

Neutrino electromagnetic properties and new effects in astrophysics

Alexander Studenikin

“Particle Physics
and Cosmology”
XXVI Rencontres de Blois
21/05/2014

Moscow State
University

&

JINR -Dubna



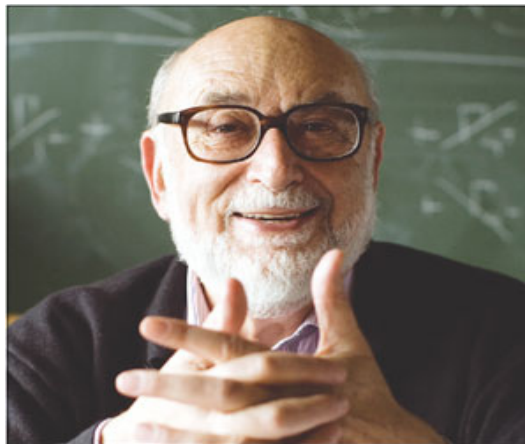
- Blois - 2014

Particle Physics and Cosmology

... a period of spectacular success of
Particle Physics ...



Robert Brout



François Englert



Peter Higgs

Observation of **Higgs boson** confirms the
symmetry breaking mechanism by

Brout-Englert-Higgs (BEH)

provides final glorious triumph of

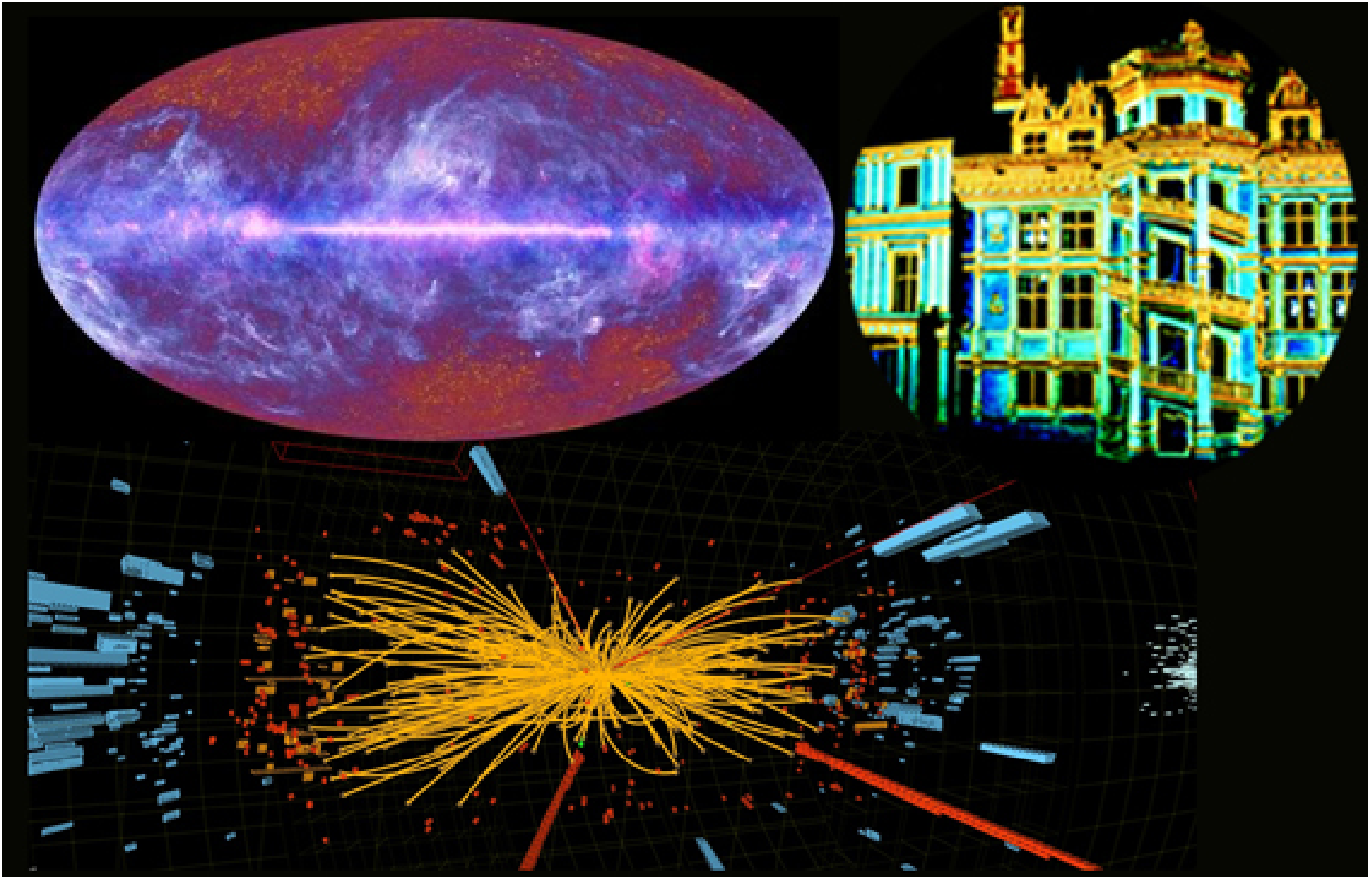
Standard Model

What

is

next ?

26th Rencontres de Blois Particle Physics and Cosmology



Château Royal de Blois

May 18-23, 2014

Prof. Jean Trân Thanh Vân
(opening speech at the conference) :

... the primary goal
of *XXVI Rencontres de Blois*
is to discuss phenomena of vast
variety scales -
from extremely *fine scales* where
fundamentals of particle physics
should manifest themselves
to *huge scales*
related to phenomenology
in astrophysics and cosmology...

- ν electromagnetic properties

- 1 C.Giunti, A.Studenikin :
“Electromagnetic interactions of neutrinos:
a window to new physics” ,
arXiv: 1403.6344 , March 25, 2014

- Astrophysical consequences of ν interactions

- 2 A. Studenikin , I.Tokarev:
“Millicharged neutrino with anomalous
magnetic moment in rotating magnetized matter” ,
Nucl.Phys.B (2014)
<http://dx.doi.org/10.1016/j.nuclphysb.2014.04.026>
- 3 I.Balantsev, A.Studenikin :
“Spin light of electron in dense neutrino fluxes” ,
arXiv: 1405.xxxx , May 19, 2014



electromagnetic
properties

(short review)

$$m_\nu \neq 0$$



① Carlo Giunti, Alexander Studenikin :
“Neutrino electromagnetic properties”
Phys.Atom.Nucl. 73, 2089 (2009)

② A. Studenikin : “Neutrino magnetic moment:
a window to new physics”
Nucl.Phys.B (Proc.Supl.) 188, 220 (2009)

③ C. Giunti, A. Studenikin :
“Electromagnetic properties of neutrino”
J.Phys.: Conf.Series. 203 (2010) 012100

④ Carlo Broggin, C. Giunti, A. Studenikin :
“Electromagnetic properties of neutrinos” ,
in: Special issue “Neutrino Physics”,
Adv. in High Energy Phys. 2012 (2012) 459526 (49 pp)

⑤ C. Giunti, A. Studenikin :
“Electromagnetic interactions of neutrinos:
a window to new physics”,
arXiv:1403.6344 , March 25, 2014, sent to Rev.Mod.Phys.

$$m_\nu \neq 0$$

... a tool for studying physics
Beyond Standard Model...

