



DARK MATTER SEARCHES IN

IceCube

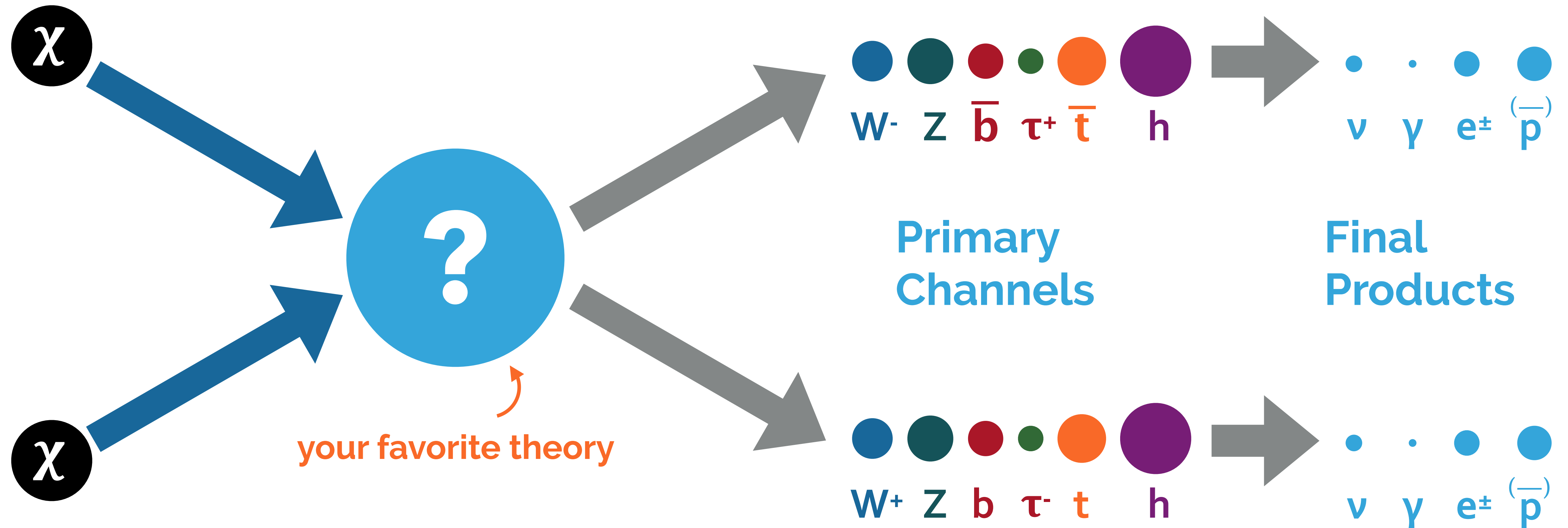
J. A. Aguilar on behalf of IceCube

Photo: Ian Reese

ULB

31st Rencontres de Blois 2019

iihe



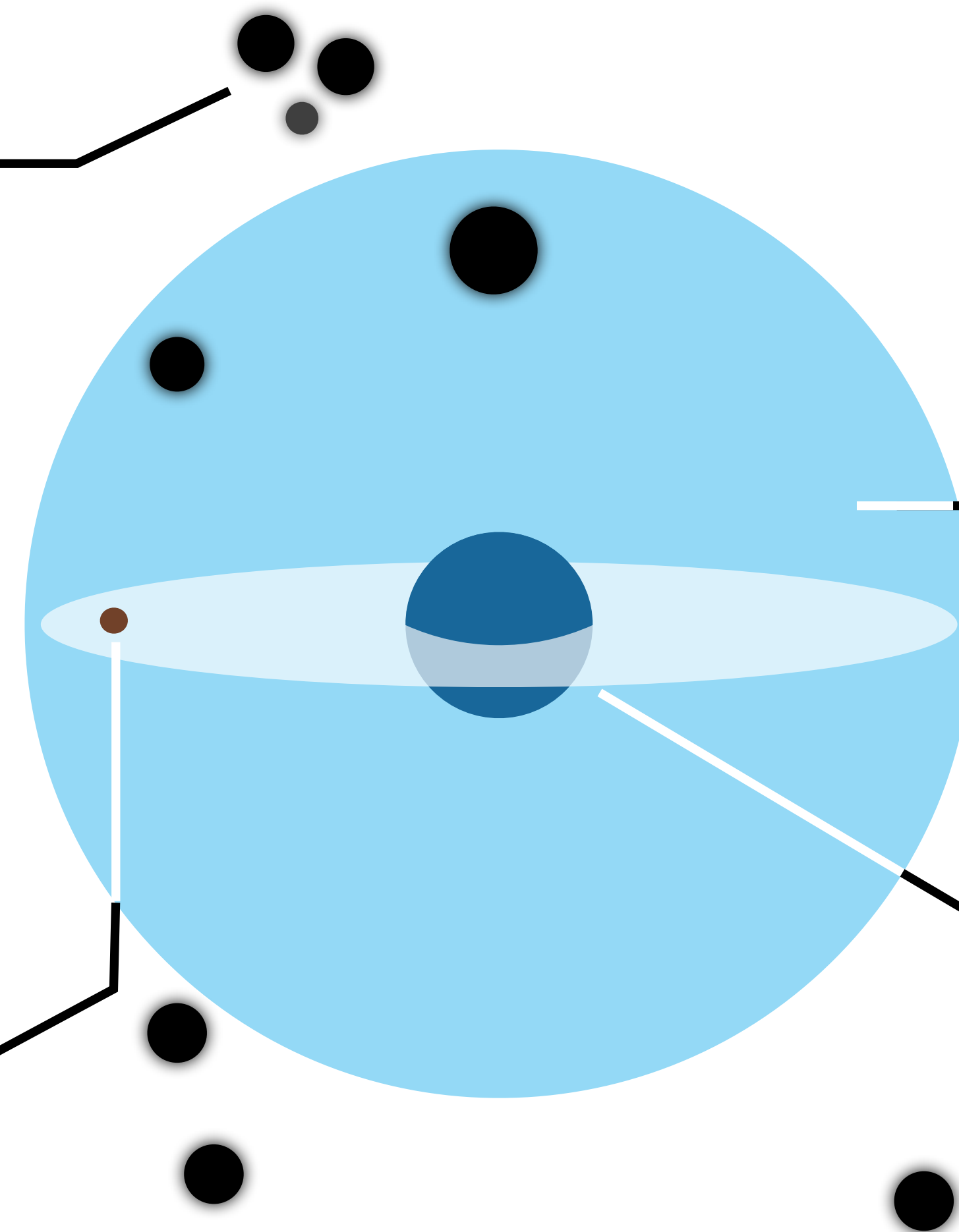
- No need of specialized detectors: **Gamma-ray telescopes, neutrino detectors, CR-experiments**
- Search for products of dark matter annihilation processes: **Focus on large reservoirs of dark matter**

