

Dark Matter at Colliders

Marco Farina

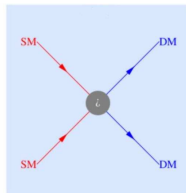
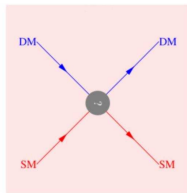
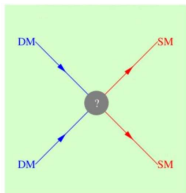
Scuola Normale Superiore & CERN

May 30 2012

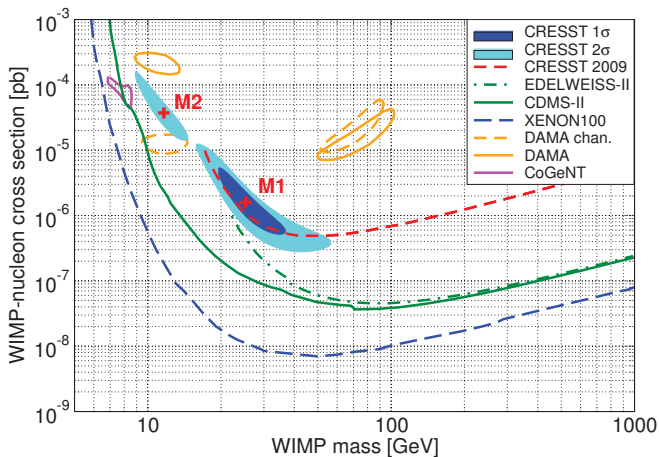
Why WIMPS?

Focusing on **WIMPS**

- Structures \rightarrow "cold" dark matter
- Wimp Miracle: $M_{DM} \sim 100 GeV$ and a typical weak scale annihilation cross section give naturally the observed Ω_{DM}
- Various detection channels:
 - Indirect (cosmic rays, ecc...): $E \sim O(10 - 100) GeV$,
 - **Direct Detection** (recoil of target nuclei): $E \sim O(1 - 10) KeV$,
 - **Production at colliders** (e.g. $pp \rightarrow \chi\chi$): $E \sim TeV$,



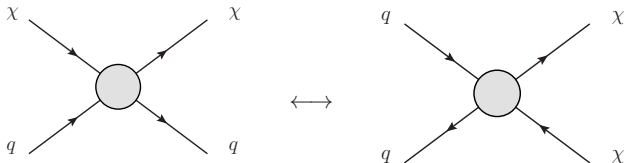
Present Direct Detection situation



CRESST-II coll. [1109.0702]

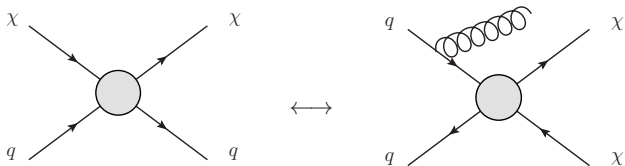
Main Idea

- If DM interacts with quarks (or leptons) and gluons both direct searches and collider production are valid



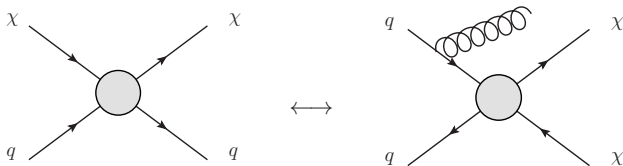
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- If DM interacts with quarks (or leptons) and gluons both direct searches and collider production are valid
- How to search for dark matter at colliders? DM(=MET) plus mono-jet or mono-photon

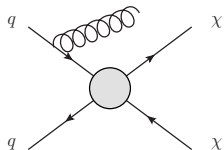


Main Idea

- If DM interacts with quarks (or leptons) and gluons both direct searches and collider production are valid
- How to search for dark matter at colliders? DM(=MET) plus mono-jet or mono-photon
- Minimal amount of assumptions (and knowledge) on the NP sector
- Easy to construct effective operators involving DM+SM