Theory (Standard Model with ν_R)

$$\mu_\nu = \frac{3eG_F}{8\sqrt{2}\pi^2} m_\nu \sim 3 \cdot 10^{-19} \mu_B \left(\frac{m_\nu}{1\text{eV}} \right), \quad \mu_B = \frac{e}{2m_e}$$

magnetic moment

$$a_e = \frac{\alpha_{QED}}{2\pi} \sim 10^{-3}$$



Lee Shrock, 1977; Fujikawa Shrock, 1980

... much greater values are desired
for astrophysical or cosmology

visualization of μ_ν

... hopes for physics BSM ...

• Astrophysical bounds

$$\mu_\nu \leq 3 \cdot 10^{-12} \mu_B$$

G. Raffelt, D. Dearborn,
J. Silk, 1989.

• Theory (Standard Model with ν_R)

$$\mu_\nu = \frac{3eG_F}{8\sqrt{2}\pi^2} m_\nu \sim 3 \cdot 10^{-19} \mu_B \left(\frac{m_\nu}{\text{eV}} \right), \quad \mu_B = \frac{e}{2m_e}$$

Lee Shrock, 1977; Fujikawa Shrock, 1980

• Limits from reactor ν - e scattering experiments

A. Beda et al. (GEMMA Coll.)
(2012):

$$\mu_\nu < 2.9 \times 10^{-11} \mu_B$$


...the present status...

to have visible $\mu \neq 0$

is not an easy task for

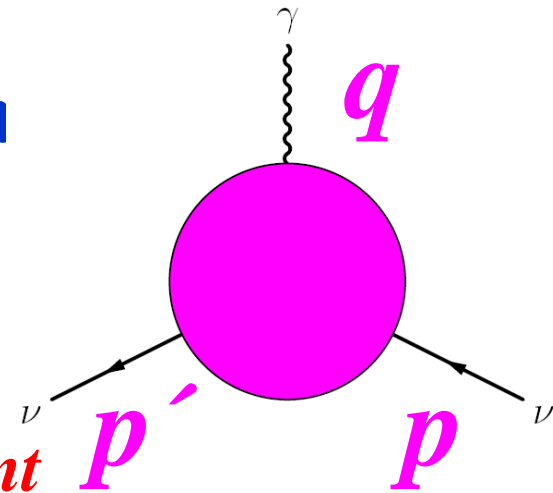
theoreticians

and experimentalists

... a bit of  electromagnetic
properties theory ...

✓ electromagnetic vertex function

$$\langle \psi(p') | J_\mu^{EM} | \psi(p) \rangle = \bar{u}(p') \Lambda_\mu(q, l) u(p)$$



Matrix element of **electromagnetic current** is a Lorentz vector

$\Lambda_\mu(q, l)$ should be constructed using

matrices $\hat{1}, \gamma_5, \gamma_\mu, \gamma_5 \gamma_\mu, \sigma_{\mu\nu},$

tensors $g_{\mu\nu}, \epsilon_{\mu\nu\sigma\gamma}$

vectors q_μ and l_μ

$$q_\mu = p'_\mu - p_\mu, \quad l_\mu = p'_\mu + p_\mu$$

Lorentz covariance (1)
and electromagnetic gauge invariance (2)



Matrix element of electromagnetic current between neutrino states

$$\langle \nu(p') | J_\mu^{EM} | \nu(p) \rangle = \bar{u}(p') \Lambda_\mu(q) u(p)$$

where vertex function generally contains 4 form factors

$$\Lambda_\mu(q) = f_Q(q^2) \gamma_\mu + f_M(q^2) i \sigma_{\mu\nu} q^\nu - f_E(q^2) \sigma_{\mu\nu} q^\nu \gamma_5 + f_A(q^2) (q^2 \gamma_\mu - q_\mu \not{q}) \gamma_5$$

1. electric dipole 2. magnetic 3. electric 4. anapole

● Hermiticity and discrete symmetries of EM current J_μ^{EM} put constraints on form factors

Dirac ✓

- 1) CP invariance + hermiticity $\implies f_E = 0$,
- 2) at zero momentum transfer **only** electric charge $f_Q(0)$ and magnetic moment $f_M(0)$ contribute $H_{int} \sim J_\mu^{EM} A^\mu$
- 3) hermiticity **itself** \implies three form factors are real: $Im f_Q = Im f_M = Im f_A = 0$

Majoran ✓

- 1) from CPT invariance (regardless CP or ~~CP~~).

$$f_Q = f_M = f_E = 0$$



...as early as 1939, W.Pauli...

EM properties \implies a way to distinguish Dirac and Majorana ✓

✓ magnetic moment
in experiments

Studies of ν - e scattering

- most sensitive method for experimental investigation of μ_ν

Cross-section:

$$\bullet \quad \frac{d\sigma}{dT}(\nu + e \rightarrow \nu + e) = \left(\frac{d\sigma}{dT}\right)_{\text{SM}} + \left(\frac{d\sigma}{dT}\right)_{\mu_\nu}$$

where the Standard Model contribution

$$\bullet \quad \left(\frac{d\sigma}{dT}\right)_{\text{SM}} = \frac{G_F^2 m_e}{2\pi} \left[(g_V + g_A)^2 + (g_V - g_A)^2 \left(1 - \frac{T}{E_\nu}\right)^2 + (g_A^2 - g_V^2) \frac{m_e T}{E_\nu^2} \right],$$

T is the electron recoil energy and

$$\bullet \quad \left(\frac{d\sigma}{dT}\right)_{\mu_\nu} = \frac{\pi \alpha_{em}^2}{m_e^2} \left[\frac{1 - T/E_\nu}{T} \right] \mu_\nu^2$$

$$\mu_\nu^2 = \sum_{j = \nu_e, \nu_\mu, \nu_\tau} | \mu_{ij} - \epsilon_{ij} |^2$$