Dwarf spheroidal Galaxies **Cluster of Galaxies**

Probe velocity-averaged DM
annihilation cross section $\langle v\sigma_A \rangle$

Local Sources (Sun, Earth)

Only accessible with neutrinos
Under equilibrium they can
probe σ_{SI} and σ_{SD}



Galactic Halo

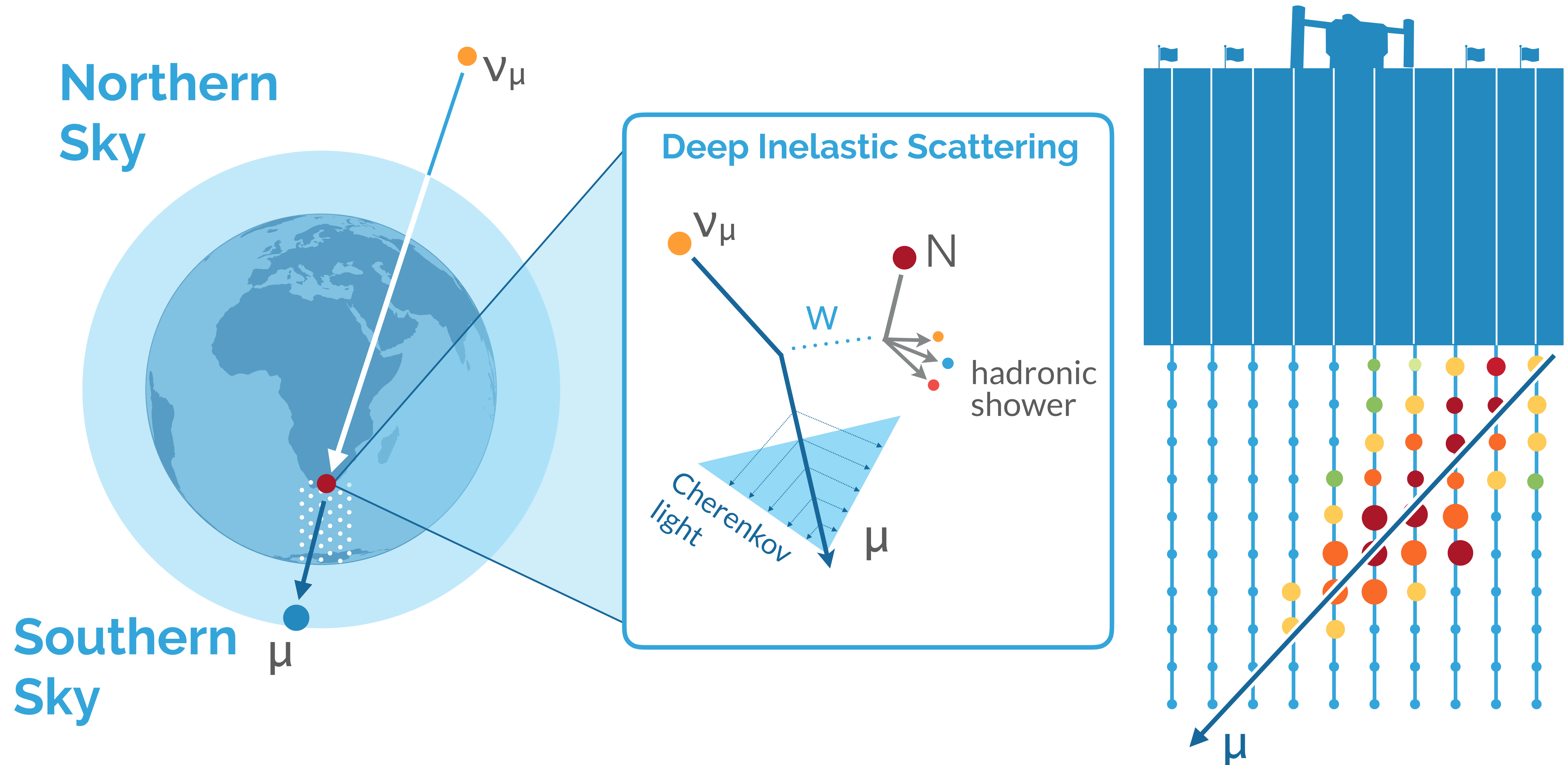
Probe velocity-averaged DM
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Galactic Center

