

Measurements of heavy flavor properties at ATLAS and CMS

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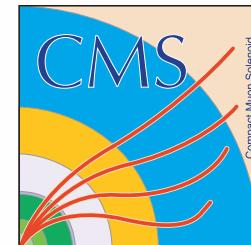
on behalf of the CMS and ATLAS collaborations

Blois 2014: 26th Rencontres de Blois on “Particle Physics and Cosmology”

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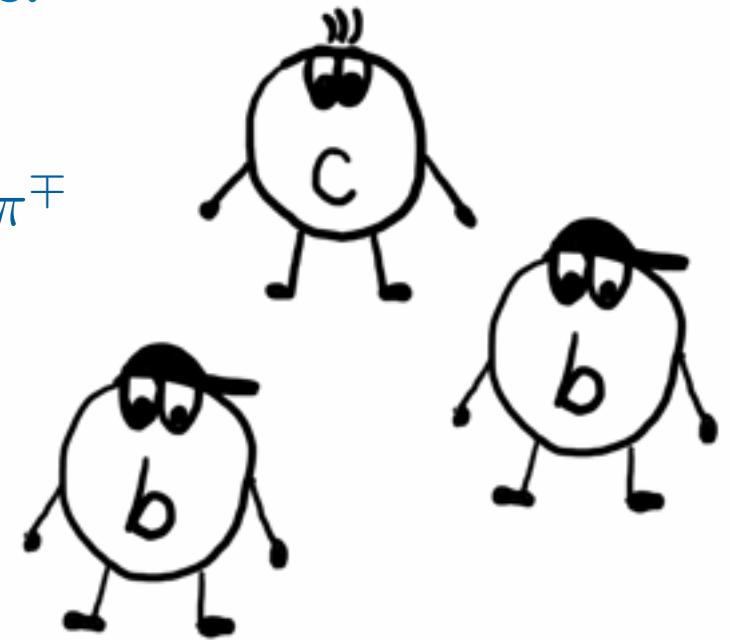
HEPHY
Institute of High Energy Physics



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Outline

- ATLAS and CMS have a rich program in heavy flavor physics
- Heavy flavor production was already covered in Konstantin Toms' talk
- I will concentrate on a few recent topics:
 - Measurement of α_b in $\Lambda_b^0 \rightarrow J/\psi \Lambda_0$
 - $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp$
 - Quarkonium polarization



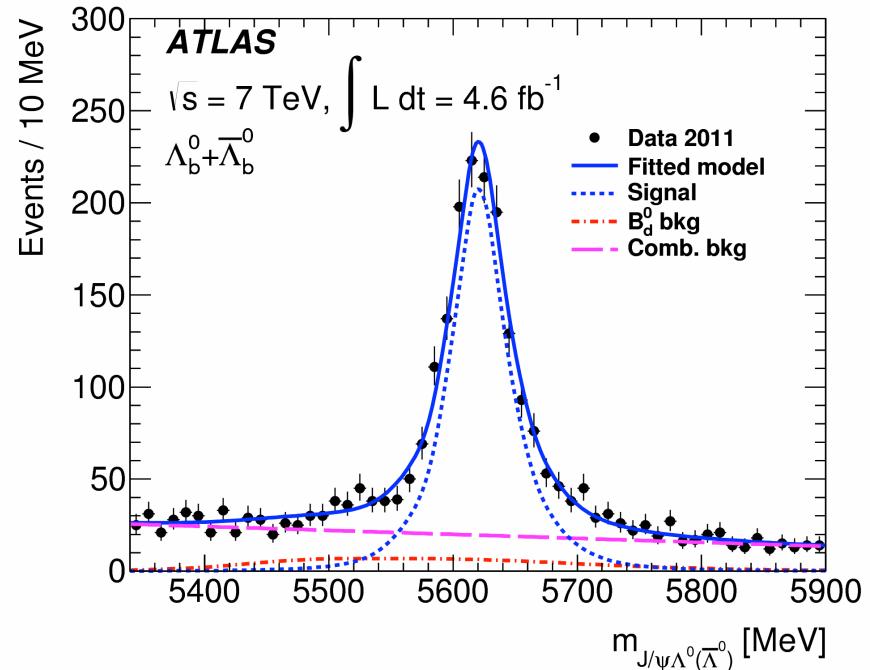
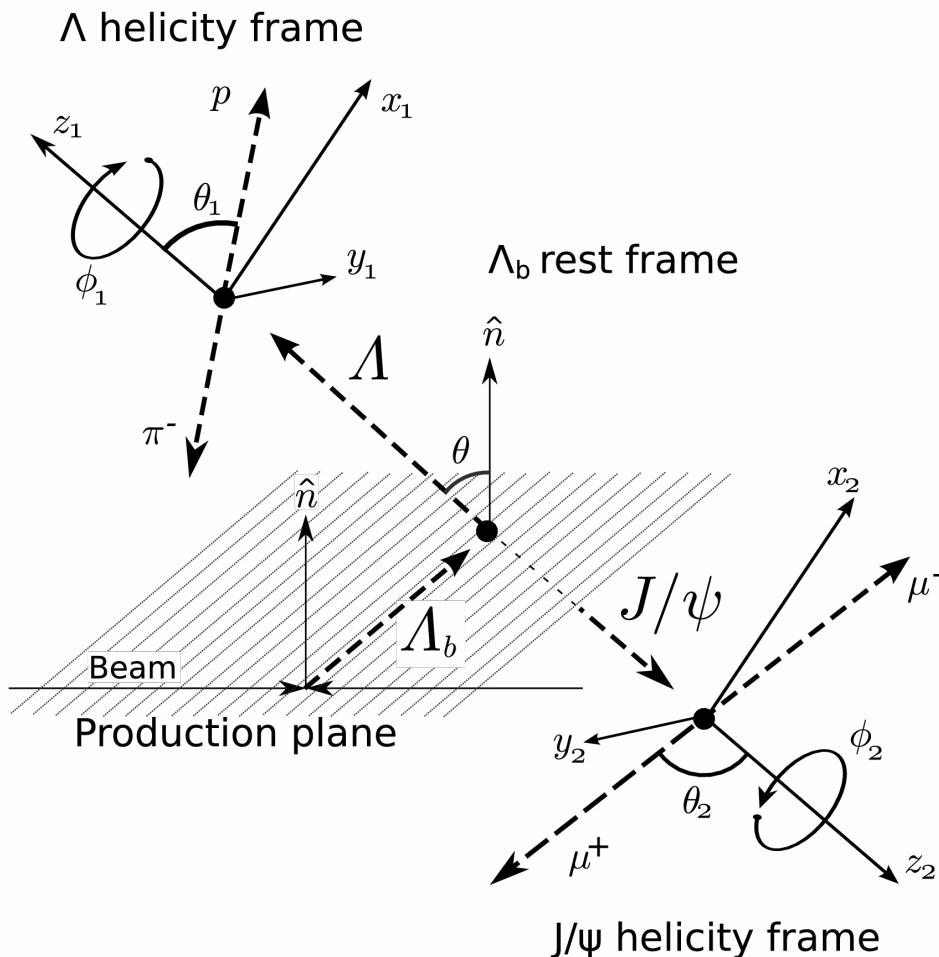
$$\alpha_b \text{ in } \Lambda_b^0 \rightarrow J/\psi \Lambda_0$$

- Parity violation in the hadronic sector depends on the constituents of the hadron because of the presence of strongly bound spectator quarks
- Measurement of the parity-violating asymmetry parameter α_b provides a test for several theoretical models:
 - Perturbative quantum chromodynamics (pQCD): $\alpha_b = -0.17$ to -0.14 (PRD 65, 074030 (2002))
 - Heavy quark effective theory (HQET): $\alpha_b = 0.78$ (PLB 614, 165 (2005))
- LHCb result compatible with pQCD calculation:
 $\alpha_b = 0.05 \pm 0.17 \text{ (stat.)} \pm 0.07 \text{ (syst.)}$
(PLB 724 (2013), 27)
- ATLAS recently performed an independent measurement



$\Lambda_b^0 \rightarrow J/\psi \Lambda_0$

- Decay angles
 $\Omega = \{\theta, \phi, \theta_1, \phi_1, \theta_2, \phi_2\}$



- Assume CP conservation
- Combination of Λ_b^0 and $\bar{\Lambda}_b^0$ baryons



Details in arXiv:1404.1071

Measurement of α_b in $\Lambda_b^0 \rightarrow J/\psi \Lambda_0$

- Parameters extracted in a least-squares fit to the distributions of decay angles
- Parity-violating asymmetry parameter

$$\alpha_b = 0.30 \pm 0.16(\text{stat.}) \pm 0.06(\text{syst.})$$

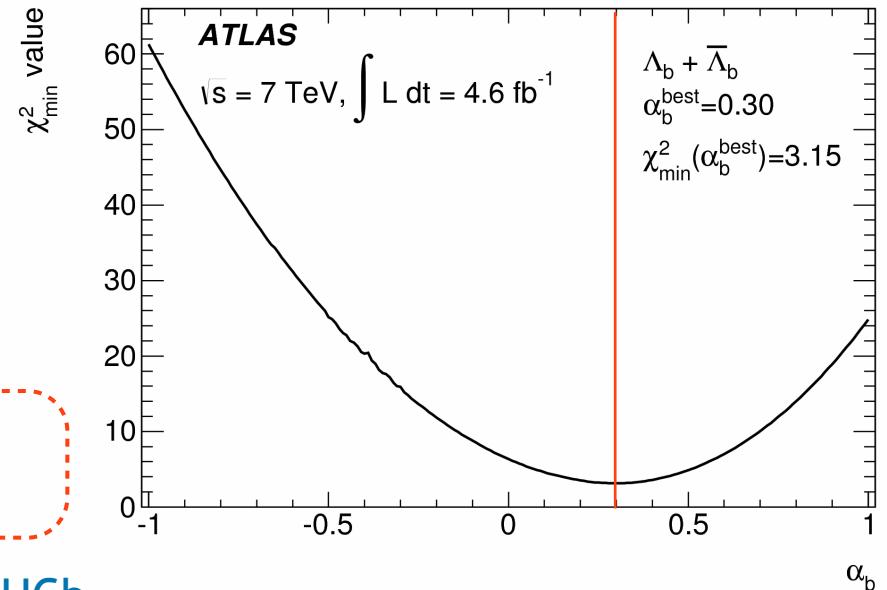
- Consistent with measurements from LHCb
- Inconsistent with pQCD and HQET at a level of 2.6 and 2.8 standard deviations
- Helicity amplitudes $A(\lambda_\Lambda, \lambda_{J/\psi})$

$$A(1/2, 0) \equiv |a_+| = 0.17^{+0.12}_{-0.17}(\text{stat.}) \pm 0.09(\text{syst.})$$

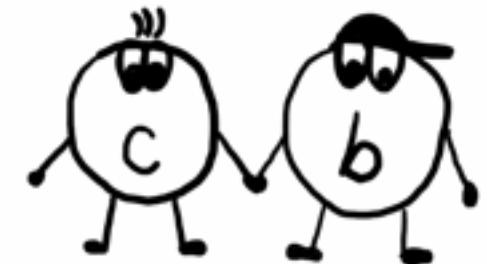
$$A(-1/2, 0) \equiv |a_-| = 0.59^{+0.06}_{-0.07}(\text{stat.}) \pm 0.03(\text{syst.})$$

$$A(-1/2, -1) \equiv |b_+| = 0.79^{+0.04}_{-0.05}(\text{stat.}) \pm 0.02(\text{syst.})$$

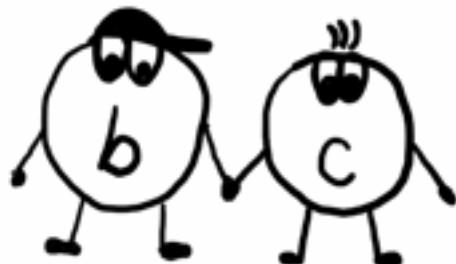
$$A(1/2, 1) \equiv |b_-| = 0.08^{+0.13}_{-0.08}(\text{stat.}) \pm 0.06(\text{syst.})$$