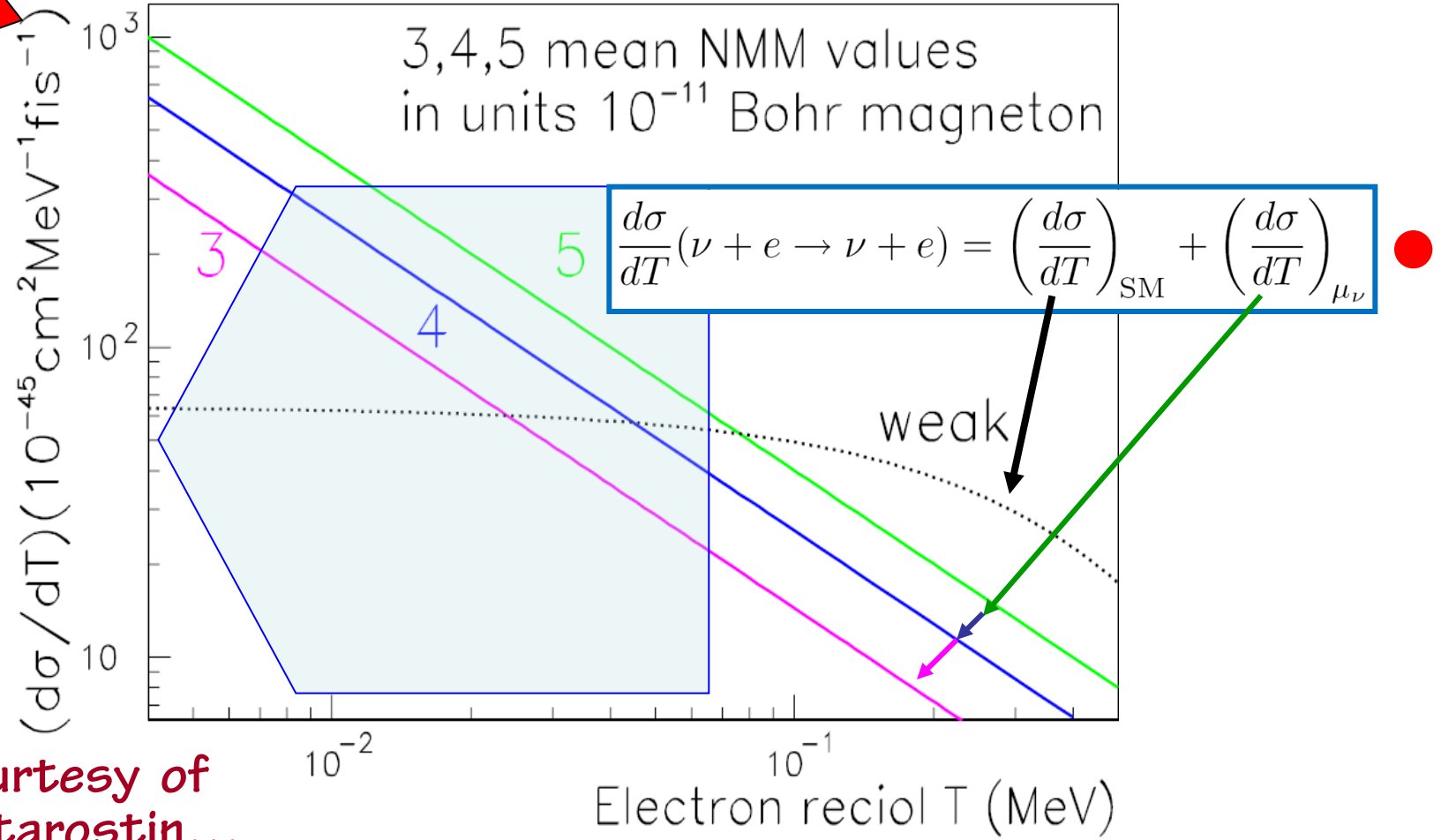
$$g_V = \begin{cases} 2 \sin^2 \theta_W + \frac{1}{2} & \text{for } \nu_e, \\ 2 \sin^2 \theta_W - \frac{1}{2} & \text{for } \nu_\mu, \nu_\tau, \end{cases} \quad g_A = \begin{cases} \frac{1}{2} & \text{for } \nu_e, \\ -\frac{1}{2} & \text{for } \nu_\mu, \nu_\tau \end{cases} \quad \begin{matrix} \text{for anti-neutrinos} \\ g_A \rightarrow -g_A \end{matrix}$$

to incorporate charge radius: $g_V \rightarrow g_V + \frac{2}{3} M_W^2 \langle r^2 \rangle \sin^2 \theta_W$

Magnetic moment contribution dominates at low electron recoil energies

recoil energies when $\left(\frac{d\sigma}{dT}\right)_{\mu\nu} > \left(\frac{d\sigma}{dT}\right)_{SM}$ and $\frac{T}{m_e} < \frac{\pi^2 \alpha_{em}}{G_F^2 m_e^4} \mu_\nu^2$

... the lower the smallest measurable electron recoil energy is, smaller values of μ_ν^2 can be probed in scattering experiments ...



... courtesy of A.Starostin...



MUNU experiment at Bugey reactor (2005)

$$\mu_{\nu} \leq 9 \times 10^{-11} \mu_B$$

TEXONO collaboration at Kuo-Sheng power plant (2006)

$$\mu_{\nu} \leq 7 \times 10^{-11} \mu_B$$

GEMMA (2007)

$$\mu_{\nu} \leq 5.8 \times 10^{-11} \mu_B$$

GEMMA I 2005 - 2007

BOREXINO (2008)

$$\mu_{\nu} \leq 5.4 \times 10^{-11} \mu_B$$

...was considered as the world best constraint..

$$\mu_{\nu} \leq 8.5 \times 10^{-11} \mu_B \quad (\nu_{\tau}, \nu_{\mu})$$

based on first release of
BOREXINO data

Montanino,
Picariello,
Pulido,
PRD 2008

... attempts to
improve bounds



GEMMA (2005-2012)

JINR (Dubna) + ITEP (Moscow) at Kalinin Nuclear Power Plant

World best experimental limit

$$\mu_\nu < 2.9 \times 10^{-11} \mu_B$$

June 2012

A. Beda et al, in: **Special Issue on “Neutrino Physics”**,
Advances in High Energy Physics (2012) 2012,
editors: J. Bernabeu, G. Fogli, A. McDonald, K. Nishikawa

... quite **realistic prospects** of the near future

$$\mu_\nu \sim 1 \times 10^{-11} \mu_B$$

(V. Brudanin, A. Starostin, priv. comm.)