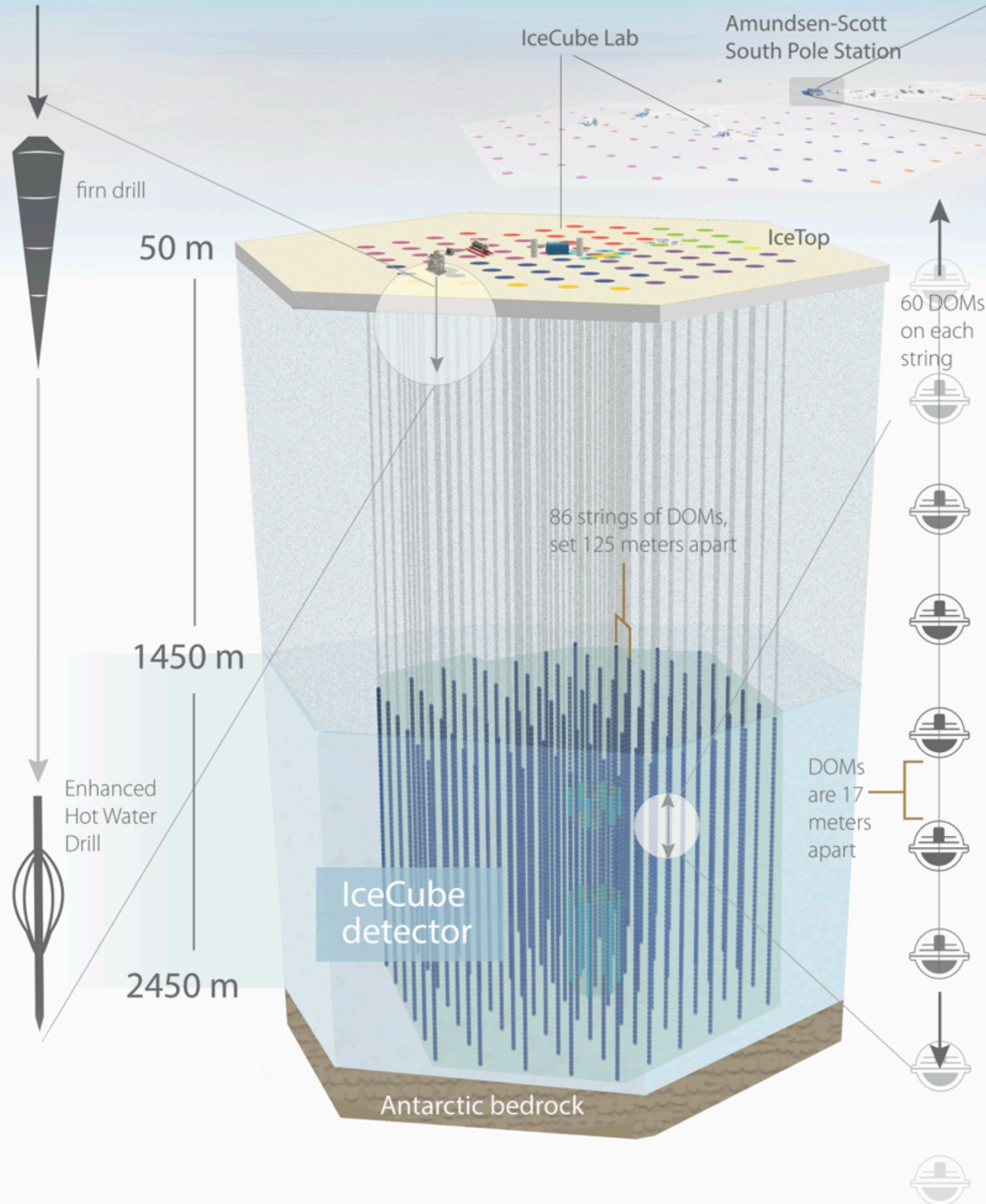
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annihilation cross section $\langle v\sigma_A \rangle$

Detection Principle How Do We Detect Neutrinos?

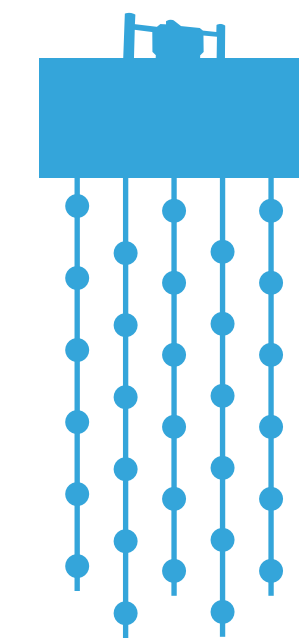
4



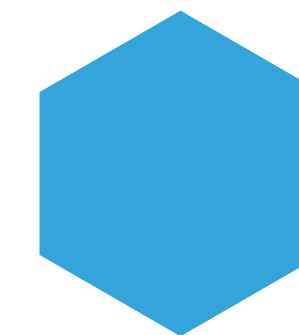
IceCube Neutrino Observatory



5,160 Digital Optical Modules (DOMs)



86 string with 60 DOMs each
6 denser strings called **DeepCore**



1 km² surface array with 324 DOMs: **IceTop**



Completion in December 2010

Neutrino Detectors **Scientific Scope**

6

And Cosmic Rays

SN Neutrinos

Astrophysical

MeV

GeV-TeV

<100 TeV

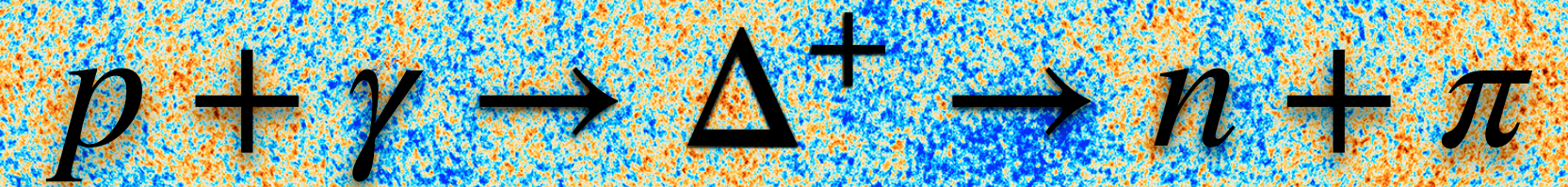
>100TeV

> 10⁶ TeV

Dark Matter

Atmos: Oscillations,
sterile ν , Prompt

Cosmogenic (GZK)



