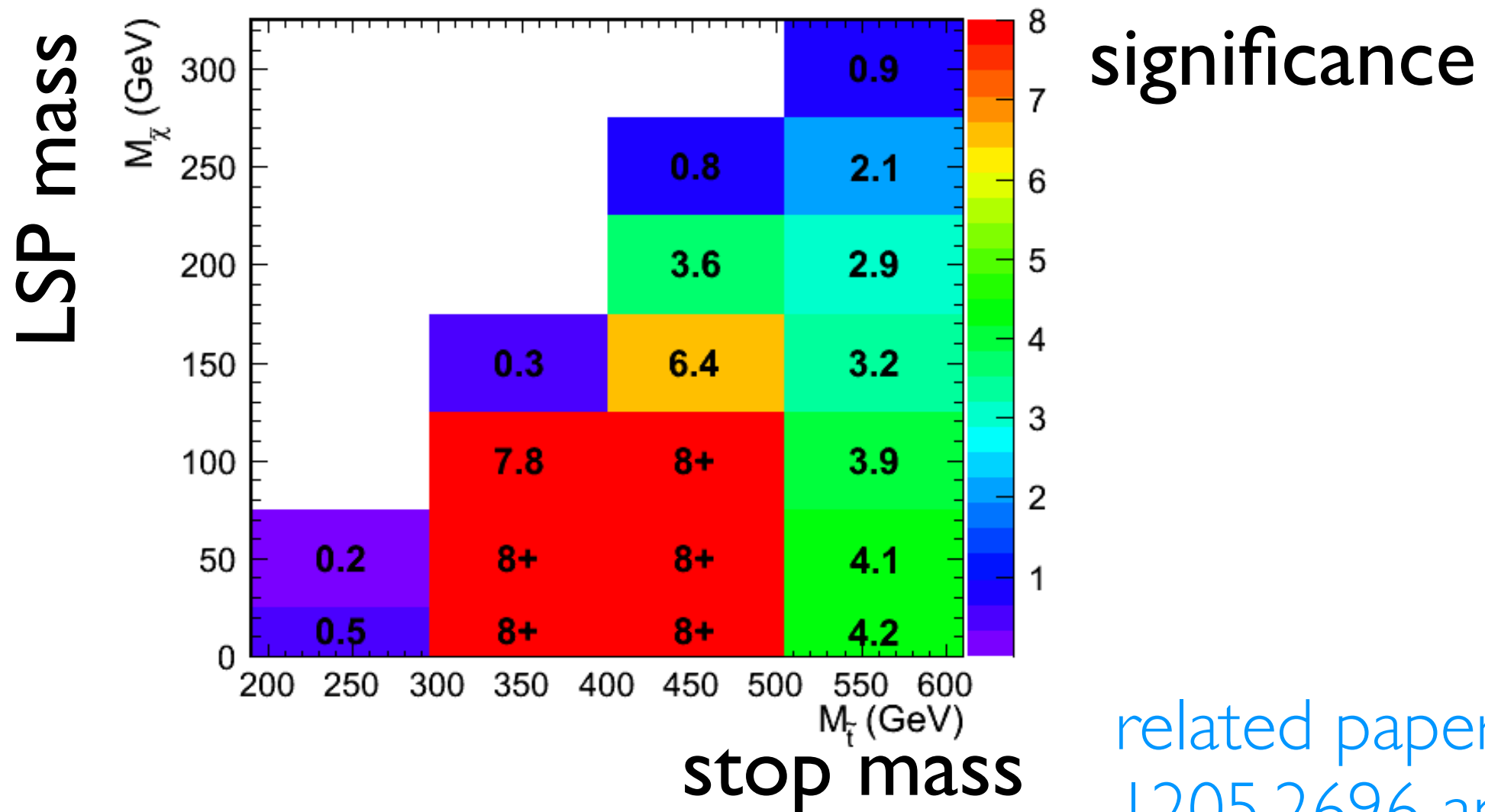


$$\delta m_H^2|_{stop} = -\frac{3}{8\pi^2} y_t^2 \left(m_{U_3}^2 + m_{Q_3}^2 + |A_t|^2 \right) \log \left(\frac{\Lambda}{\text{TeV}} \right)$$

Sensitivity at 8 TeV & 20/fb

Kaplan, Rehermann, Stolarski

$$pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow (t\chi^0) (\bar{t}\chi^0) \rightarrow (bjj\chi^0) (\bar{b}jj\chi^0)$$



related papers:
[1205.2696](#) and
 today(!): [1205.5805](#),
[1205.5808](#)

Conclusions

- Non-degenerate 1st & 2nd generation squarks poorly constrained, surprises in data? Dedicated experimental study needed.
- Next frontier: **Heavy flavor themed naturalness, EW-inos**
- Natural SUSY not in trouble yet and won't be for years to come
- Gluinos > 900 GeV, Stops $> 200-300$ GeV, Higgsinos above 100 GeV (LEP) is a completely viable spectrum