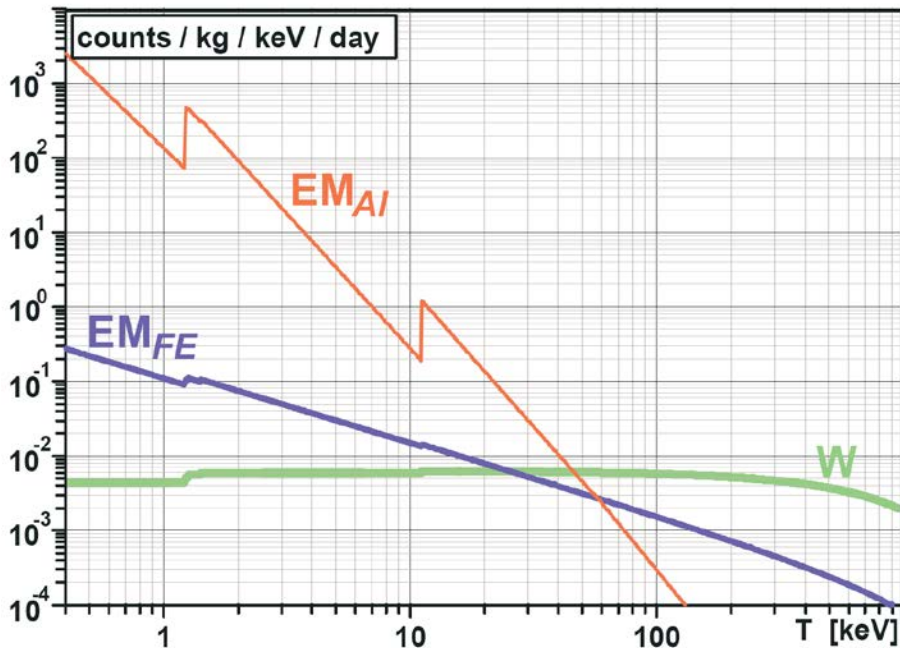
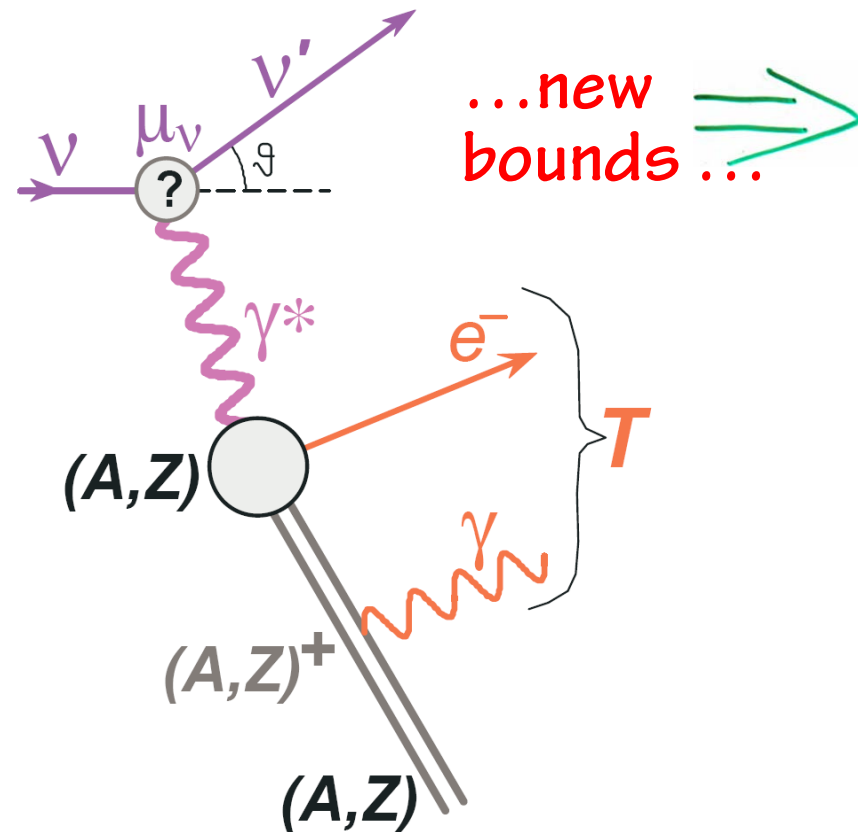
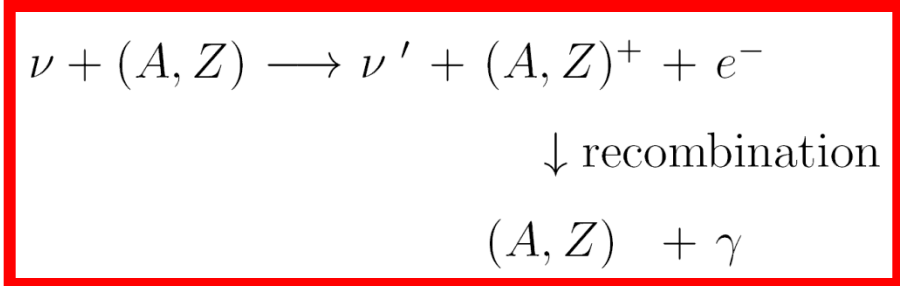
... quite recent **claim**
 that ν - e cross section
 should be increased by
Atomic Ionization Effect:



H.Wong et al. (TEXONO Coll.),
 PRL 105 (2010)
 061801

(ν scattering on bound e)

... an interesting hypothetical
 possibility to improve bounds...



K.Kouzakov, A.Studenikin,

- “Magnetic neutrino scattering on atomic electrons revisited”
Phys.Lett. B 105 (2011) 061801,
- “Electromagnetic neutrino-atom collisions: The role of electron binding”
Nucl.Phys.B (Proc.Suppl.) 217 (2011) 353

K.Kouzakov, A.Studenikin, M.Voloshin,

- “Neutrino electromagnetic properties and new bounds on neutrino magnetic moments” J.Phys.: Conf.Ser. 375 (2012) 042045
 - “Neutrino-impact ionization of atoms in search for neutrino magnetic moment”, Phys.Rev.D 83 (2011) 113001
 - “On neutrino-atom scattering in searches for neutrino magnetic moments” Nucl.Phys.B (Proc.Supp.) 2011 (Proc. of Neutrino 2010 Conf.)
 - “Testing neutrino magnetic moment in ionization of atoms by neutrino impact”, JETP Lett. 93 (2011) 699
- M.Voloshin,
- “Neutrino scattering on atomic electrons in search for neutrino magnetic moment”
Phys.Rev.Lett. 105 (2010) 201801

No important effect of
Atomic Ionization on cross section in
 μ experiments once all possible final
electronic states accounted for

...free electron approximation ...

● M.Voloshin, 23 Aug 2010;

K.Kouzakov, A.Studenikin, 26 Nov 2010;

H.Wong et al, arXiv: 1001.2074 V3, 28 Nov 2010

Experimental limits for different effective neutrino magnetic moments

Method	Experiment	Limit	CL	Reference
Reactor $\bar{\nu}_e-e^-$	Krasnoyarsk	$\mu_{\nu_e} < 2.4 \times 10^{-10} \mu_B$	90%	Vidyakin <i>et al.</i> (1992)
	Rovno	$\mu_{\nu_e} < 1.9 \times 10^{-10} \mu_B$	95%	Derbin <i>et al.</i> (1993)
	● MUNU	$\mu_{\nu_e} < 0.9 \times 10^{-10} \mu_B$	90%	Daraktchieva <i>et al.</i> (2005)
	● TEXONO	$\mu_{\nu_e} < 7.4 \times 10^{-11} \mu_B$	90%	Wong <i>et al.</i> (2007)
	● GEMMA	$\mu_{\nu_e} < 2.9 \times 10^{-11} \mu_B$	90%	Beda <i>et al.</i> (2012)
Accelerator ν_e-e^-	LAMPF	$\mu_{\nu_e} < 10.8 \times 10^{-10} \mu_B$	90%	Allen <i>et al.</i> (1993)
Accelerator $(\nu_\mu, \bar{\nu}_\mu)-e^-$	BNL-E734	$\mu_{\nu_\mu} < 8.5 \times 10^{-10} \mu_B$	90%	Ahrens <i>et al.</i> (1990)
	LAMPF	$\mu_{\nu_\mu} < 7.4 \times 10^{-10} \mu_B$	90%	Allen <i>et al.</i> (1993)
	LSND	$\mu_{\nu_\mu} < 6.8 \times 10^{-10} \mu_B$	90%	Auerbach <i>et al.</i> (2001)
Accelerator $(\nu_\tau, \bar{\nu}_\tau)-e^-$	DONUT	$\mu_{\nu_\tau} < 3.9 \times 10^{-7} \mu_B$	90%	Schwienhorst <i>et al.</i> (2001)
Solar ν_e-e^-	Super-Kamiokande	$\mu_S(E_\nu \gtrsim 5 \text{ MeV}) < 1.1 \times 10^{-10} \mu_B$	90%	Liu <i>et al.</i> (2004)
	● Borexino	$\mu_S(E_\nu \lesssim 1 \text{ MeV}) < 5.4 \times 10^{-11} \mu_B$	90%	Arpesella <i>et al.</i> (2008)

C. Giunti, A. Studenikin, [arXiv: 1403.6344](https://arxiv.org/abs/1403.6344)

... A remark on electric charge of ν ... Beyond Standard Model...