Neutrino Detectors **Scientific Scope**

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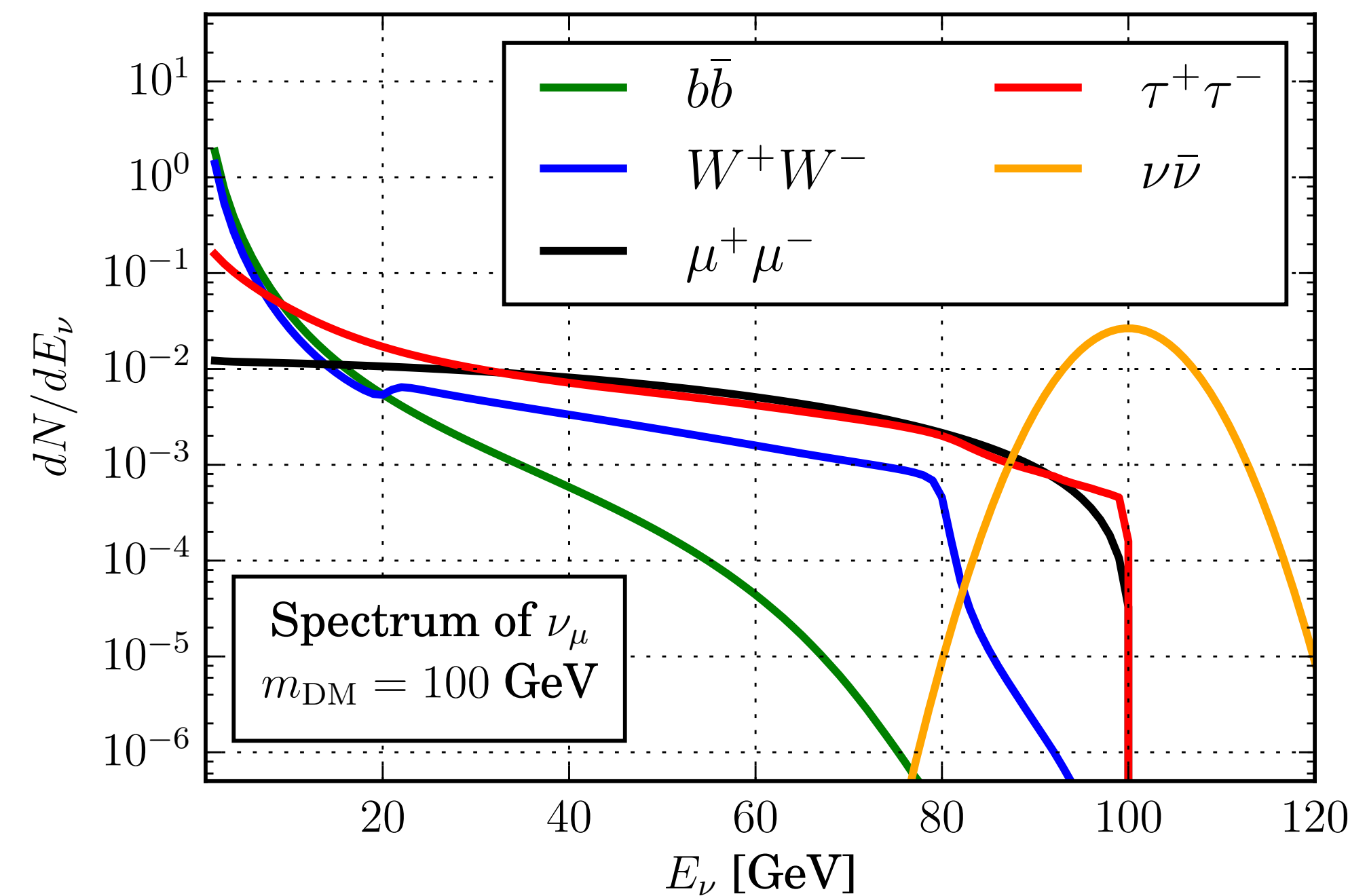
Cosmogenic (GZK)



$$\frac{d\Phi_\nu}{dE_\nu} = \frac{1}{4\pi} \frac{\langle \sigma_A v \rangle}{2m_\chi^2} \frac{dN_\nu}{dE_\nu} \int_0^{\Delta\Omega} d\Omega \int_{l.o.s} \rho_\chi^2(r(s, \Psi, \theta)) ds$$

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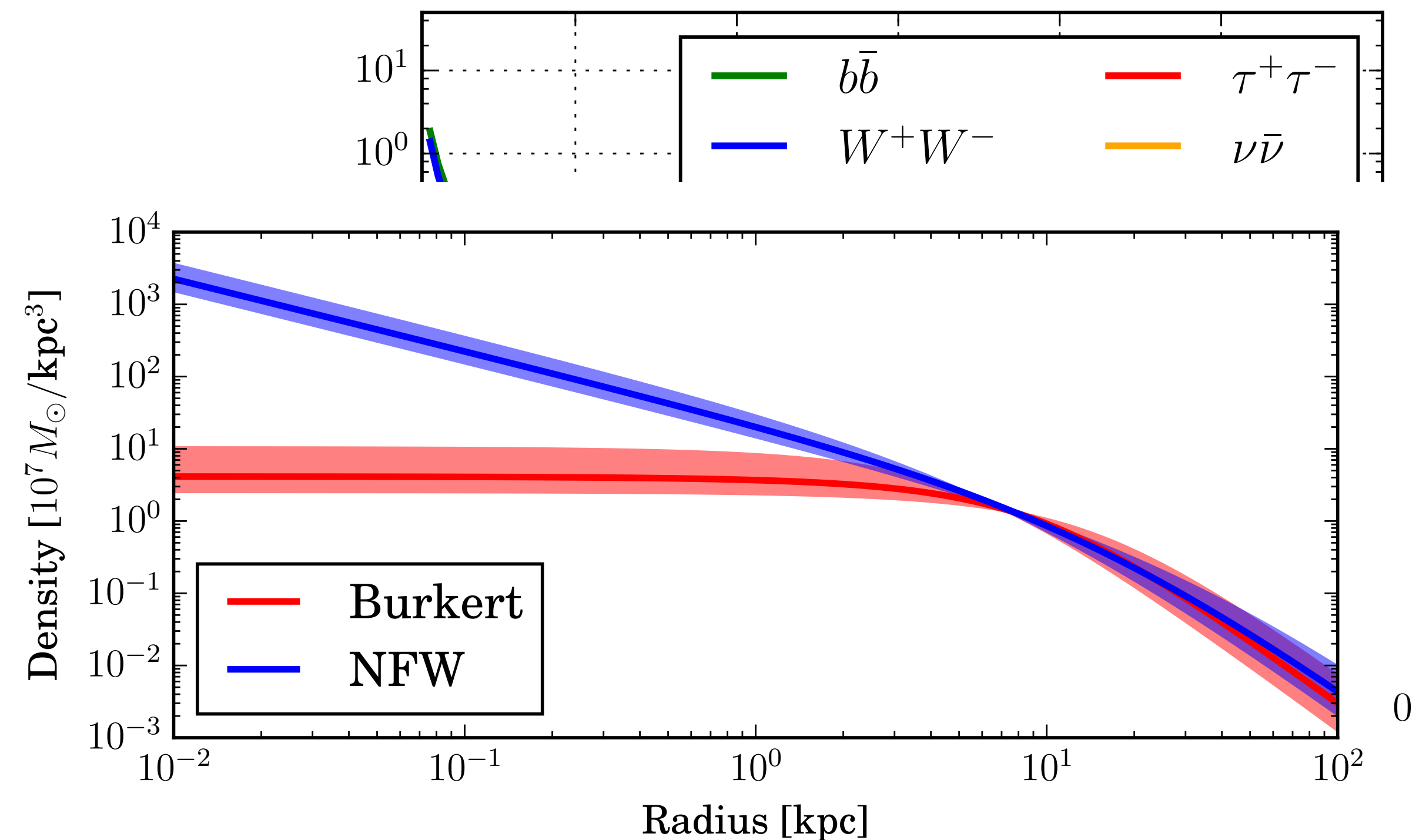
Theory input: SUSY?