B_c^\pm meson

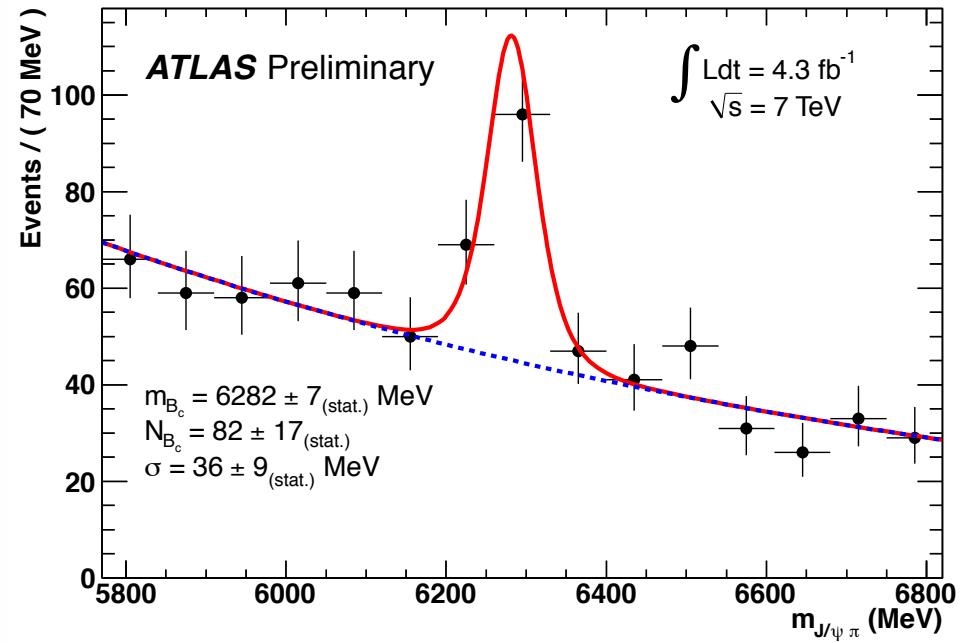


- First seen at CDF through the semileptonic decay $B_c^\pm \rightarrow J/\psi l^+ \nu$
- Ground state of the $\bar{b}c$ ($b\bar{c}$) system
- Unique in that any of its quarks can decay weakly, leaving the other as a spectator
- Offers possibility to study the heavy quark dynamics that is inaccessible through the investigations on the $b\bar{b}$ and $c\bar{c}$ quarkonia



$B_c^\pm \rightarrow J/\psi \pi^\pm$

- ATLAS observed B_c^\pm mesons through their decay to $J/\psi \pi^\pm$ for $p_T > 15$ GeV
- The yield is extracted using an unbinned maximum likelihood fit to the mass distribution



yield = 82 ± 17 (stat.)events
 $m(B_c^\pm) = 6282 \pm 7$ (stat.)MeV

- Consistent with PDG average: $m(B_c^\pm) = 6277 \pm 6$ MeV

Details in ATLAS-CONF-2012-028



$B_c^\pm \rightarrow J/\psi \pi^\pm$ and $B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp$

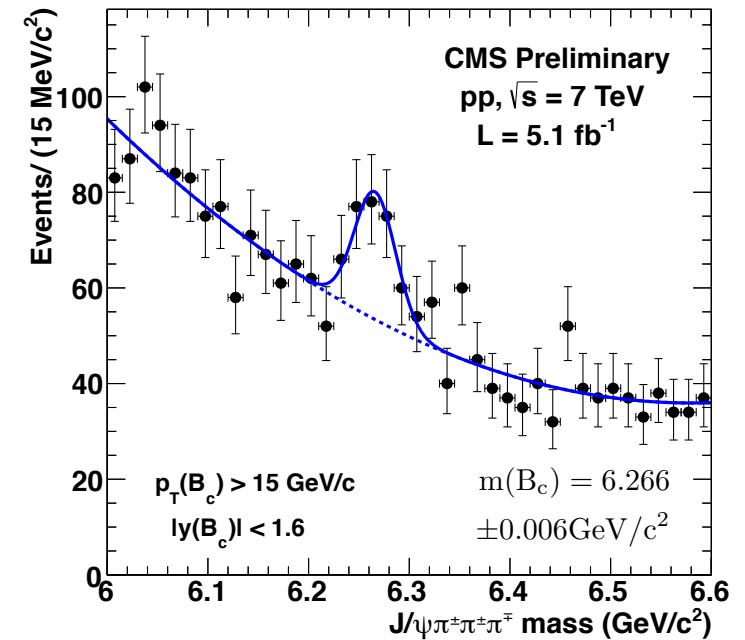
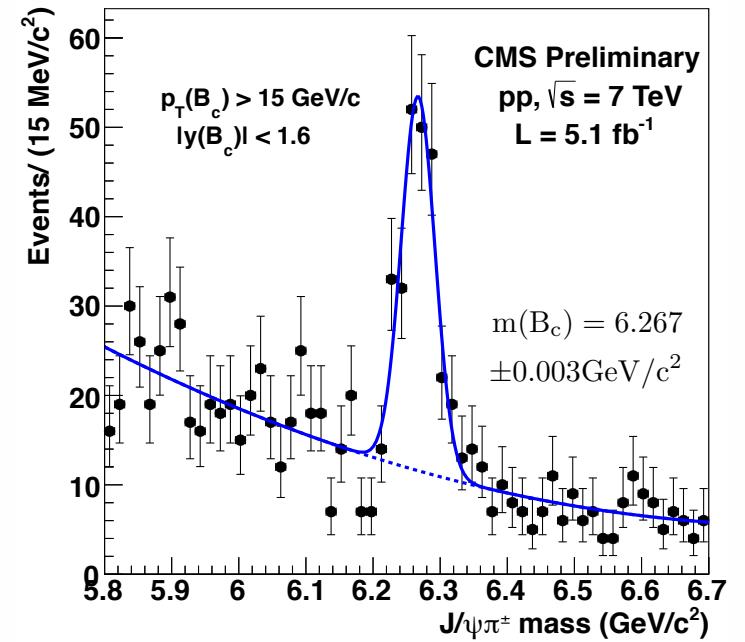
- CMS studied the $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $J/\psi \pi^\pm \pi^\pm \pi^\mp$ decay modes for $p_T > 15 \text{ GeV}$ and $|y| < 1.6$

$$\frac{\sigma(B_c^\pm) \times \text{Br}(B_c^\pm \rightarrow J/\psi \pi^\pm)}{\sigma(B^\pm) \times \text{Br}(B^\pm \rightarrow J/\psi K^\pm)} = (0.48 \pm 0.05 \text{ (stat.)} \pm 0.04 \text{ (syst.)})_{-0.03}^{+0.05} (\tau_{B_c}) \times 10^{-2}$$

$$\frac{\text{Br}(B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp)}{\text{Br}(B_c^\pm \rightarrow J/\psi \pi^\pm)} = 2.43 \pm 0.76 \text{ (stat.)}_{-0.44}^{+0.46} \text{ (syst.)}$$

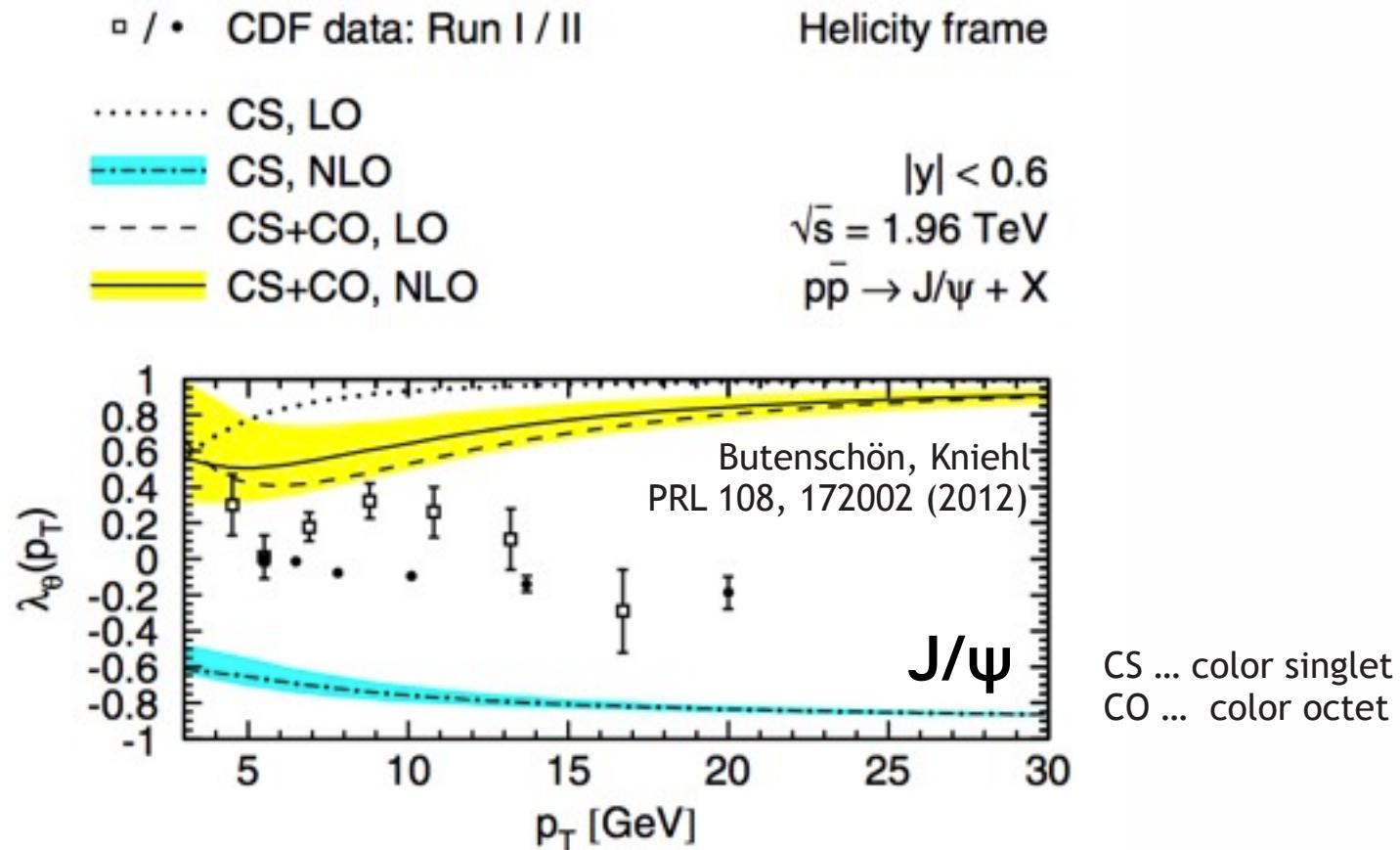
- Branching fraction ratio in good agreement with LHCb result:
 $2.41 \pm 0.30 \text{ (stat.)} \pm 0.33 \text{ (syst.)}$
 (PRL 109 (2012), 232001)

Details in CMS-PAS-BPH-12-011



Quarkonium polarization

- Long standing puzzle: No theoretical model has explained experimental measurements of quarkonium polarization
- Pre-LHC measurements showed inconsistencies in the determination of the polarization



Quarkonium polarization

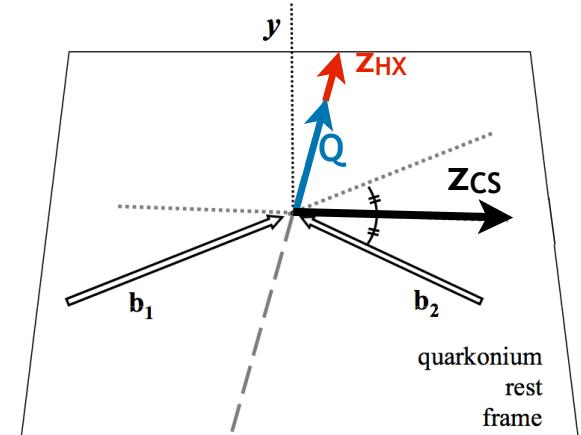
- Polarization is measured through the average angular decay distribution - for vector particles most generally written as

$$W(\cos \vartheta, \varphi | \vec{\lambda}) = \frac{3/(4\pi)}{(3 + \lambda_\vartheta)} (1 + \lambda_\vartheta \cos^2 \vartheta + \lambda_\varphi \sin^2 \vartheta \cos 2\varphi + \lambda_{\vartheta\varphi} \sin 2\vartheta \cos \varphi)$$

where $\lambda_\vartheta, \lambda_\varphi, \lambda_{\vartheta\varphi}$ are the polarization parameters