ν neutrality $Q=0$ is attributed to

gauge invariance
+
anomaly cancellation constraints

imposed in SM of electroweak interactions

*Foot, Joshi, Lew, Volkas, 1990;
Foot, Lew, Volkas, 1993;
Babu, Mohapatra, 1989, 1990*



...General proof:

$$SU(2)_L \times U(1)_Y$$

$$Q = I_3 + \frac{Y}{2}$$

In SM :

In SM (without ν_R) triangle anomalies cancellation constraints \xrightarrow{Y} certain relations among particle hypercharges Y , that is enough to fix all Y so that they, and consequently Q , are quantized



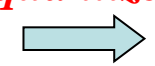
$Q=0$

$Q=0$ is proven also by direct calculation in SM within different gauges and methods

... However, strict requirements for Q quantization may disappear in extensions of standard $SU(2)_L \times U(1)_Y$ EW model if ν_R with $Y \neq 0$ are included : in the absence of Y quantization electric charges Q gets dequantized

*Bardeen, Gastmans, Lautrup, 1972;
Cabral-Rosetti, Bernabeu, Vidal, Zepeda, 2000;
Beg, Marciano, Ruderman, 1978;
Marciano, Sirlin, 1980; Sakakibara, 1981;
M.Dvornikov, A.S., 2004 (for extended SM in one-loop calculations)*

millicharged ν



Experimental limits for different effective neutrino millicharges

Limit	Method	Reference
$ \mathbf{q}_{\nu\tau} \lesssim 3 \times 10^{-4} e$	SLAC e^- beam dump	Davidson <i>et al.</i> (1991)
$ \mathbf{q}_{\nu\tau} \lesssim 4 \times 10^{-4} e$	BEBC beam dump	Babu <i>et al.</i> (1994)
$ \mathbf{q}_{\nu} \lesssim 6 \times 10^{-14} e$	Solar cooling (plasmon decay)	Raffelt (1999a)
$ \mathbf{q}_{\nu} \lesssim 2 \times 10^{-14} e$	Red giant cooling (plasmon decay)	Raffelt (1999a)
$ \mathbf{q}_{\nu_e} \lesssim 3 \times 10^{-21} e$	Neutrality of matter	Raffelt (1999a)
$ \mathbf{q}_{\nu_e} \lesssim 3.7 \times 10^{-12} e$	Nuclear reactor	Gninenko <i>et al.</i> (2007)
$ \mathbf{q}_{\nu_e} \lesssim 1.5 \times 10^{-12} e$	Nuclear reactor	Studenikin (2013)

C.Giunti, A.Studenikin, [arXiv: 1403.6344](https://arxiv.org/abs/1403.6344)

Bounds on ν millicharge q_ν from μ_ν

A.S.,
arXiv: 1302.1168,
18 March, 2014

ν - e cross-section

$$\left(\frac{d\sigma}{dT}\right)_{\nu-e} = \left(\frac{d\sigma}{dT}\right)_{SM} + \left(\frac{d\sigma}{dT}\right)_{\mu_\nu} + \left(\frac{d\sigma}{dT}\right)_{q_\nu}$$

two not seen contributions:

$$\left(\frac{d\sigma}{dT}\right)_{\mu_\nu^a} \approx \pi\alpha^2 \frac{1}{m_e^2 T} \left(\frac{\mu_\nu^a}{\mu_B}\right)^2$$

$$\left(\frac{d\sigma}{dT}\right)_{q_\nu} \approx 2\pi\alpha \frac{1}{m_e T^2} q_\nu^2$$

Bounds on q_ν from

$$R = \frac{\left(\frac{d\sigma}{dT}\right)_{q_\nu}}{\left(\frac{d\sigma}{dT}\right)_{\mu_\nu^a}} = \frac{2m_e}{T} \frac{\left(\frac{q_\nu}{e_0}\right)^2}{\left(\frac{\mu_\nu^a}{\mu_B}\right)^2} \ll 1$$

... no
observable
effects of
New
Physics

Constraints on μ_ν from GEMMA: Constraints on q_ν

now $\mu_\nu^a < 2.9 \times 10^{-11} \mu_B$ ($T \sim 2.8$ keV)

$$|q_\nu| < 1.5 \times 10^{-12} e_0$$

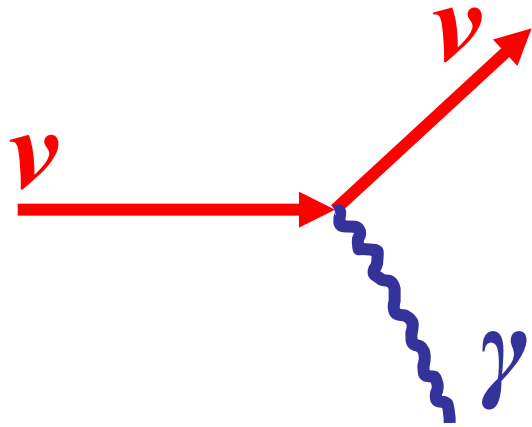
2015 (expected) $\mu_\nu^a \sim 1.5 \times 10^{-11} \mu_B$ ($T = 1.5$ keV)

$$|q_\nu| < 3.7 \times 10^{-13} e_0$$

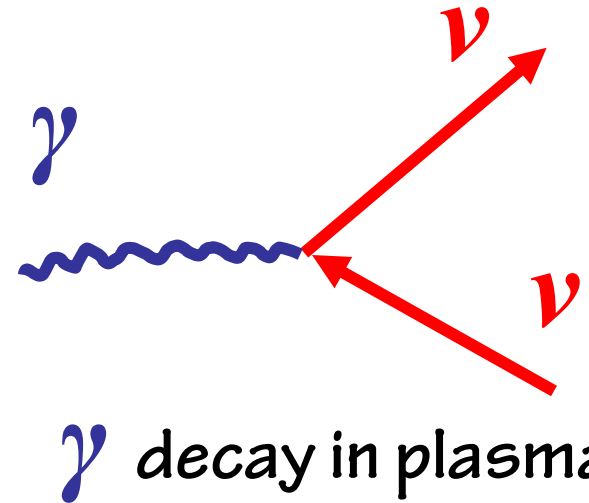
2018 (expected) $\mu_\nu^a \sim 0.9 \times 10^{-12} \mu_B$ ($T = 350$ eV)