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Theory input: SUSY?

Astrophysics input

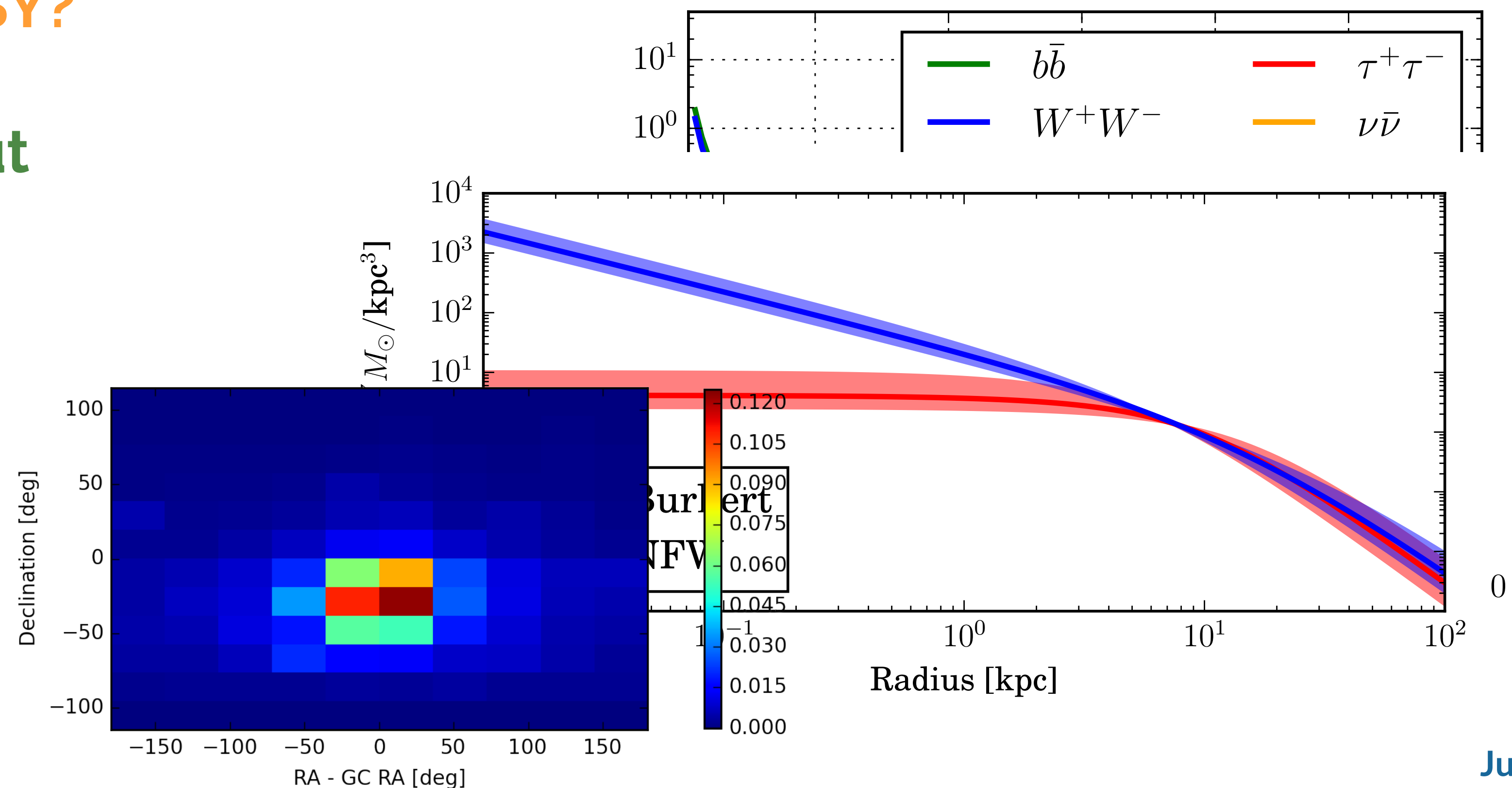


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Theory input: SUSY?