Main Idea



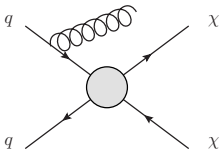
Representative sample of operators:

$$\mathcal{O}_1 = \frac{m_q}{\Lambda_1^2} (\bar{\chi}\chi) (\bar{q}q) \quad \mathcal{O}_2 = \frac{1}{\Lambda_2^3} (\bar{\chi}\chi) \left(\frac{\alpha_s}{12\pi} G^{\mu\nu} G_{\mu\nu} \right)$$

$$\mathcal{O}_3 = \frac{1}{\Lambda_3^2} (\bar{\chi}\gamma^\mu\chi) (\bar{q}\gamma_\mu q) \quad \mathcal{O}_4 = \frac{1}{\Lambda_4^2} (\bar{\chi}\gamma^\mu\gamma^5\chi) (\bar{q}\gamma_\mu\gamma^5 q)$$

E.g. s-channel exchange $\Lambda \sim M/\sqrt{g_q g_\chi}$

Main Idea



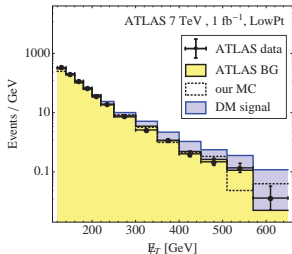
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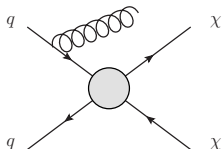
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E.g. s-channel exchange $\Lambda \sim M/\sqrt{g_q g_\chi}$

	ATLAS LowPT 1.0 fb ⁻¹	ATLAS HighPT 1.0 fb ⁻¹	ATLAS veryHighPT 1.0 fb ⁻¹
Expected	15100 ± 700	1010 ± 75	193 ± 25
Observed	15740	965	167



Main Idea

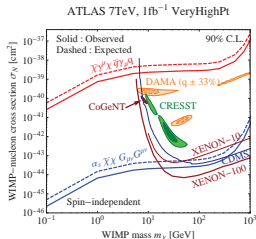
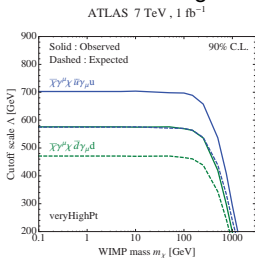


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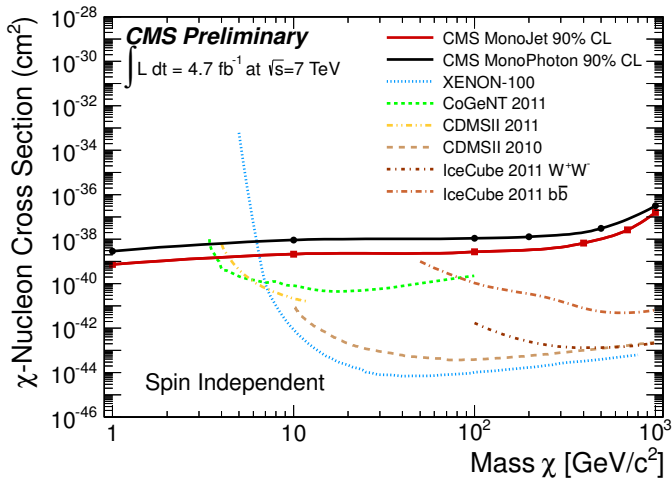
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E.g. s-channel exchange $\Lambda \sim M/\sqrt{g_q g_\chi}$



Fox et al. [1109.4398]

Recent results (CMS)



Uncertainties and Issues

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(independence between each other and possible cancellations)

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D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\mu\nu}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$
D15	$\bar{\chi}\sigma^{\mu\nu}\chi F_{\mu\nu}$	M
D16	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi F_{\mu\nu}$	D
M1	$\bar{\chi}\chi\bar{q}q$	$m_q/2M_*^3$
M2	$\bar{\chi}\gamma^5\chi\bar{q}q$	$im_q/2M_*^3$