$$|q_\nu| < 1.8 \times 10^{-13} e_0$$

ν electromagnetic interactions

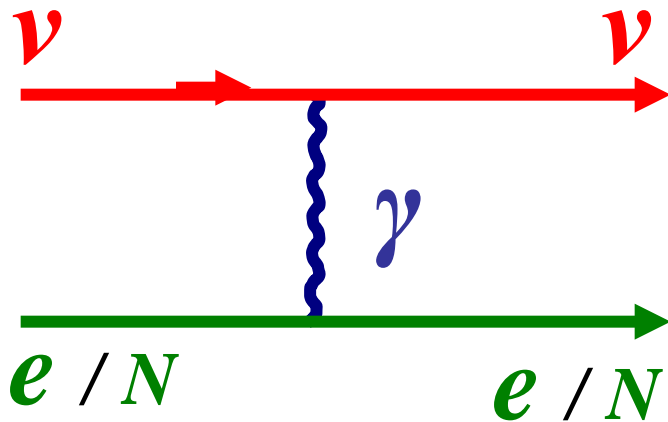


ν decay, Cherenkov radiation

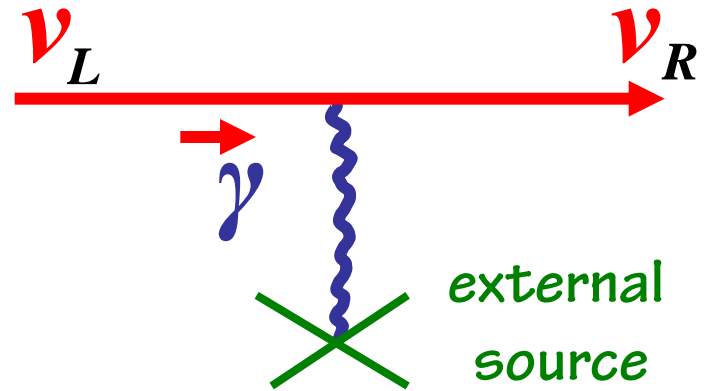


γ decay in plasma

!!!



Scattering

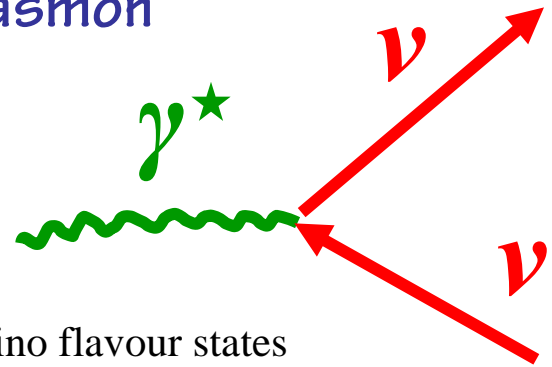


Spin precession

Astrophysical bound on μ_ν

G.Raffelt, PRL 1990

comes from cooling of **red giant** stars by plasmon decay
 decay $\gamma^* \longrightarrow \nu \bar{\nu}$



$$L_{int} = \frac{1}{2} \sum_{a,b} \left(\mu_{a,b} \bar{\psi}_a \sigma_{\mu\nu} \psi_b + \epsilon_{a,b} \bar{\psi}_a \sigma_{\mu\nu} \gamma_5 \psi_b \right)$$

neutrino flavour states

$$\epsilon_\alpha k^\alpha = 0$$

Matrix element

$$|M|^2 = M_{\alpha\beta} p^\alpha p^\beta, \quad M_{\alpha\beta} = 4\mu^2 (2k_\alpha k_\beta - 2k^2 \epsilon_\alpha^* \epsilon_\beta - k^2 g_{\alpha,\beta}),$$

Decay rate

$$\Gamma_{\gamma \rightarrow \nu \bar{\nu}} = \frac{\mu^2 (\omega^2 - k^2)^2}{24\pi \omega}$$

= 0 in vacuum $\omega = k$

In the classical limit γ^* - like a massive particle with $\omega^2 - k^2 = \omega_{pl}^2$

Energy-loss rate per unit volume

$$Q_\mu = g \int \frac{d^3k}{(2\pi)^3} \omega f_{BE} \Gamma_{\gamma \rightarrow \nu \bar{\nu}}$$

distribution function of plasmons

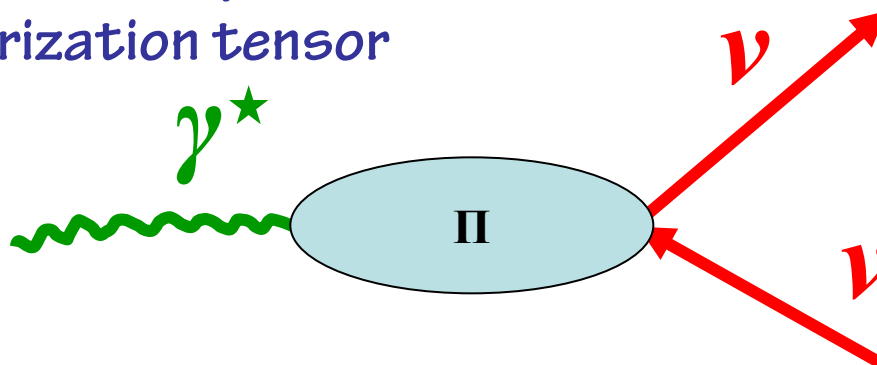
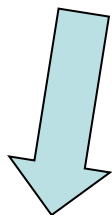
$$\mu^2 \rightarrow \sum_{a,b} (|\mu_{a,b}|^2 + |\epsilon_{a,b}|^2)$$

Astrophysical bound on μ_ν

$$Q_\mu = g \int \frac{d^3k}{(2\pi)^3} \omega f_{BE} \Gamma_{\gamma \rightarrow \nu\bar{\nu}}$$

Energy-loss rate
per unit volume

Magnetic moment **plasmon** decay
enhances the Standard Model photo-neutrino
cooling by photon polarization tensor



more fast star cooling

In order not to delay helium ignition ($\leq 5\%$ in Q)

... best
astrophysical
limit on

$$\mu \leq 3 \times 10^{-12} \mu_B$$

G.Raffelt, PRL 1990

✓ magnetic moment...

$$\mu^2 \rightarrow \sum_{a,b} (|\mu_{a,b}|^2 + |\epsilon_{a,b}|^2)$$

3.10

Dobroliubov, Ignatiev (1990); Babu, Volkas (1992);

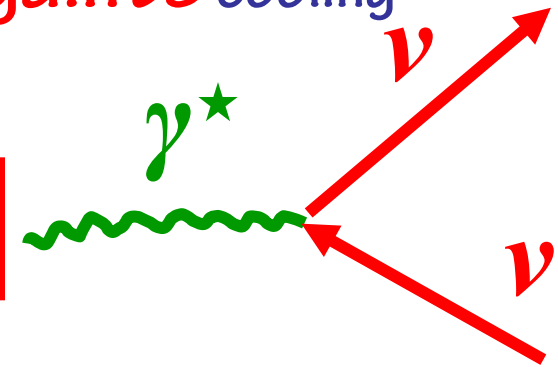
↙ Mohapatra, Nussinov (1992) ...