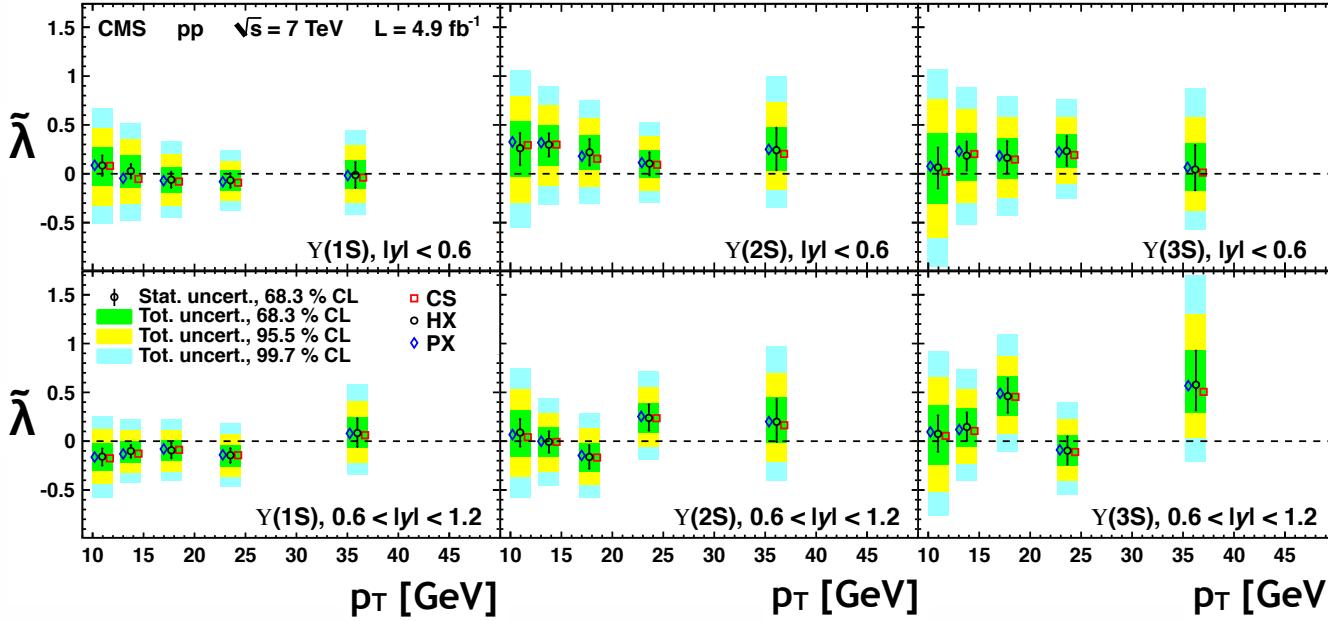
- CMS measured $\lambda_\vartheta, \lambda_\varphi, \lambda_{\vartheta\varphi}$ and the frame invariant parameter $\tilde{\lambda} = (\lambda_\vartheta + 3\lambda_\varphi)/(1 - \lambda_\varphi)$ for J/ψ , $\psi(2S)$, $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$ mesons decaying to $\mu^+\mu^-$ three different reference frames:

- center-of-mass helicity HX (polar axis z_{HX}
 \approx direction of quarkonium momentum)
- Collins-Soper CS ($z_{CS} \approx$ direction of relative velocity of colliding particles)
- perpendicular helicity PX ($z_{PX} \perp z_{CS}$)

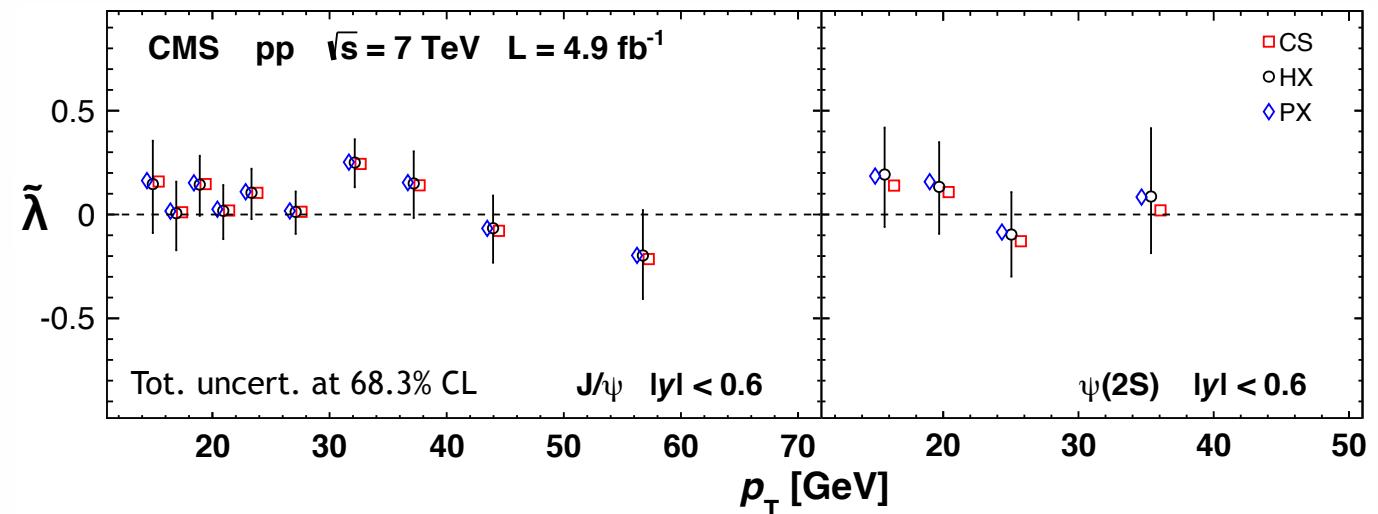
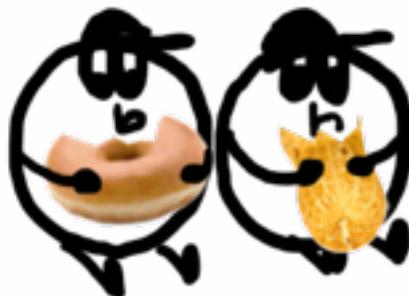


Details in PRL 110 (2013), 081802 and PLB 727 (2013), 381

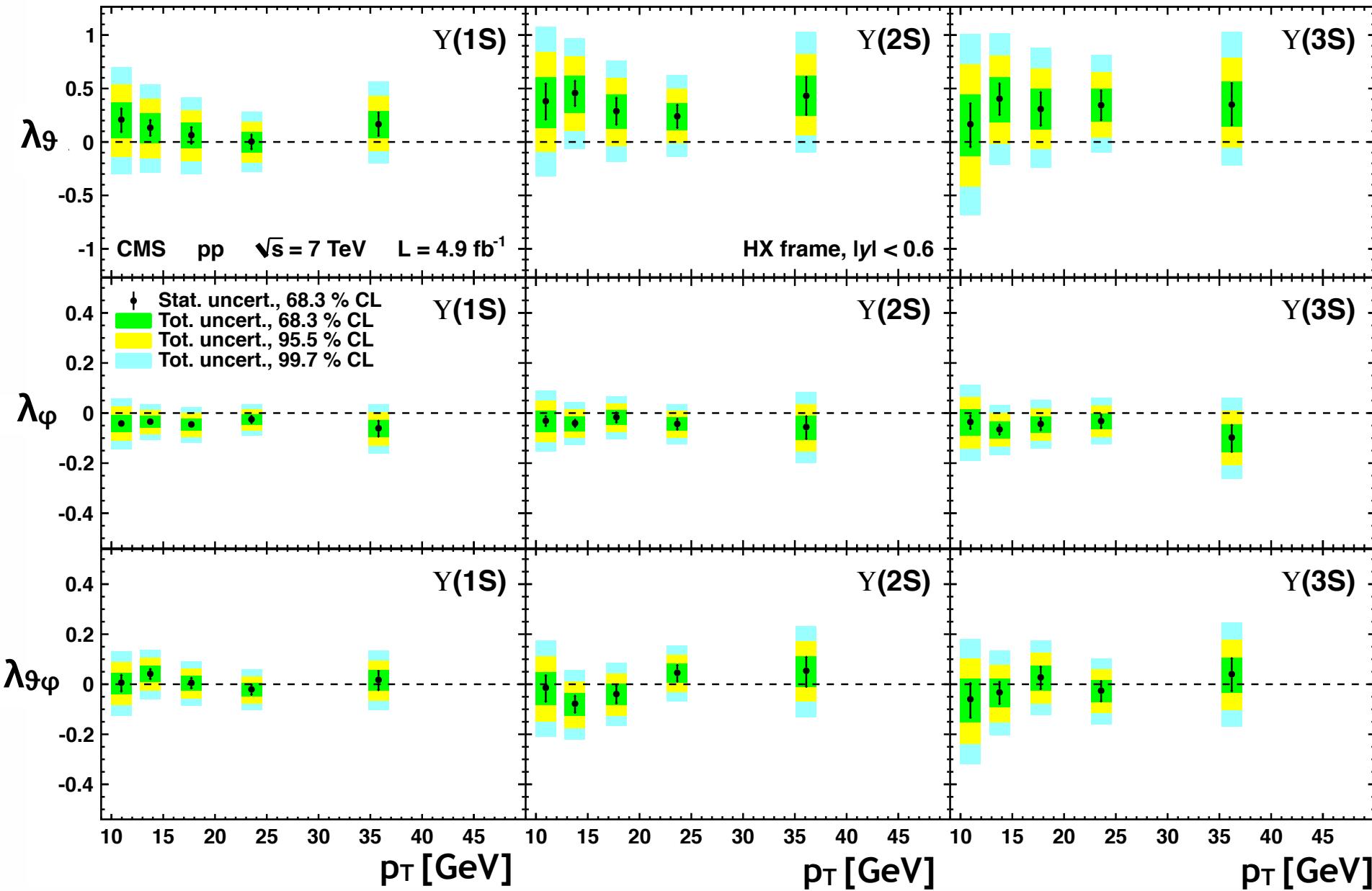
Frame Invariant Parameter $\tilde{\lambda}$



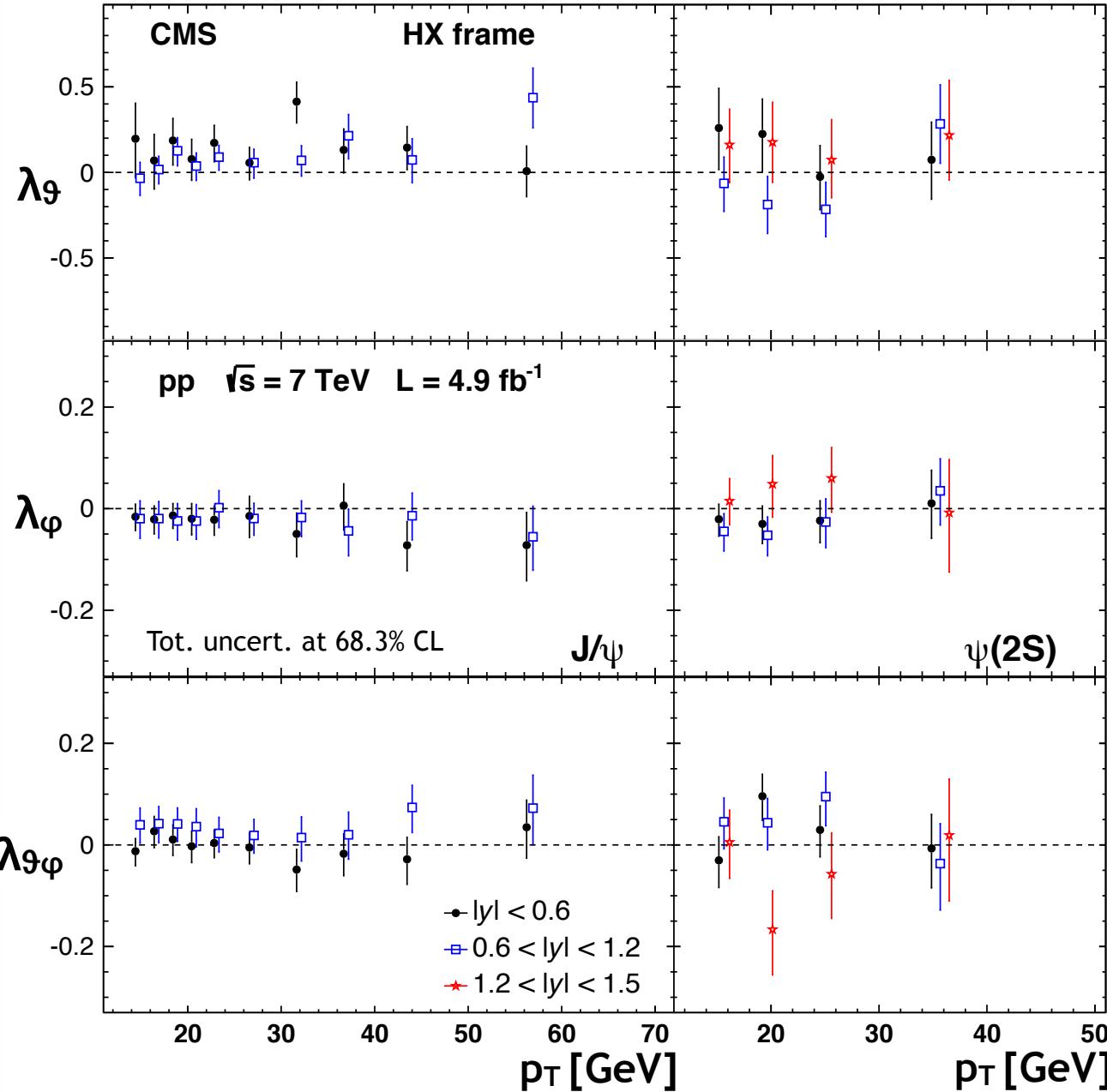
Good agreement between the $\tilde{\lambda}$ parameters in the three reference frames shows that the results are consistent