Astrophysics input

Measurement



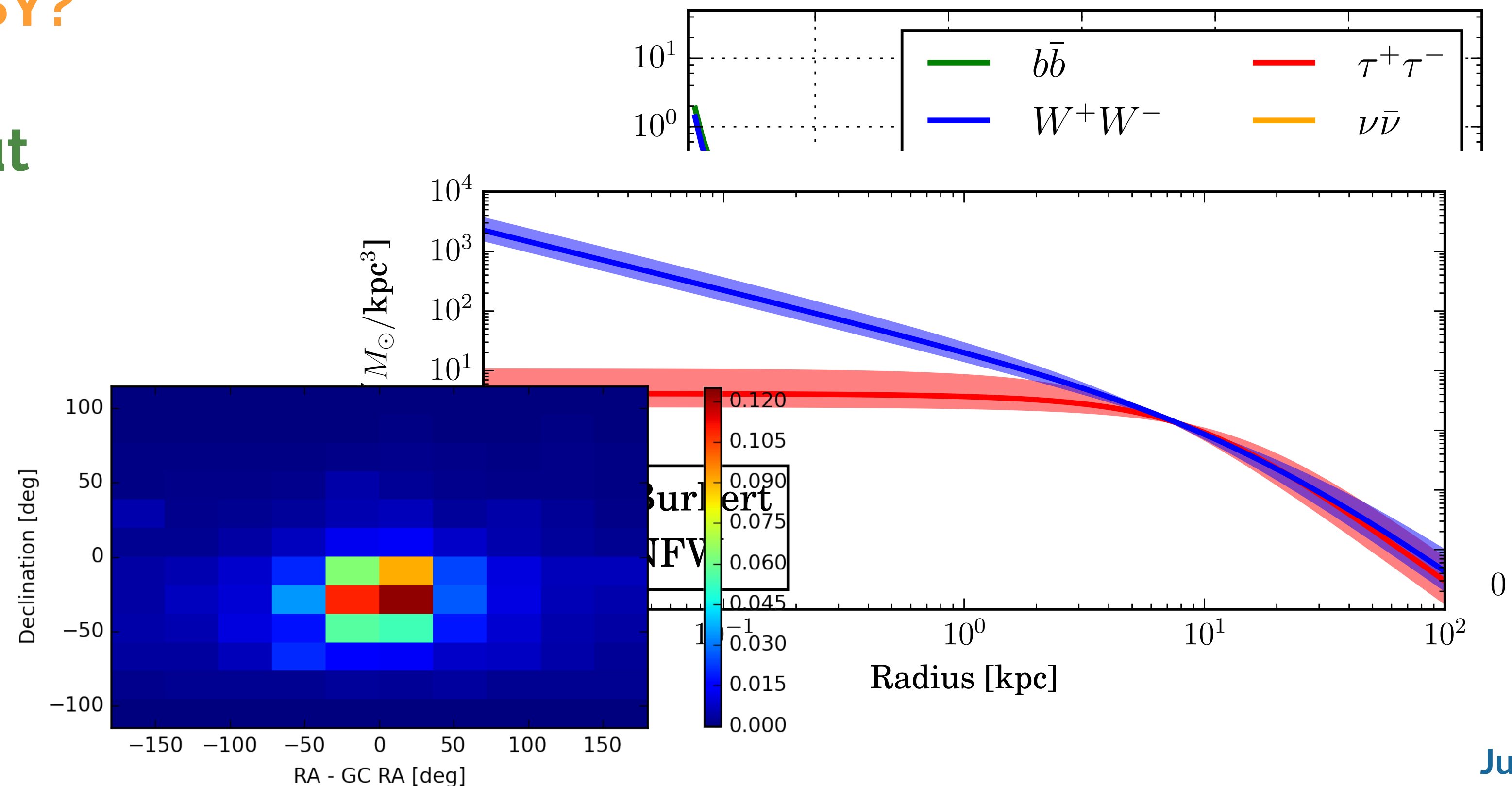
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