- Constraints on neutrino millicharge from red giants cooling



Interaction Lagrangian

$$L_{int} = -iq_\nu \bar{\psi}_\nu \gamma^\mu \psi_\nu A^\mu$$



↖ millicharge

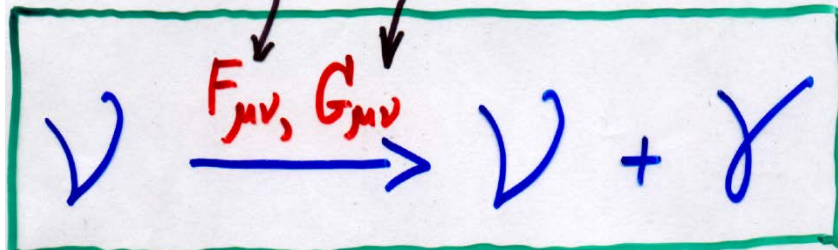
Decay rate

$$\Gamma_{q_\nu} = \frac{q_\nu^2}{12\pi} \omega_{pl} \left(\frac{\omega_{pl}}{\omega} \right)$$

- $q_\nu \leq 2 \times 10^{-14} e$...to avoid helium ignition in low-mass red giants Halt, Raffelt, Weiss, PRL1994
- $q_\nu \leq 3 \times 10^{-17} e$... absence of anomalous energy-dependent dispersion of SN1987A ✓ signal, most model independent
- ... from “charge neutrality” of neutron... $q_\nu \leq 3 \times 10^{-21} e$

• New mechanism of electromagnetic radiation

"Spin light of neutrino"
in matter and
electromagnetic fields



A.Lobanov, A.Studenikin,
Phys.Lett. B 564 (2003) 27
Phys.Lett. B 601 (2004) 171

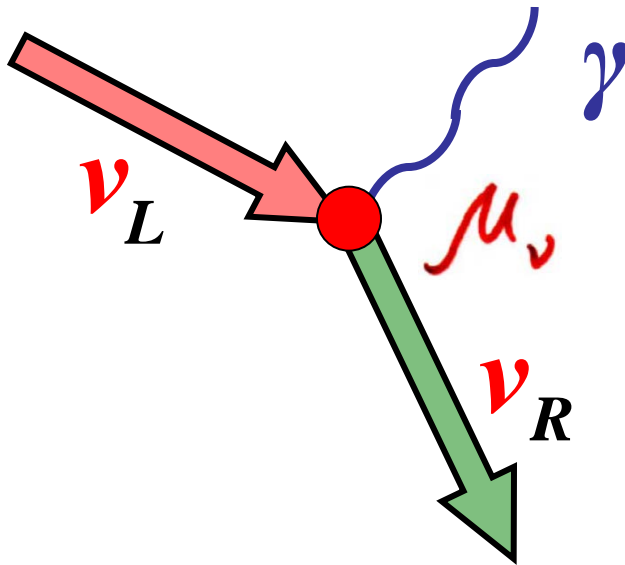
Studenikin, A.Ternov,
Phys.Lett. B 608 (2005) 107

A.Grigoriev, A.S., Ternov,
Phys.Lett. B 622 (2005) 199

Studenikin,
J.Phys.A: Math.Gen. 39 (2006) 6769,
J.Phys.A: Math.Theor. 41 (2008) 16402,

A.Grigoriev, A.Lokhov,
A.Studenikin, A.Ternov,
Nuovo Cim. 35 C (2012) 57,
Phys.lett.B 718 (2012) 512

Neutrino – photon coupling



broad neutrino lines
account for interaction
with environment

“Spin light of neutrino in matter”

$SL\nu$

- ... within the quantum treatment based on
method of exact solutions ...

in extreme environments

- A. Studenikin,
“Quantum treatment of neutrino in background matter”,
J. Phys. A: Math. Gen. 39 (2006) 6769–6776
- “Method of wave equations exact solutions in studies of neutrinos and electron interactions in dense matter”,
J.Phys.A: Math.Theor. 41 (2008) 164047
- “Neutrinos and electrons in background matter: a new approach”,
Ann.Fond. de Broglie 31 (2006) 289-316

...«method of exact solutions»

- ν quantum states in dense magnetized matter

... new effect of ...

Spin Light of ν
in matter



... phenomenological

ν energy quantization in rotating matter

consequences in astrophysics (pulsars)

ν in matter treated within
«method of exact solutions»

(Dirac equation with matter potential for ν)

✓ energy quantization in rotating magnetized media

Grigoriev, Savochkin, Studenikin, Russ.Phys.J. 50 (2007) 845

Studenikin, J.Phys. A: Math.Theor. 41 (2008) 164047

Balantsev, Popov, Studenikin,

J.Phys. A:Math.Theor. 44 (2011) 255301

Balantsev, Studenikin, Tokarev, Phys.Part.Nucl. 43 (2012), 727

Phys.Atom.Nucl. 76 (2013) 489

Studenikin, Tokarev, Nucl.Phys.B (2014)

Millicharged ψ in rotating magnetized matter

Balatsev, Tokarev, Studenikin,
 Phys.Part.Nucl., 2012,
 Phys.Atom.Nucl., Nucl.Phys. B, 2013,
 Studenikin, Tokarev, Nucl.Phys.B (2014) •

Modified Dirac equation for ψ wave function

$$\left(\gamma_\mu (p^\mu + q_0 A^\mu) - \frac{1}{2} \gamma_\mu (c_l + \gamma_5) f^\mu - \frac{i}{2} \mu \sigma_{\mu\nu} F^{\mu\nu} - m \right) \Psi(x) = 0$$

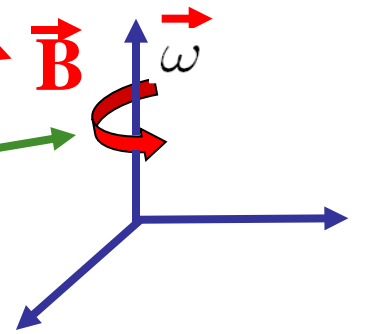
external magnetic field

$$V_m = \frac{1}{2} \gamma_\mu (c_l + \gamma_5) f^\mu \quad c_l = 1$$

matter potential

rotating matter

$$f^\mu = -G n_n (1, -\epsilon y \omega, \epsilon x \omega, 0)$$



*rotation
 angular
 frequency*

✓ energy is quantized in rotating matter

$$G = \frac{G_F}{\sqrt{2}}$$

$$p_0 = \sqrt{p_3^2 + 2N|2Gn_n\omega - \epsilon q_\nu B| + m^2} - Gn_n - q\phi$$

$$N = 0, 1, 2, \dots$$

integer number

matter rotation
frequency

scalar potential
of electric field

✓ energy is quantized in rotating matter
like electron energy in magnetic field
(Landau energy levels):

$$p_0^{(e)} = \sqrt{m_e^2 + p_3^2 + 2\gamma N}, \quad \gamma = eB, \quad N = 0, 1, 2, \dots$$

- In quasi-classical approach
- ✓ quantum states in rotating matter
- ✓ motion in circular orbits

$$R = \int_0^\infty \Psi_L^\dagger r \Psi_L d\mathbf{r} = \sqrt{\frac{2N}{|2Gn_n\omega - \epsilon q_0 B|}}$$

- due to **effective Lorentz force**

Studenikin,
J.Phys. A
(2008)

$$\mathbf{F}_{eff} = q_{eff} \mathbf{E}_{eff} + q_{eff} [\boldsymbol{\beta} \times \mathbf{B}_{eff}]$$

$$q_{eff} \mathbf{E}_{eff} = q_m \mathbf{E}_m + q_0 \mathbf{E} \quad q_{eff} \mathbf{B}_{eff} = |q_m B_m + q_0 B| \mathbf{e}_z$$

where

$$q_m = -G, \quad \mathbf{E}_m = -\nabla n_n, \quad \mathbf{B}_m = 2n_n \boldsymbol{\omega}$$

- matter induced “charge”, “electric” and fields “magnetic”

... we predict :

$$E \sim 1 \text{ eV}$$

1) low-energy ν are trapped in circular orbits inside rotating neutron stars

$$R = \sqrt{\frac{2N}{Gn\omega}} < R_{NS} = 10 \text{ km}$$

$$\begin{aligned} R_{NS} &= 10 \text{ km} \\ n &= 10^{37} \text{ cm}^{-3} \\ \omega &= 2\pi \times 10^3 \text{ s}^{-1} \end{aligned}$$

2) rotating neutron stars as

filters for low-energy relic ν ?

$$T_\nu \sim 10^{-4} \text{ eV}$$

... we predict :

A.Studenikin, I.Tokarev,
Nucl.Phys.B (2014)

3) high-energy ν are deflected inside
a rotating **astrophysical transient sources**
(GRBs, SNe, AGNs)

absence of light in correlation with
 ν signal reported by ANTARES Coll.