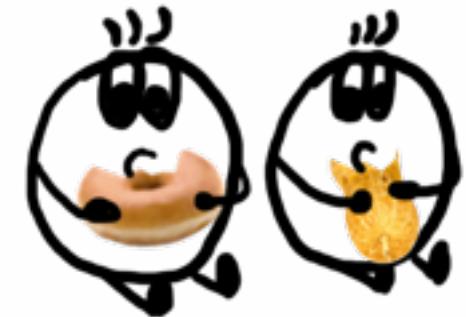
$\Upsilon(nS)$ Polarization in the HX Frame, $|y| < 0.6$



Prompt $\psi(nS)$ Polarization in the HX Frame

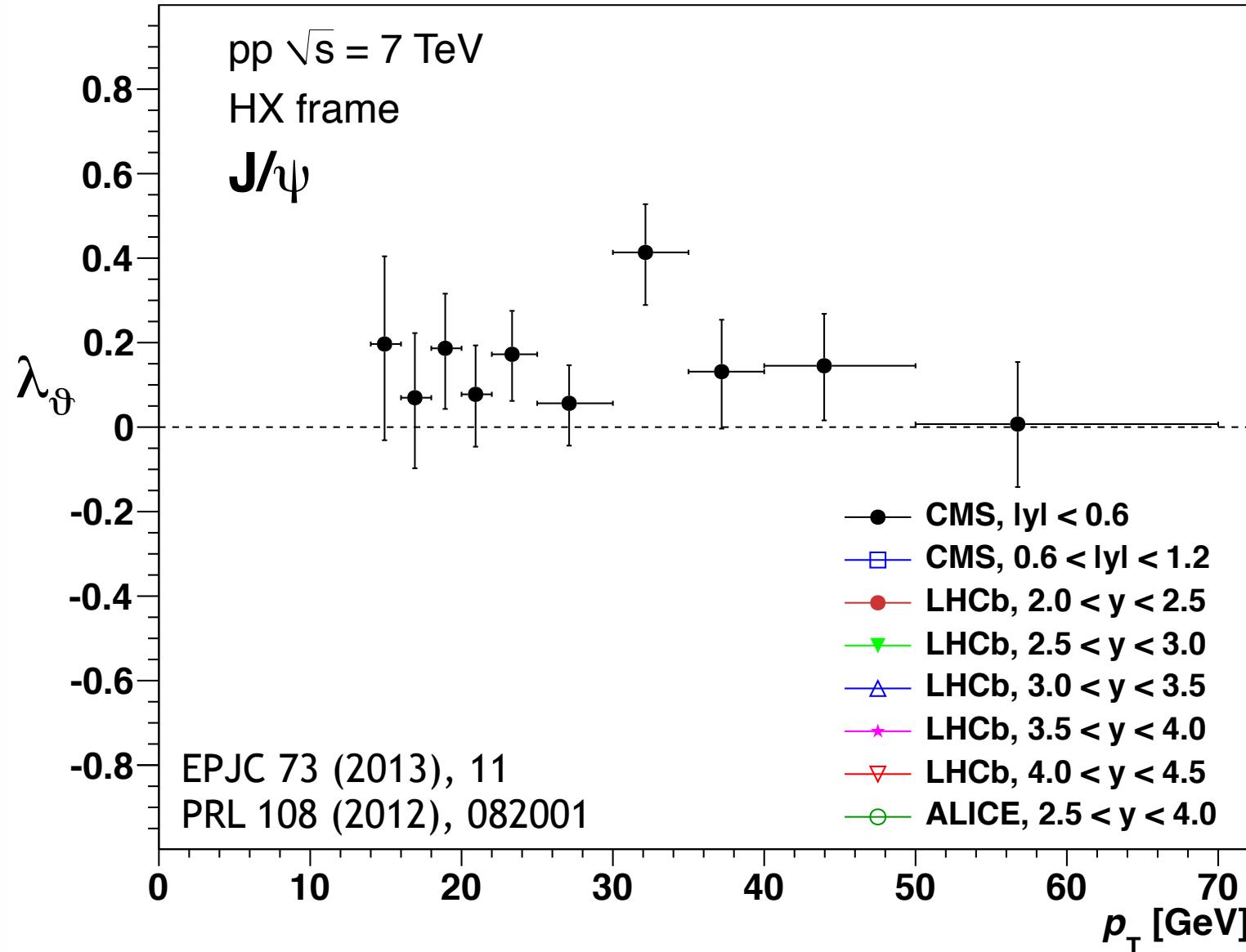


- $\psi(2S)$ is not affected by feed-down decays from higher states
- No sign of strong polarization



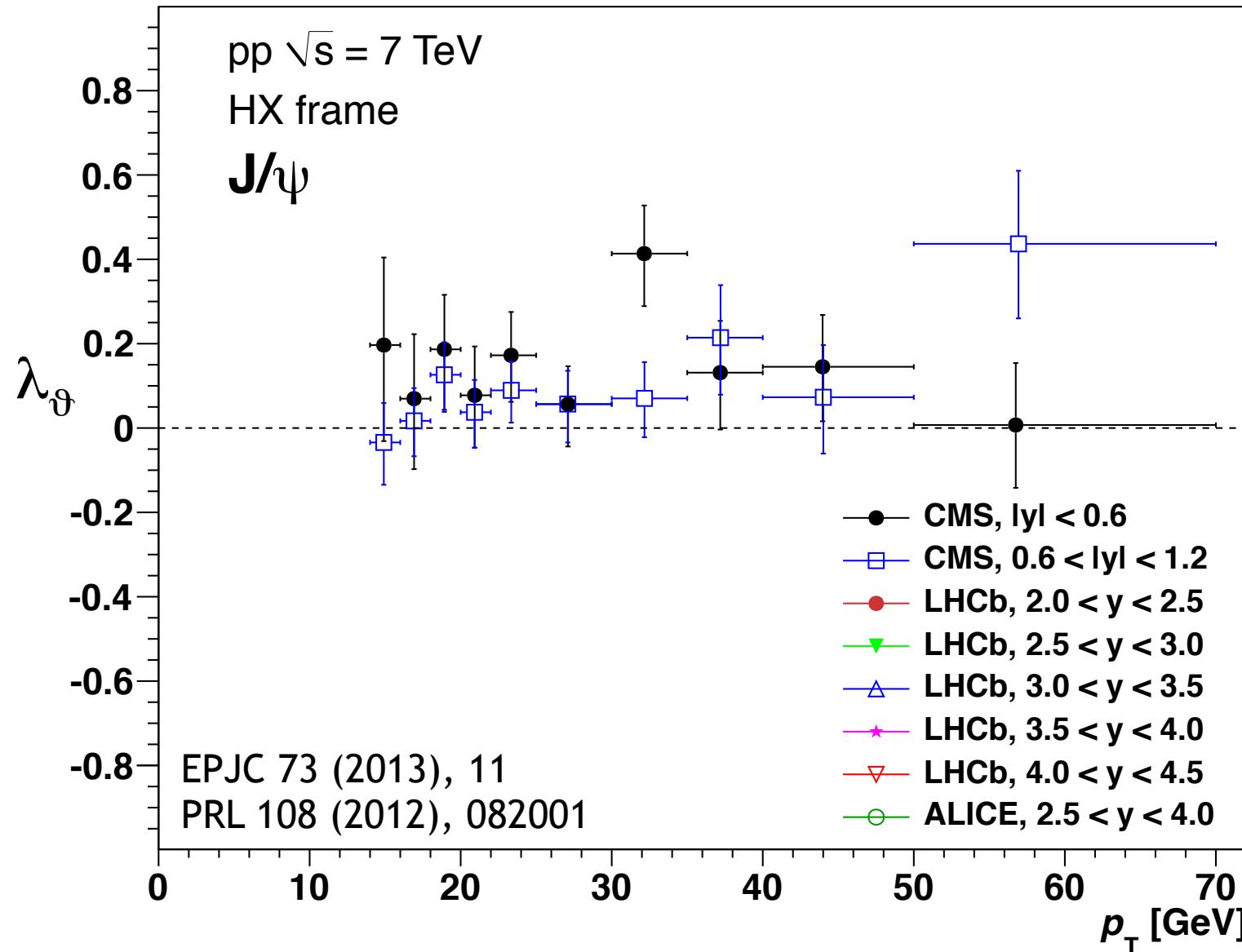
Comparison to Other LHC Experiments: J/ ψ

All LHC results are compatible with each other.