Name	Operator	Coefficient
M3	$\bar{\chi}\chi\bar{q}\gamma^5q$	$im_q/2M_*^3$
M4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	$m_q/2M_*^3$
M5	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/2M_*^2$
M6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/2M_*^2$
M7	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/8M_*^3$
M8	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/8M_*^3$
M9	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^3$
M10	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/8M_*^3$
C1	$\chi^\dagger\chi\bar{q}q$	m_q/M_*^2
C2	$\chi^\dagger\chi\bar{q}\gamma^5q$	im_q/M_*^2
C3	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu q$	$1/M_*^2$
C4	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu\gamma^5q$	$1/M_*^2$
C5	$\chi^\dagger\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^2$
C6	$\chi^\dagger\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^2$
R1	$\chi^2\bar{q}q$	$m_q/2M_*^2$
R2	$\chi^2\bar{q}\gamma^5q$	$im_q/2M_*^2$
R3	$\chi^2 G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/8M_*^2$
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Goodman et Al. [1009.0008]

Uncertainties and Issues

- Number of operators
(independence between each other and possible cancellations)
- Perturbativity ($\Lambda \gtrsim 2\pi m_\chi$), light mediators, ecc...

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D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\mu\nu}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$
D15	$\bar{\chi}\sigma^{\mu\nu}\chi F_{\mu\nu}$	M
D16	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi F_{\mu\nu}$	D
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M5	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/2M_*^2$
M6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/2M_*^2$
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- QCD and background uncertainties

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- QCD and background uncertainties
- Translation and comparison on the Direct Detection plane (reintroduces atomic and astro uncertainties)

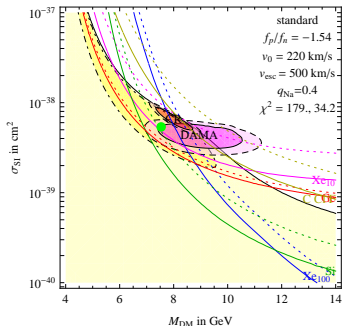
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M5	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/2M_*^2$
M6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/2M_*^2$
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M8	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/8M_*^3$
M9	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^3$
M10	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/8M_*^3$
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C2	$\chi^\dagger\chi\bar{q}\gamma^5q$	im_q/M_*^2
C3	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu q$	$1/M_*^2$
C4	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu\gamma^5q$	$1/M_*^2$
C5	$\chi^\dagger\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^2$
C6	$\chi^\dagger\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^2$
R1	$\chi^2\bar{q}q$	$m_q/2M_*^2$
R2	$\chi^2\bar{q}\gamma^5q$	$im_q/2M_*^2$
R3	$\chi^2 G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/8M_*^2$
R4	$\chi^2 G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^2$

Goodman et Al. [1009.0008]

Three good examples I: isospin violating couplings

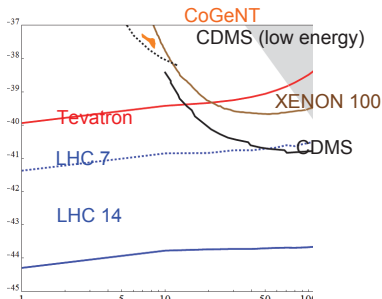
- DAMA and CoGeNT can be accommodated by different couplings to u and d quarks
- It is required $f_p/f_n = -1.54$



MF et al. [1107.0715]

Three good examples I: isospin violating couplings

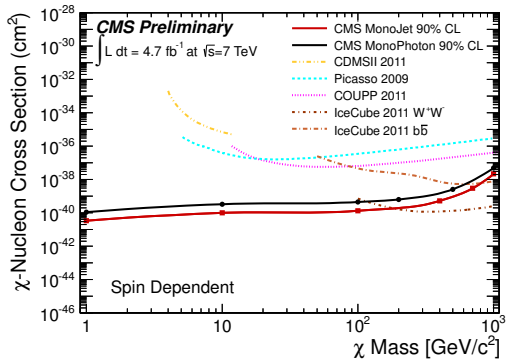
- DAMA and CoGeNT can be accommodated by different couplings to u and d quarks
- It is required $f_p/f_n = -1.54$