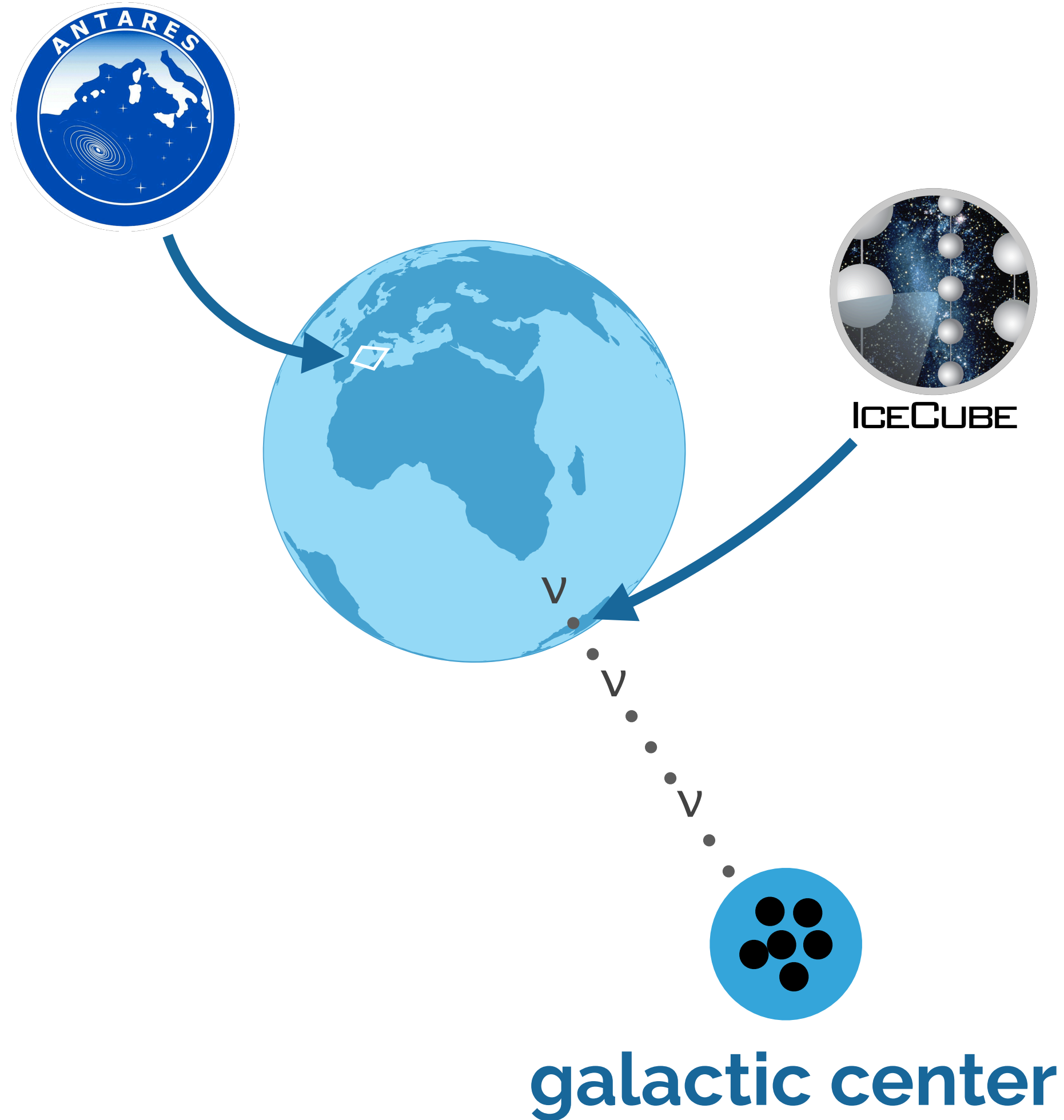
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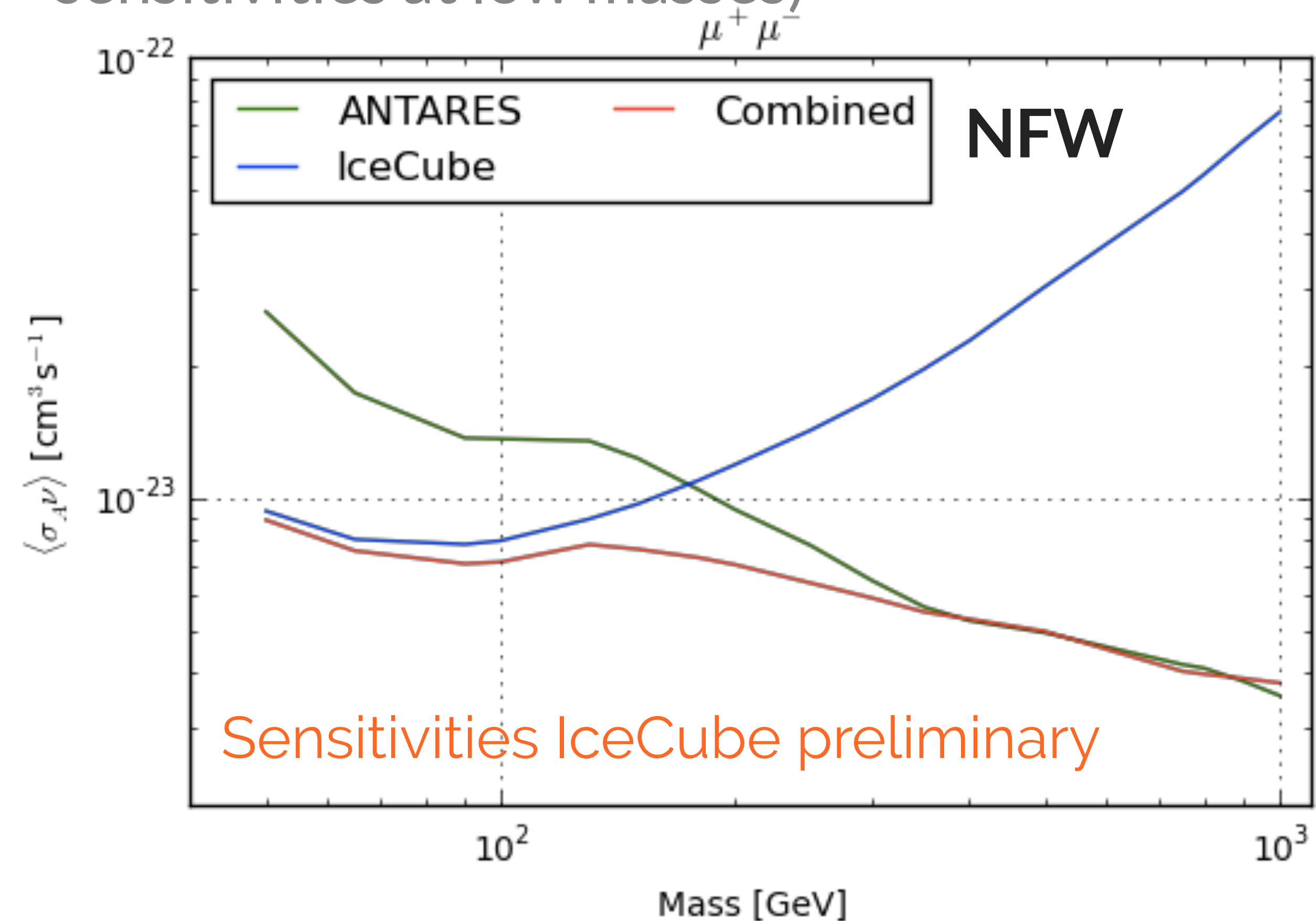
Measurement