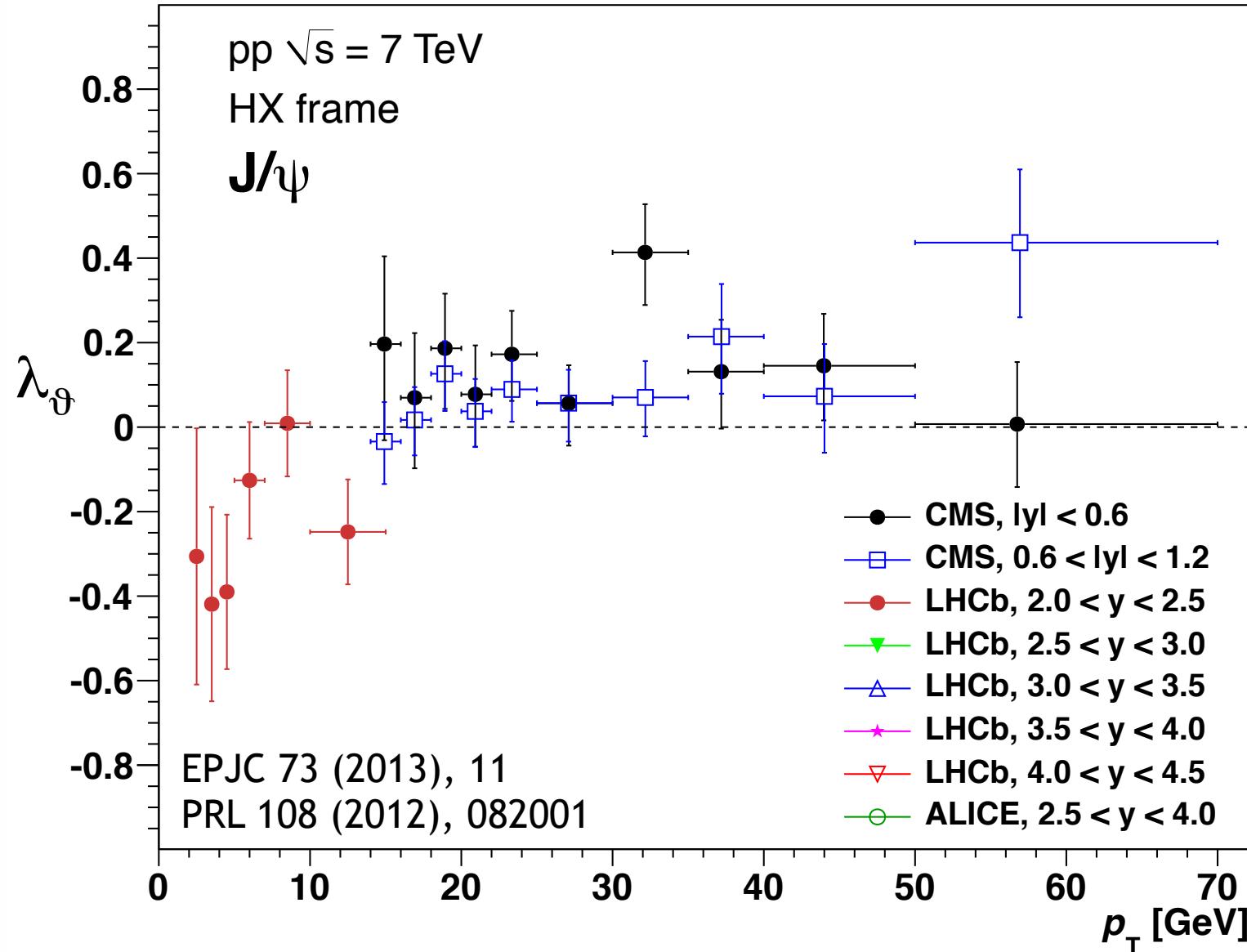
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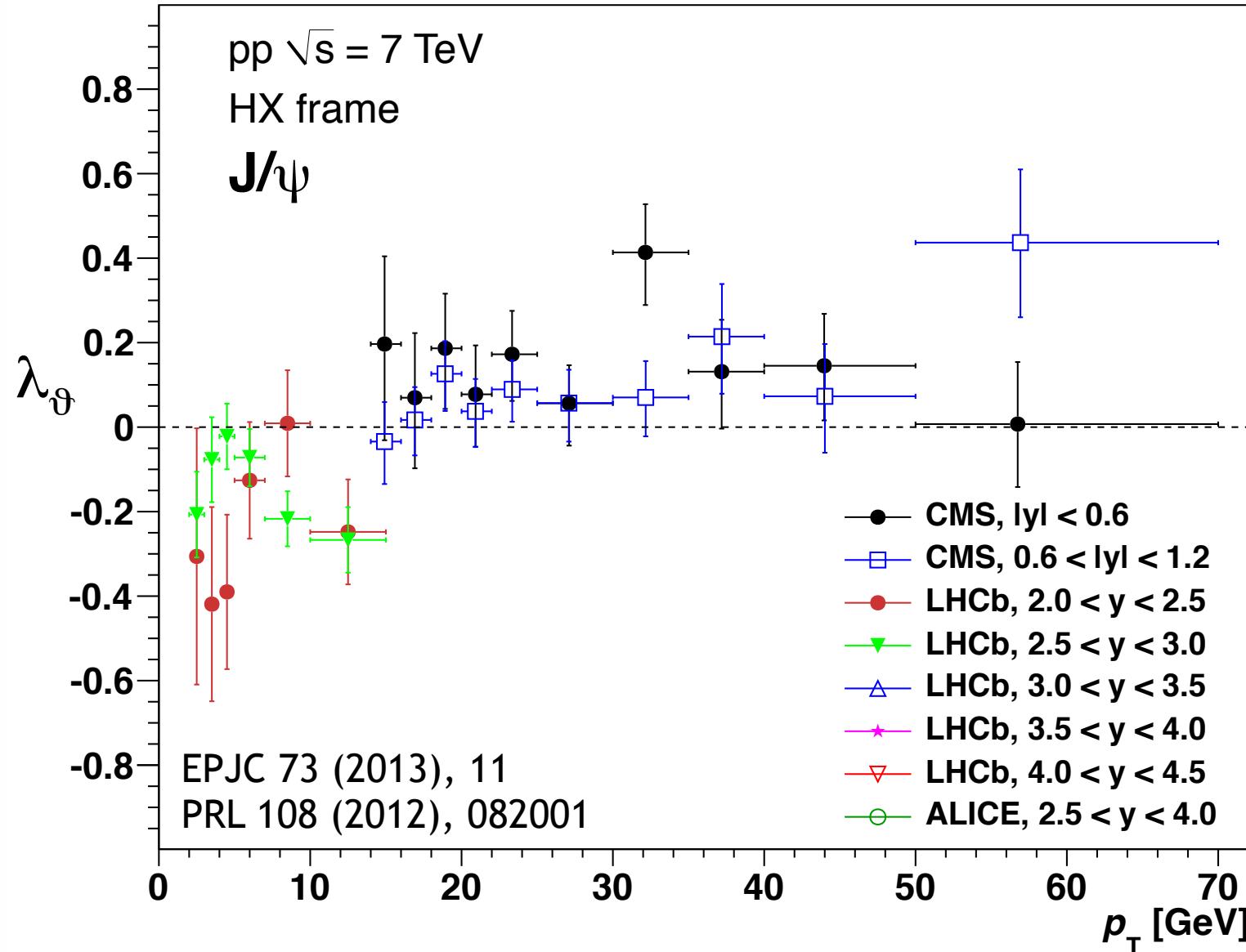
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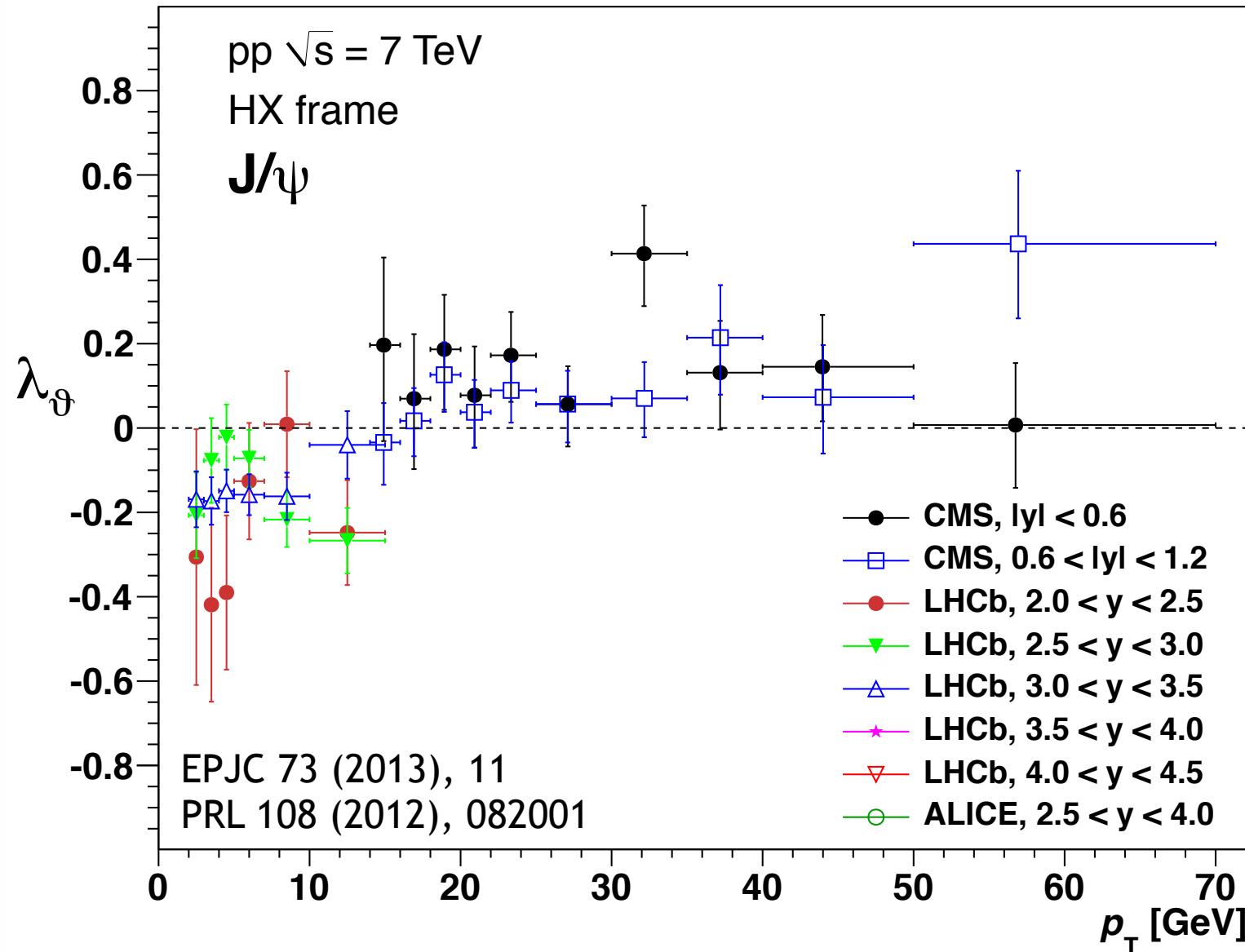
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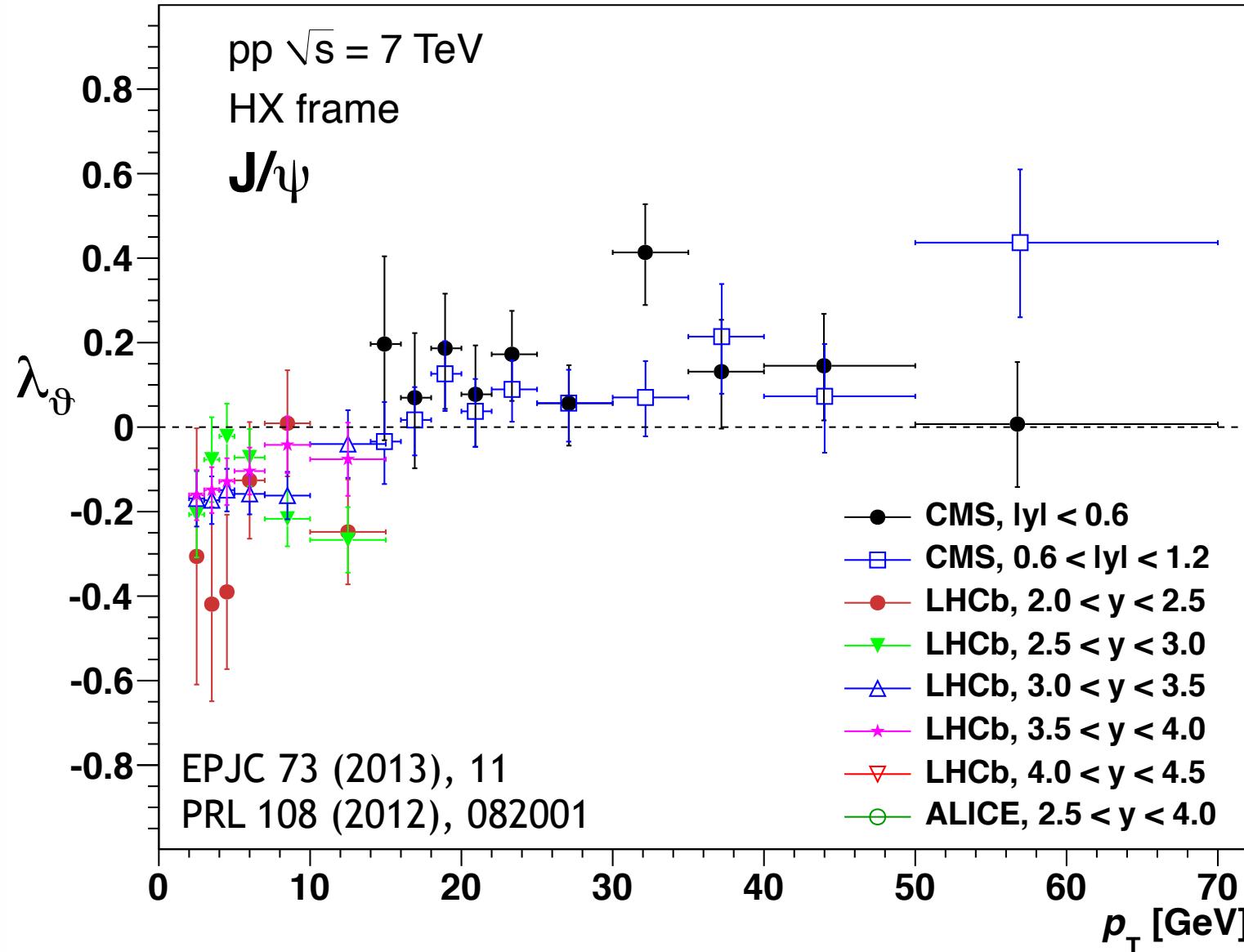
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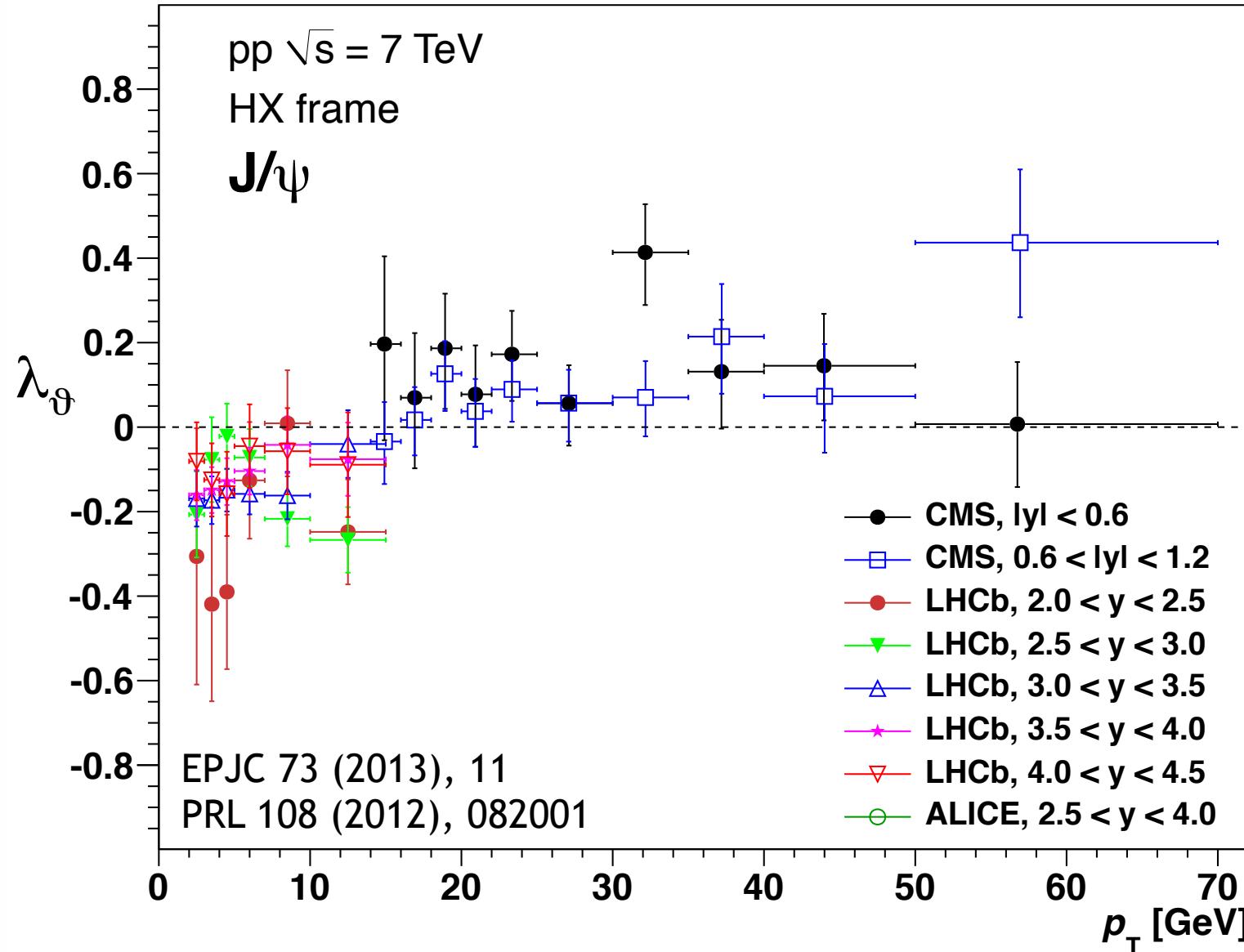
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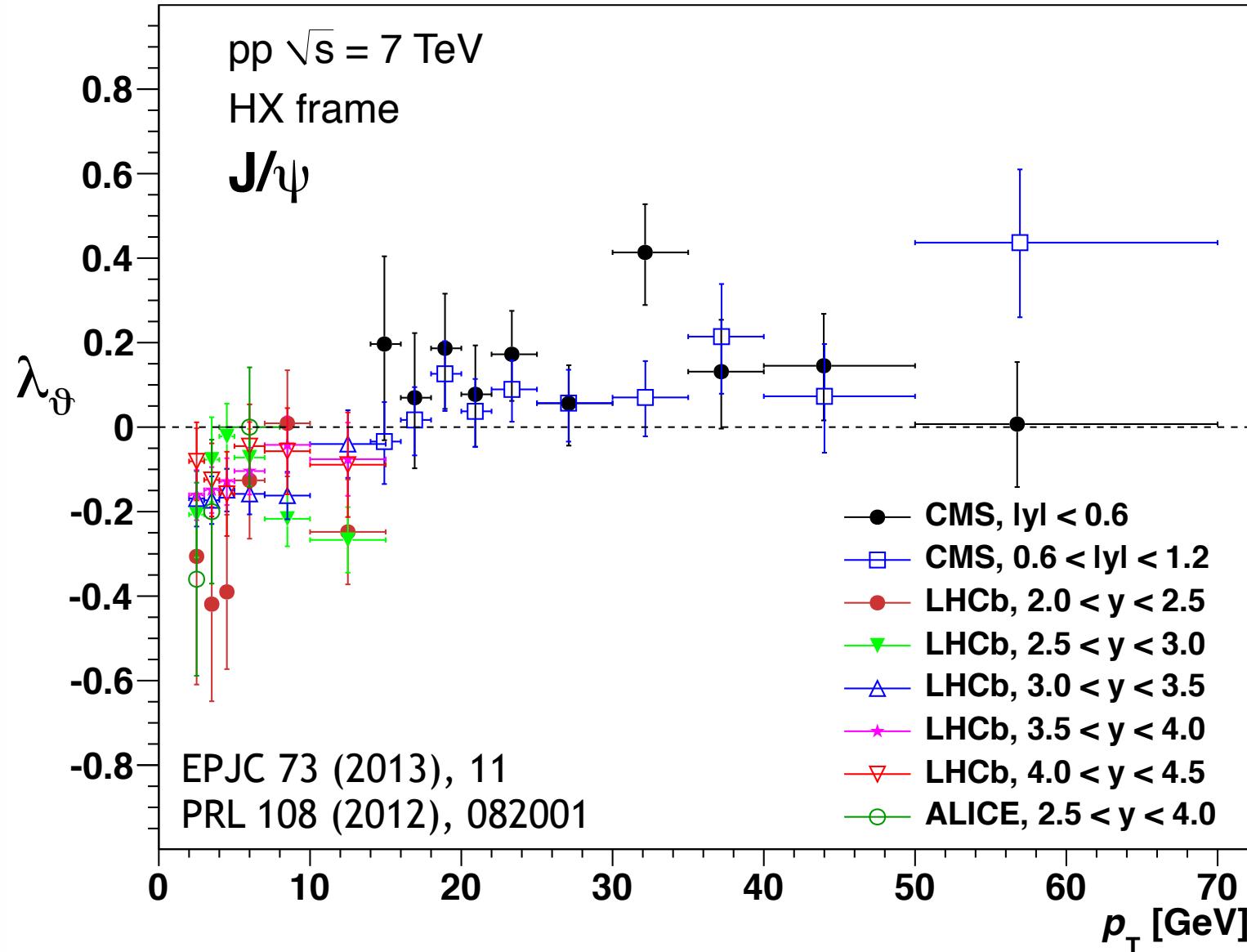
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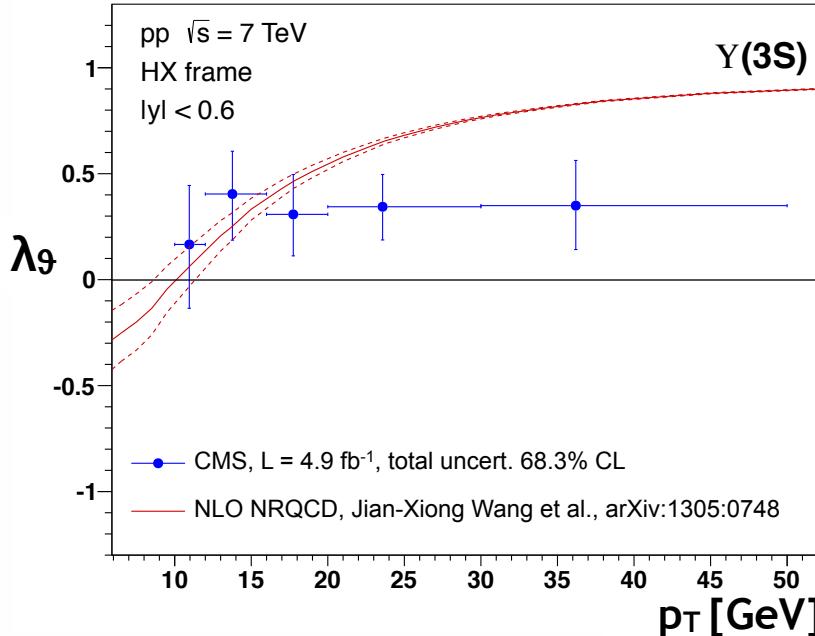
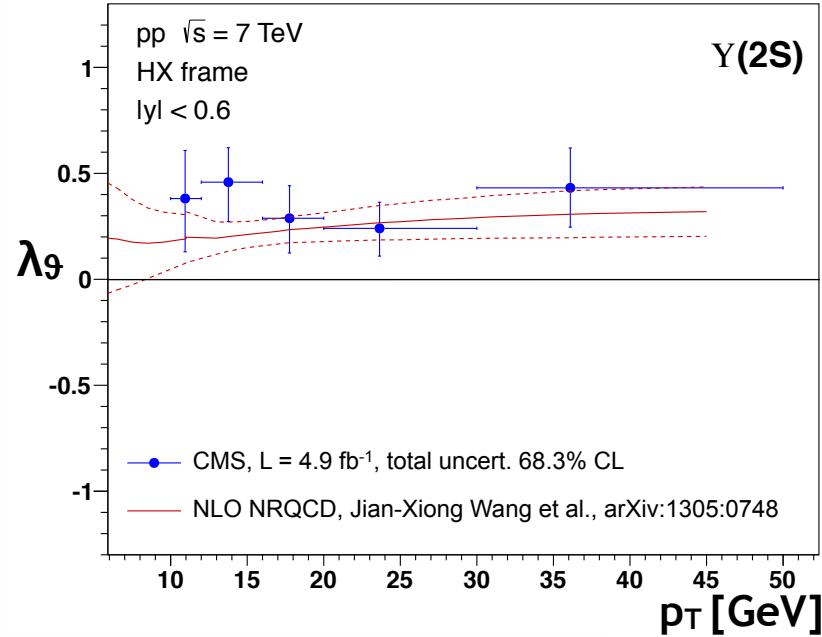
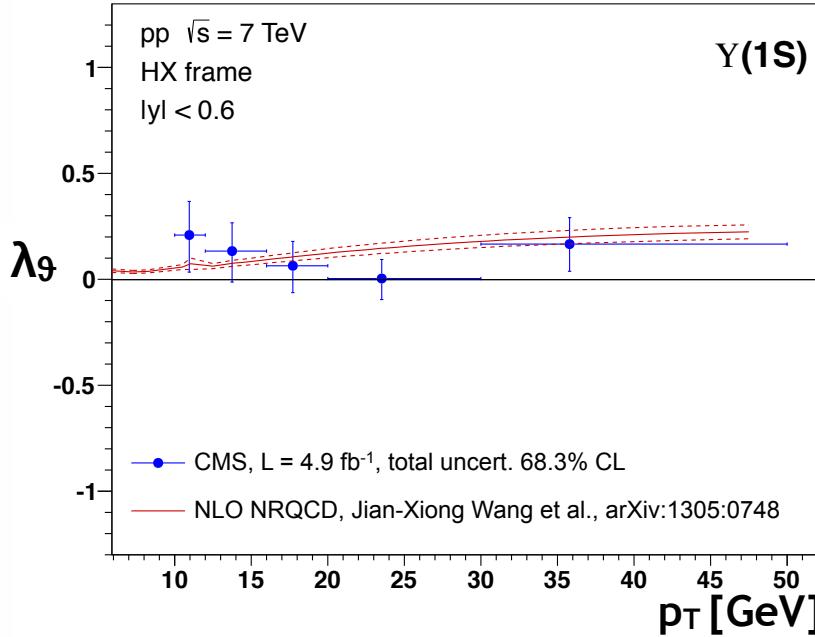