Rajaraman et al. [1108.1196]

Three good examples II: spin dependent scattering

Collider bounds do not suffer from spin-spin suppression

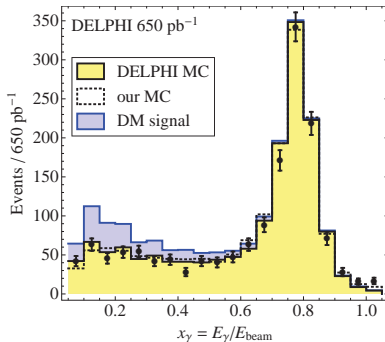
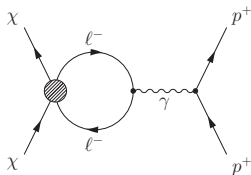


Three good examples III: LEP(tons)

- Yet another way to alleviate the tension between experiments
- Coupling to leptons only e.g.

$$\mathcal{O}_l = \frac{1}{\Lambda^2} (\bar{\chi}\chi) (\bar{l}l)$$

- LEP searches for $\gamma + MET$



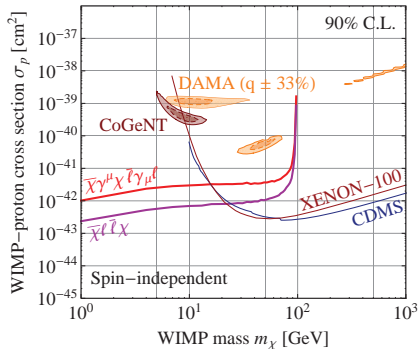
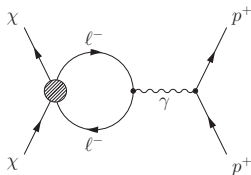
Fox et al. [1103.0204]

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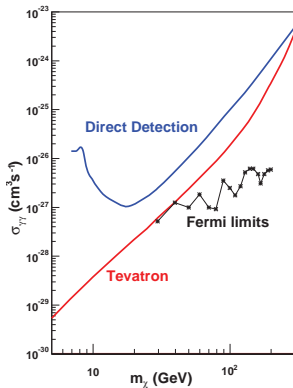
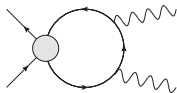


Fox et al. [1103.0204]

FERMI and Colliders?

Comparison is possible also with Indirect Detection experiments.

For example FERMI data ($\chi\chi \rightarrow \gamma\gamma$, $\chi\chi \rightarrow \gamma Z$)

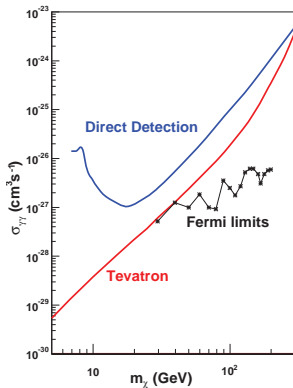
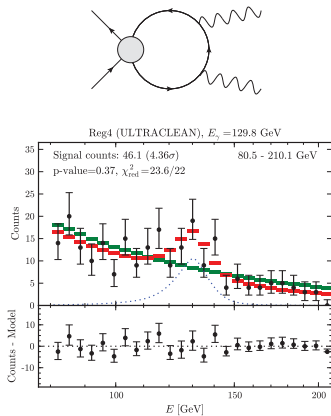


Goodman et al. [1009.0008]

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Weniger [1204.2797]

Goodman et al. [1009.0008]

Conclusions

- Comparison of different types of experiments is possible with a minimal set of assumptions
- Effective Operators are powerful (if handled with care)
- Is it all in the hands of experimental collaborations?
No! We just have to scratch our heads more (new signals, new channels, old colliders)

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- Questions? Ask Tim on Friday