M.Ageron et al,
Nucl.Instrum.Meth. A692 (2012) 184

Millicharged ν as star rotation engine

- Single ν generates feedback force with projection on rotation plane

- $F = (q_0 B + 2Gn_n \omega) \sin \theta$

single ν torque

- $M_0(t) = \sqrt{1 - \frac{r^2(t)\Omega^2 \sin^2 \theta}{4}} F r(t) \sin \theta$

total N_ν torque

$$M(t) = \frac{N_\nu}{4\pi} \int M_0(t) \sin \theta d\theta d\varphi$$

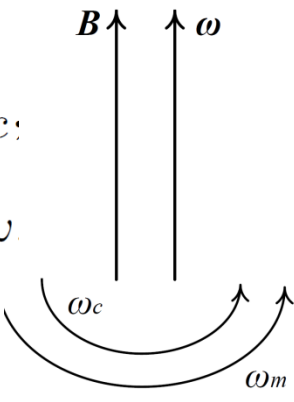
- Shift of star angular velocity

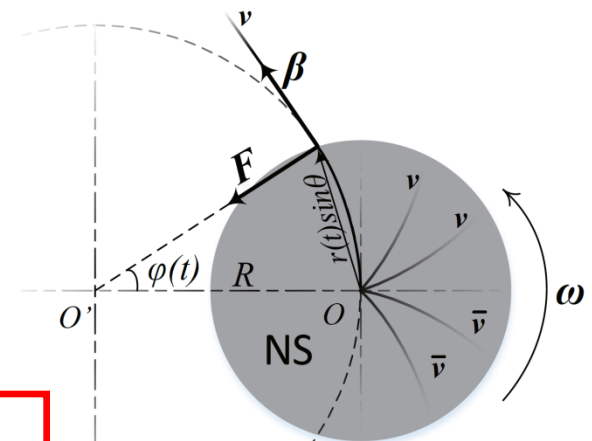
$$|\Delta\omega| = \frac{5N_\nu}{6M_S} (q_0 B + 2Gn_n \omega_0)$$

$$\Delta\omega = \omega - \omega_0$$

$$\Omega = \omega_m + \omega_c;$$

$$\omega_m = \frac{2Gn_n}{p_0 + Gn_n} \omega$$

$$\omega_c = \frac{q_0 B}{p_0 + Gn_n}$$




A. Studenikin, I. Tokarev,
Nucl. Phys. B (2014)

• ν Star Turning mechanism (ν ST)

A.Studenikin, I.Tokarev, Nucl.Phys.B (2014)

Escaping ν s move on curved orbits inside magnetized rotating star and feedback of effective Lorentz force should effect initial star rotation

- **New** astrophysical constraint on ν millicharge

$$\frac{|\Delta\omega|}{\omega_0} = 7.6\varepsilon \times 10^{18} \left(\frac{P_0}{10 \text{ s}} \right) \left(\frac{N_\nu}{10^{58}} \right) \left(\frac{1.4M_\odot}{M_S} \right) \left(\frac{B}{10^{14}G} \right)$$

- $|\Delta\omega| < \omega_0$! ...to avoid contradiction of ν ST impact with observational data on pulsars ...

$$q_0 < 1.3 \times 10^{-19} e_0$$

... **best** astrophysical bound ...

Spin light of e in dense neutrino fluxes

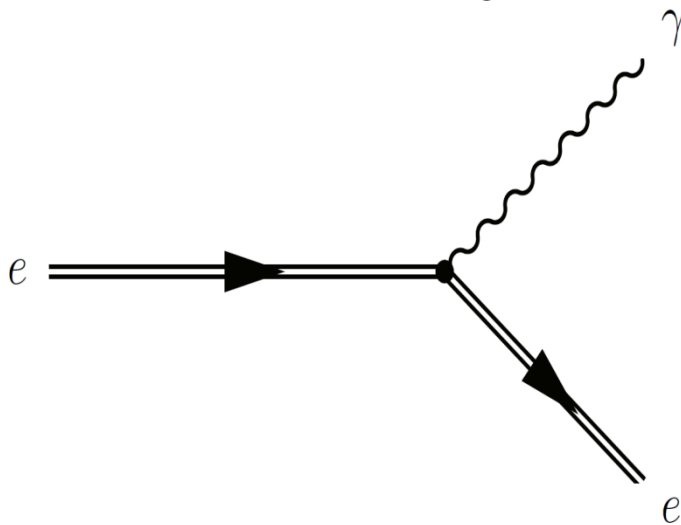
SLe_ν
I. Balantsev,
A. Studenikin,
(2014)

Dirac eq for e in dense relativistic flux of ν

$$\left(\gamma_\mu p^\mu + \gamma_\mu \frac{c + \delta_e \gamma^5}{2} f^\mu - m \right) \Psi(x) = 0$$

$$c = \delta_e - 12 \sin^2 \theta_W$$

$$\delta_e = \frac{n_\mu + n_\tau - n_e}{n}$$



$$f^\mu = G(n, 0, 0, n)$$

background matter
(ν flux) potential

Conclusion

✓ ***e.m. vertex function*** \Rightarrow **4 form factors** 

charge dipole magnetic and electric

● $\Lambda_\mu(q) = f_Q(q^2)\gamma_\mu + f_M(q^2)i\sigma_{\mu\nu}q^\nu + f_E(q^2)\sigma_{\mu\nu}q^\nu\gamma_5$
 $f_A(q^2)(q^2\gamma_\mu - q_\mu\not{q})\gamma_5$ ***anapole***

● ***EM properties*** \Rightarrow **a way to distinguish Dirac and Majorana** ✓

● **Standard Model with ν_R ($m_\nu \neq 0$):** $\mu_e = \frac{3eG_F}{8\sqrt{2}\pi^2} m_{\nu_e} \sim 3 \cdot 10^{-19} \mu_B \left(\frac{m_{\nu_e}}{1\text{eV}}\right)$

● ***In extensions of SM*** \rightarrow **enhancement of magnetic moment**, even **electrically millicharged**

● **Limits from reactor ν - e scattering experiments (2012):**

$\mu_\nu < 2.9 \times 10^{-11} \mu_B$

A.Beda et al. (GEMMA Coll.)

● **Limits from astrophysics, star cooling (1990):**

$\mu_\nu < 3 \times 10^{-12} \mu_B$

G.Raffelt

$|q_\nu| < 1.5 \times 10^{-12} e_0$

$q_0 < 1.3 \times 10^{-19} e_0$

ν ST mechanism