Comparison to Other LHC Experiments: J/ ψ

All LHC results are compatible with each other.



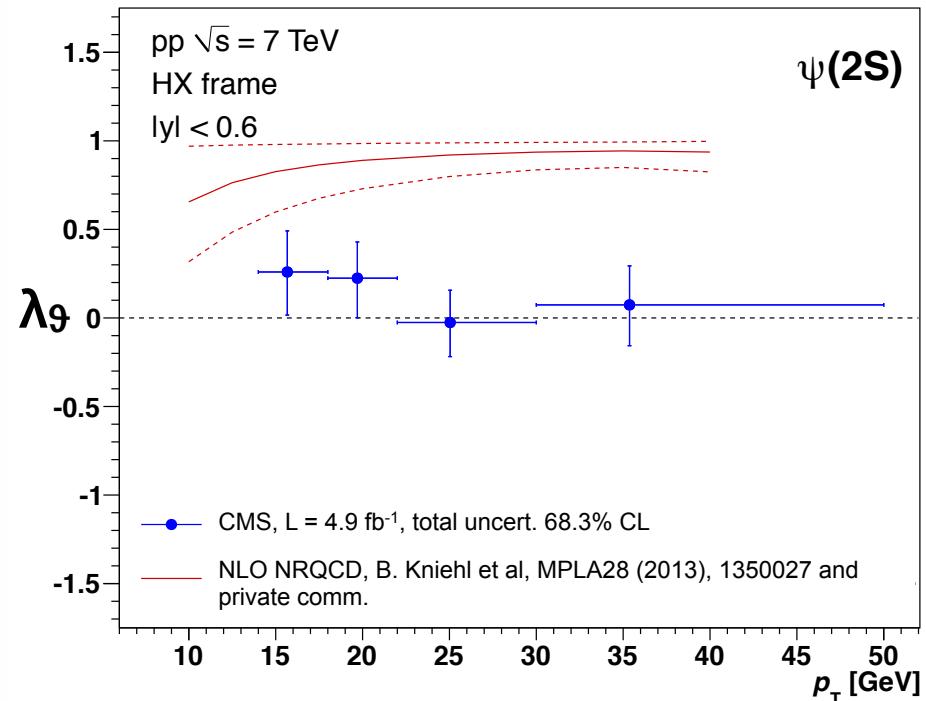
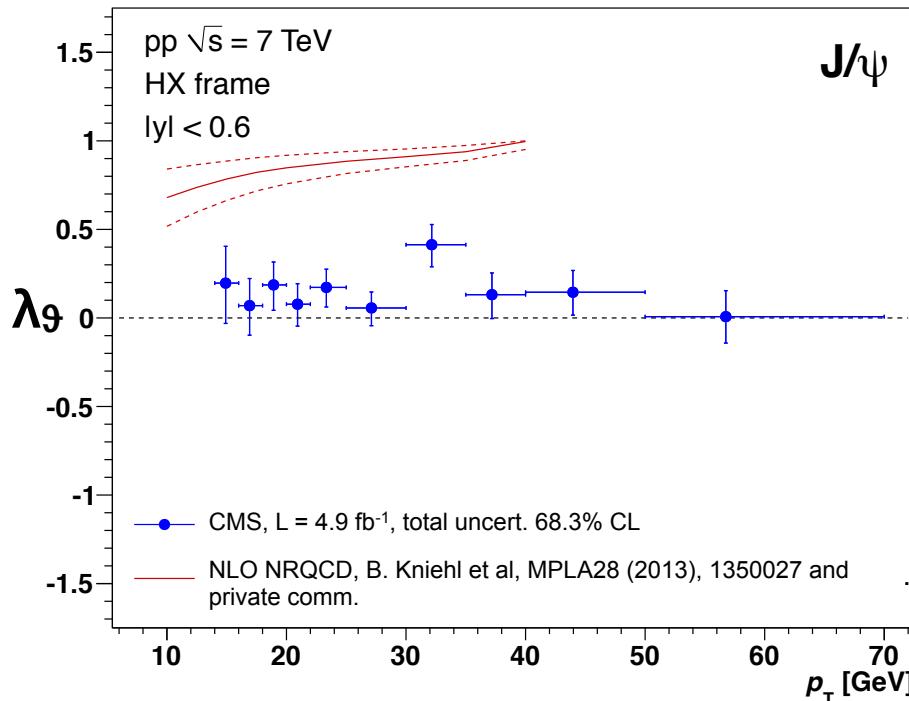
NLO NRQCD Comparison: $\Upsilon(nS)$



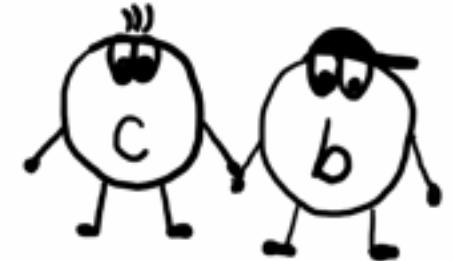
- Theory calculation accounts for feed-down contributions to $\Upsilon(1S)$ and $\Upsilon(2S)$ states
- Prediction for $\Upsilon(3S)$ may change when including feed-down from $\chi_b(3P)$ states
- Color octet matrix elements are fit to hadroproduction data only, including CMS $\Upsilon(nS)$ polarization results

NLO NRQCD Comparison: $\psi(nS)$

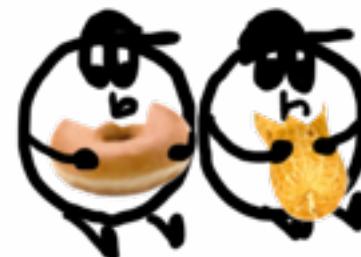
- Calculations use a global fit of color octet matrix elements to photo- as well as hadroproduction data, excluding polarization results
- NLO NRQCD calculations fail to describe CMS results
- Theory predictions only consider directly produced J/ ψ 's
- $\psi(2S)$ result can be directly compared to theory since it has no feed-down contribution



Summary

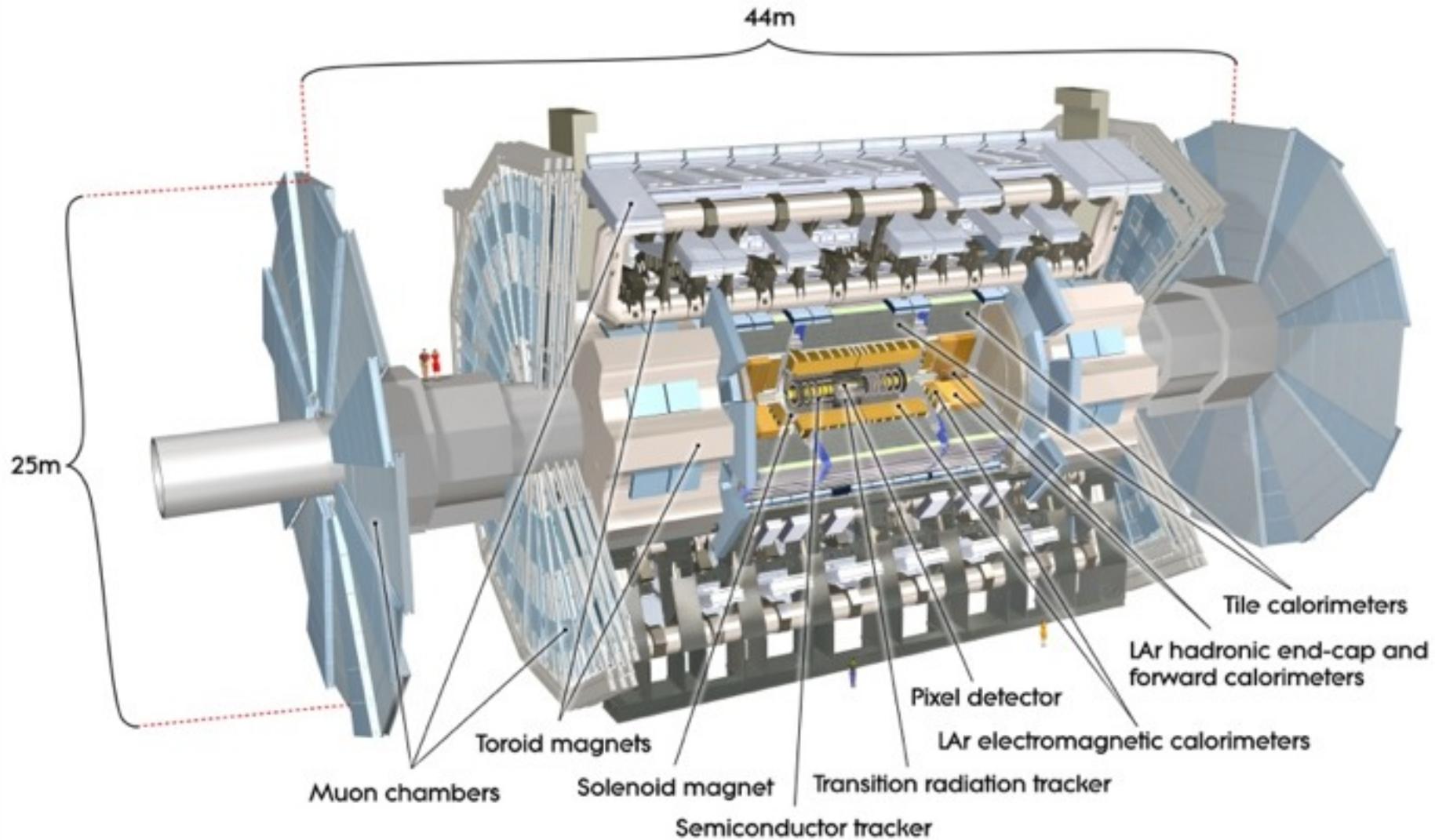


- ATLAS and CMS have a rich program in heavy flavor physics
- Topics shown today:
 - Measurement of α_b in $\Lambda_b^0 \rightarrow J/\psi \Lambda_0$ inconsistent with pQCD and HQET at a level of 2.6 and 2.8 standard deviations
 - $B_c^\pm \rightarrow J/\psi \pi^\pm$ and $B_c^\pm \rightarrow J/\psi \pi^\pm \pi^\pm \pi^\mp$ modes observed
 - Quarkonium polarization measurements in disagreement with NLO NRQCD calculations
- More results are expected to come in the near future

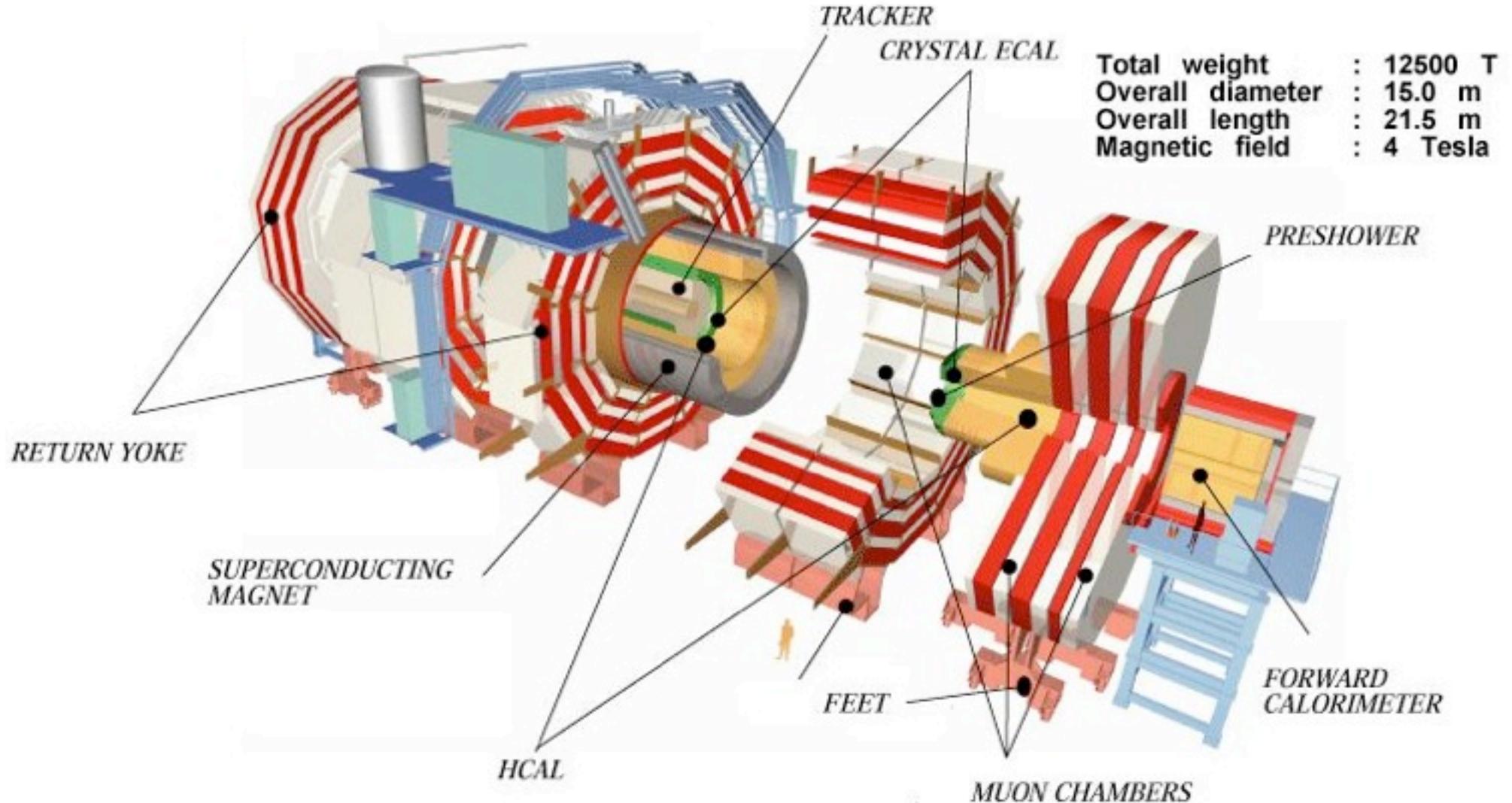


BACKUP

ATLAS Detector



CMS Detector



$$\alpha_b \text{ in } \Lambda_b^0 \rightarrow J/\psi \Lambda_0$$

Full angular probability density function of the decay angles

$$w(\Omega, \vec{A}, P) = \frac{1}{(4\pi)^3} \sum_{i=0}^{19} f_{1i}(\vec{A}) f_{2i}(P, \alpha_\Lambda) F_i(\Omega)$$

Ω : decay angles

\vec{A} : helicity amplitudes

α_Λ : decay asymmetry parameter

P: polarization of Λ_b^0



α_b in $\Lambda_b^0 \rightarrow J/\psi \Lambda_0$

| i | f_{1i} | f_{2i} | F_i |
|-----|--|-------------------|---|
| 0 | $a_+a_+^* + a_-a_-^* + b_+b_+^* + b_-b_-^*$ | 1 | 1 |
| 1 | $a_+a_+^* - a_-a_-^* + b_+b_+^* - b_-b_-^*$ | P | $\cos\theta$ |
| 2 | $a_+a_+^* - a_-a_-^* - b_+b_+^* + b_-b_-^*$ | α_Λ | $\cos\theta_1$ |
| 3 | $a_+a_+^* + a_-a_-^* - b_+b_+^* - b_-b_-^*$ | $P\alpha_\Lambda$ | $\cos\theta \cos\theta_1$ |
| 4 | $-a_+a_+^* - a_-a_-^* + \frac{1}{2}b_+b_+^* + \frac{1}{2}b_-b_-^*$ | 1 | $\frac{1}{2}(3\cos^2\theta_2 - 1)$ |
| 5 | $-a_+a_+^* + a_-a_-^* + \frac{1}{2}b_+b_+^* - \frac{1}{2}b_-b_-^*$ | P | $\frac{1}{2}(3\cos^2\theta_2 - 1) \cos\theta$ |
| 6 | $-a_+a_+^* + a_-a_-^* - \frac{1}{2}b_+b_+^* + \frac{1}{2}b_-b_-^*$ | α_Λ | $\frac{1}{2}(3\cos^2\theta_2 - 1) \cos\theta_1$ |
| 7 | $-a_+a_+^* - a_-a_-^* - \frac{1}{2}b_+b_+^* - \frac{1}{2}b_-b_-^*$ | $P\alpha_\Lambda$ | $\frac{1}{2}(3\cos^2\theta_2 - 1) \cos\theta \cos\theta_1$ |
| 8 | $-3Re(a_+a_-^*)$ | $P\alpha_\Lambda$ | $\sin\theta \sin\theta_1 \sin^2\theta_2 \cos\varphi_1$ |
| 9 | $3Im(a_+a_-^*)$ | $P\alpha_\Lambda$ | $\sin\theta \sin\theta_1 \sin^2\theta_2 \sin\varphi_1$ |
| 10 | $-\frac{3}{2}Re(b_-b_+^*)$ | $P\alpha_\Lambda$ | $\sin\theta \sin\theta_1 \sin^2\theta_2 \cos(\varphi_1 + 2\varphi_2)$ |
| 11 | $\frac{3}{2}Im(b_-b_+^*)$ | $P\alpha_\Lambda$ | $\sin\theta \sin\theta_1 \sin^2\theta_2 \sin(\varphi_1 + 2\varphi_2)$ |
| 12 | $-\frac{3}{\sqrt{2}}Re(b_-a_+^* + a_-b_+^*)$ | $P\alpha_\Lambda$ | $\sin\theta \cos\theta_1 \sin\theta_2 \cos\theta_2 \cos\varphi_2$ |
| 13 | $\frac{3}{\sqrt{2}}Im(b_-a_+^* + a_-b_+^*)$ | $P\alpha_\Lambda$ | $\sin\theta \cos\theta_1 \sin\theta_2 \cos\theta_2 \sin\varphi_2$ |
| 14 | $-\frac{3}{\sqrt{2}}Re(b_-a_-^* + a_+b_+^*)$ | $P\alpha_\Lambda$ | $\cos\theta \sin\theta_1 \sin\theta_2 \cos\theta_2 \cos(\varphi_1 + \varphi_2)$ |
| 15 | $\frac{3}{\sqrt{2}}Im(b_-a_-^* + a_+b_+^*)$ | $P\alpha_\Lambda$ | $\cos\theta \sin\theta_1 \sin\theta_2 \cos\theta_2 \sin(\varphi_1 + \varphi_2)$ |
| 16 | $\frac{3}{\sqrt{2}}Re(a_-b_+^* - b_-a_+^*)$ | P | $\sin\theta \sin\theta_2 \cos\theta_2 \cos\varphi_2$ |
| 17 | $-\frac{3}{\sqrt{2}}Im(a_-b_+^* - b_-a_+^*)$ | P | $\sin\theta \sin\theta_2 \cos\theta_2 \sin\varphi_2$ |
| 18 | $\frac{3}{\sqrt{2}}Re(b_-a_-^* - a_+b_+^*)$ | α_Λ | $\sin\theta_1 \sin\theta_2 \cos\theta_2 \cos(\varphi_1 + \varphi_2)$ |
| 19 | $-\frac{3}{\sqrt{2}}Im(b_-a_-^* - a_+b_+^*)$ | α_Λ | $\sin\theta_1 \sin\theta_2 \cos\theta_2 \sin(\varphi_1 + \varphi_2)$ |



Fit model parametrization

$$\alpha_b = |a_+|^2 - |a_-|^2 + |b_+|^2 - |b_-|^2$$

$$k_+ = \frac{|a_+|}{\sqrt{|a_+|^2 + |b_+|^2}}$$

$$k_- = \frac{|b_-|}{\sqrt{|a_-|^2 + |b_-|^2}}$$

$$\Delta_+ = \rho_+ - w_+$$

$$\Delta_- = \rho_- - w_-$$

a,b: helicity amplitudes

ρ , w : phases of helicity amplitudes



Quarkonium polarization

Definition of the PPD

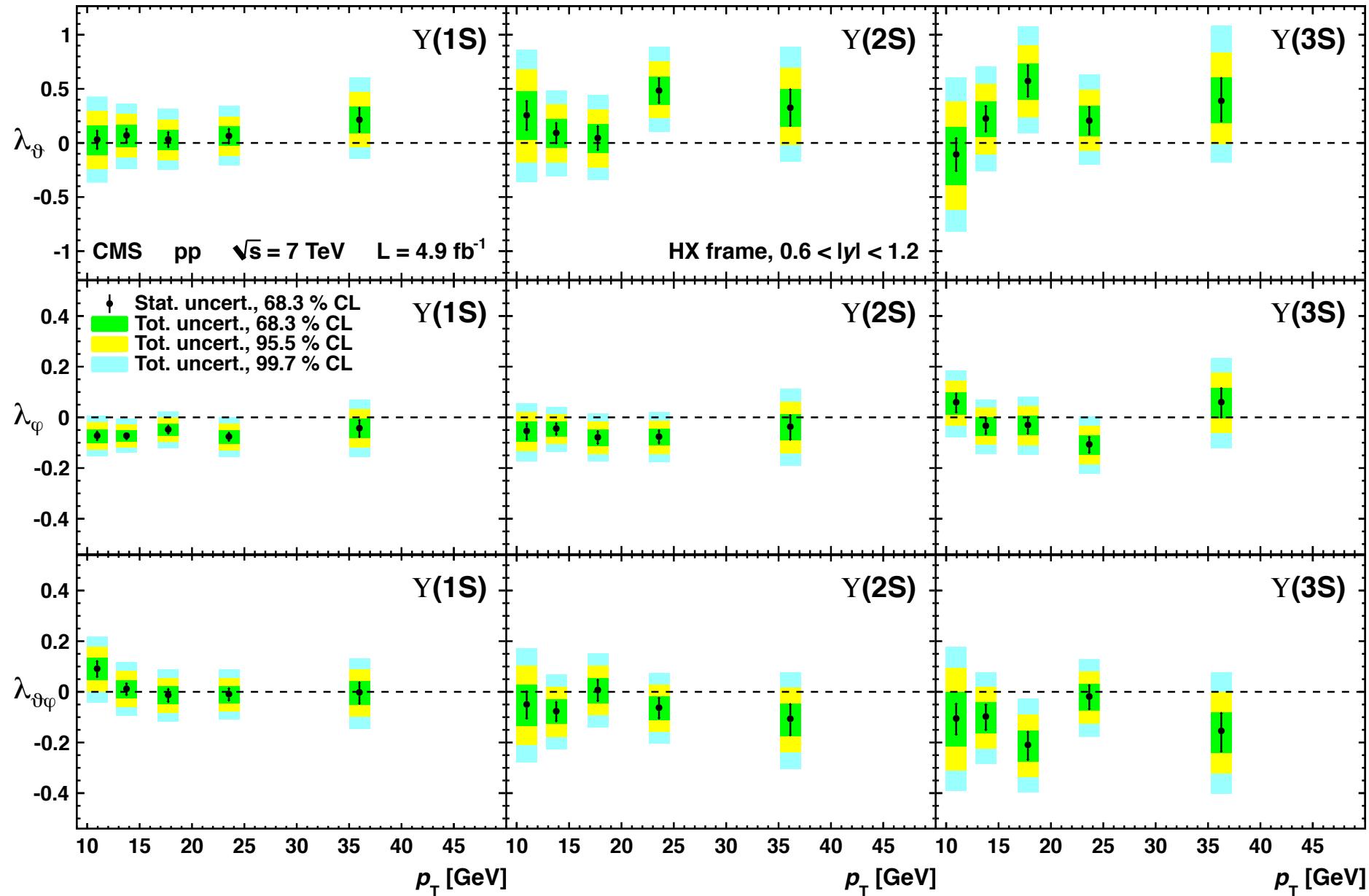
$$\mathcal{P}(\vec{\lambda}) \propto \prod_i \frac{1}{\mathcal{N}(\vec{\lambda})} W(\cos \theta^{(i)}, \phi^{(i)} | \vec{\lambda}) \varepsilon(p_1^{(i)} p_2^{(i)})$$

\mathcal{N} : normalization

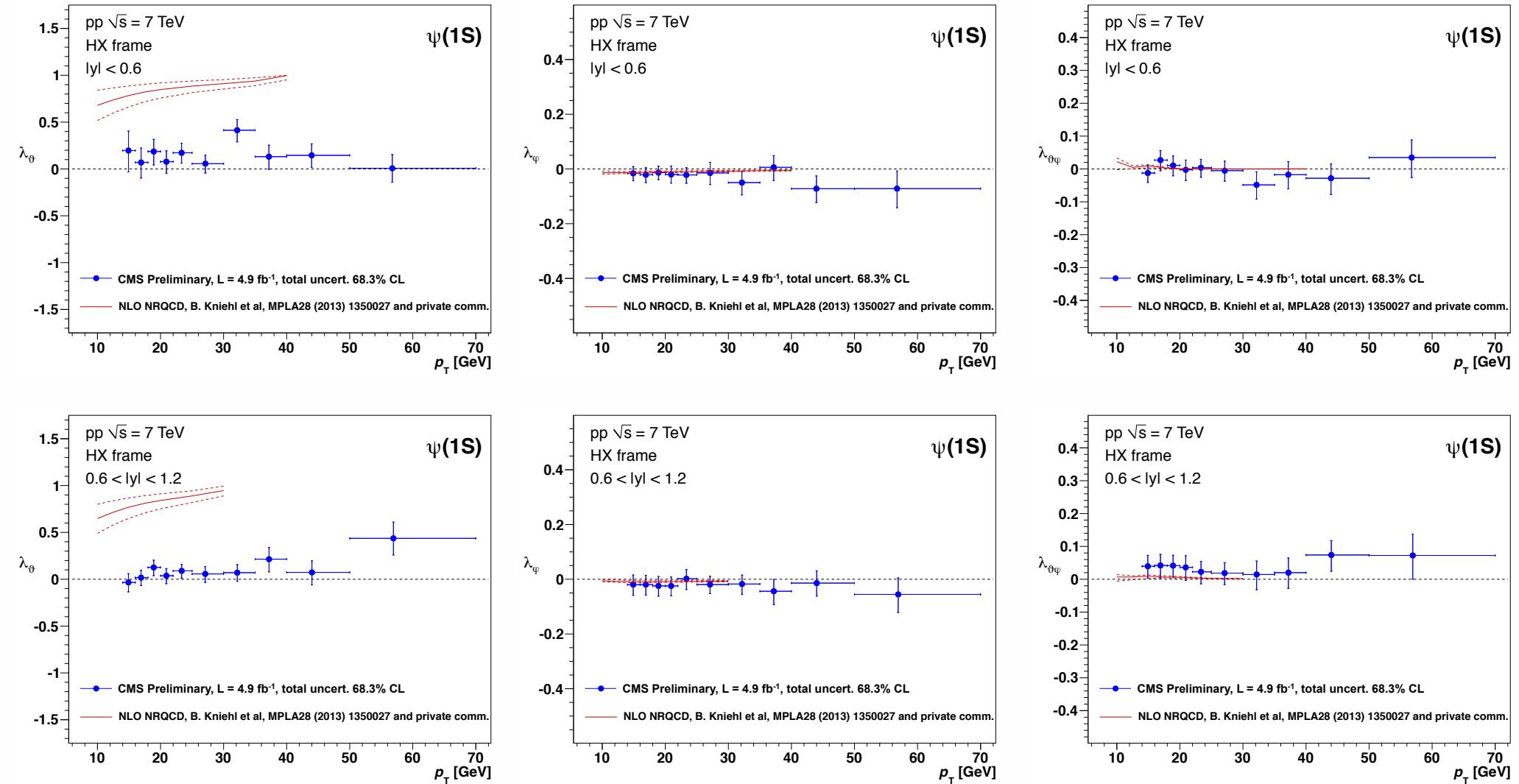
W : general angular distribution

ε : dimuon efficiency as a function of the muon momenta

$\Upsilon(nS)$ Polarization in the HX Frame, $0.6 < |y| < 1.2$



$\psi(1S)$: Comparison to NLO NRQCD



$\psi(2S)$: Comparison to NLO NRQCD