Constrain!

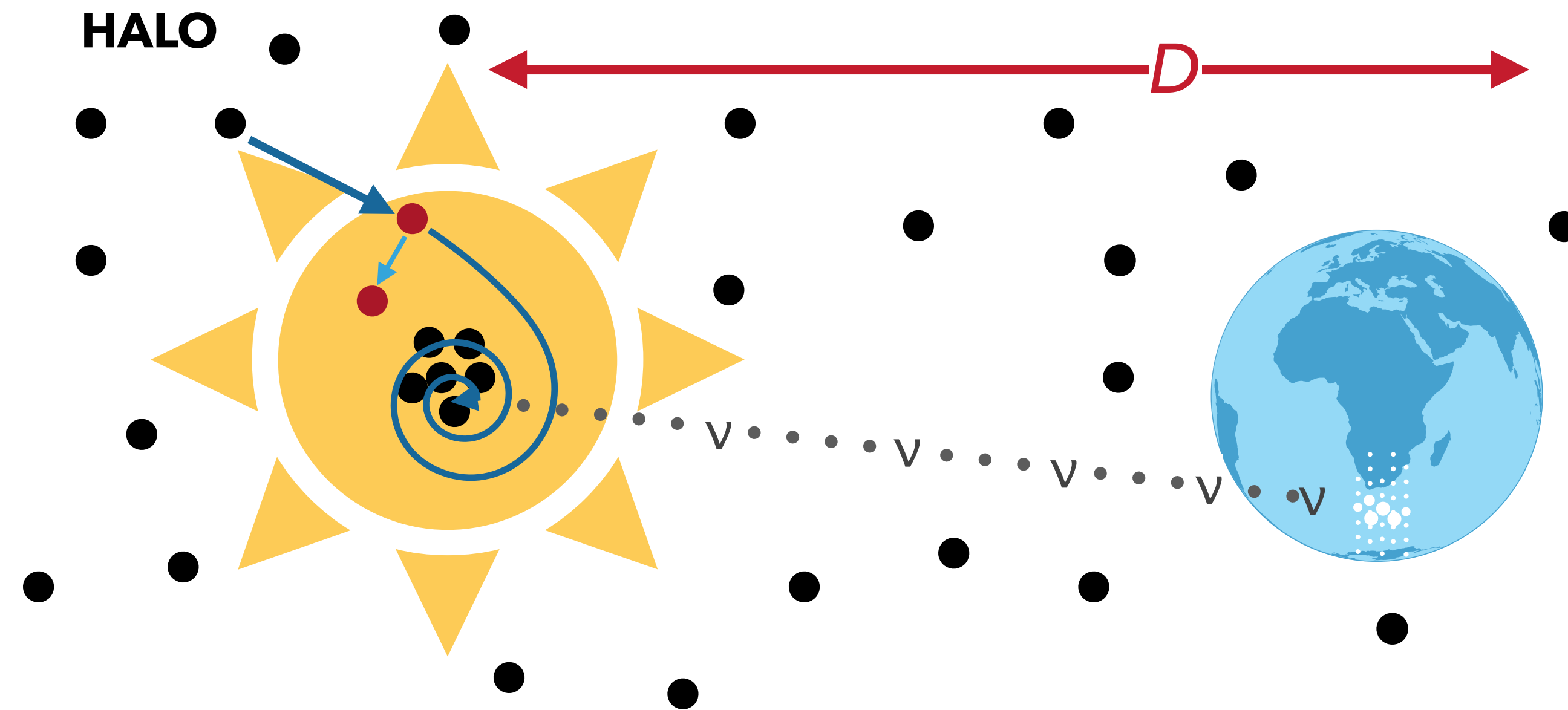




- ANTARES can use the Earth as a shield to observe the Galactic Center.
- IceCube needs a self-veto technique (better sensitivities at low masses)

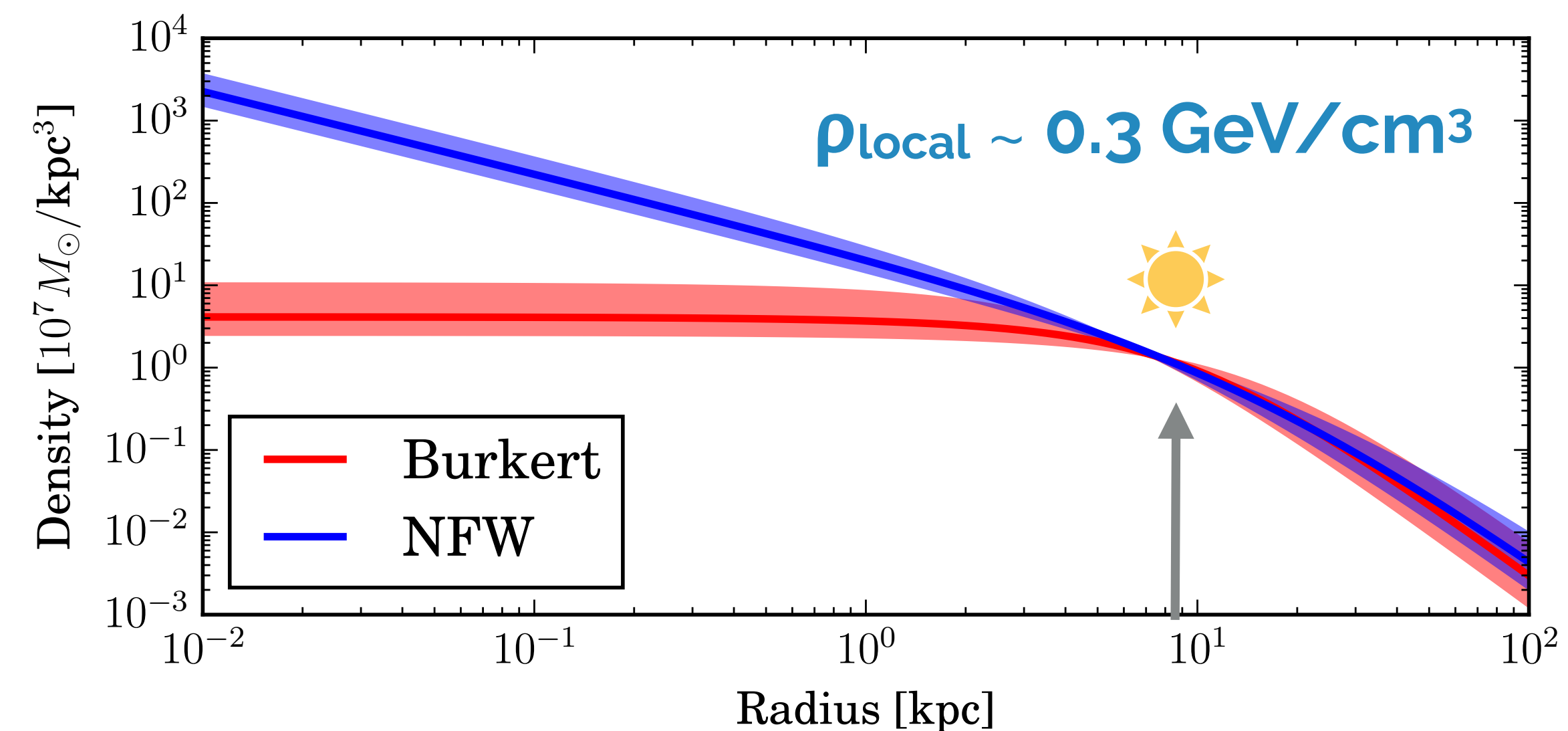


[Neutrino 2018, doi:10.5281/zenodo.1300924]



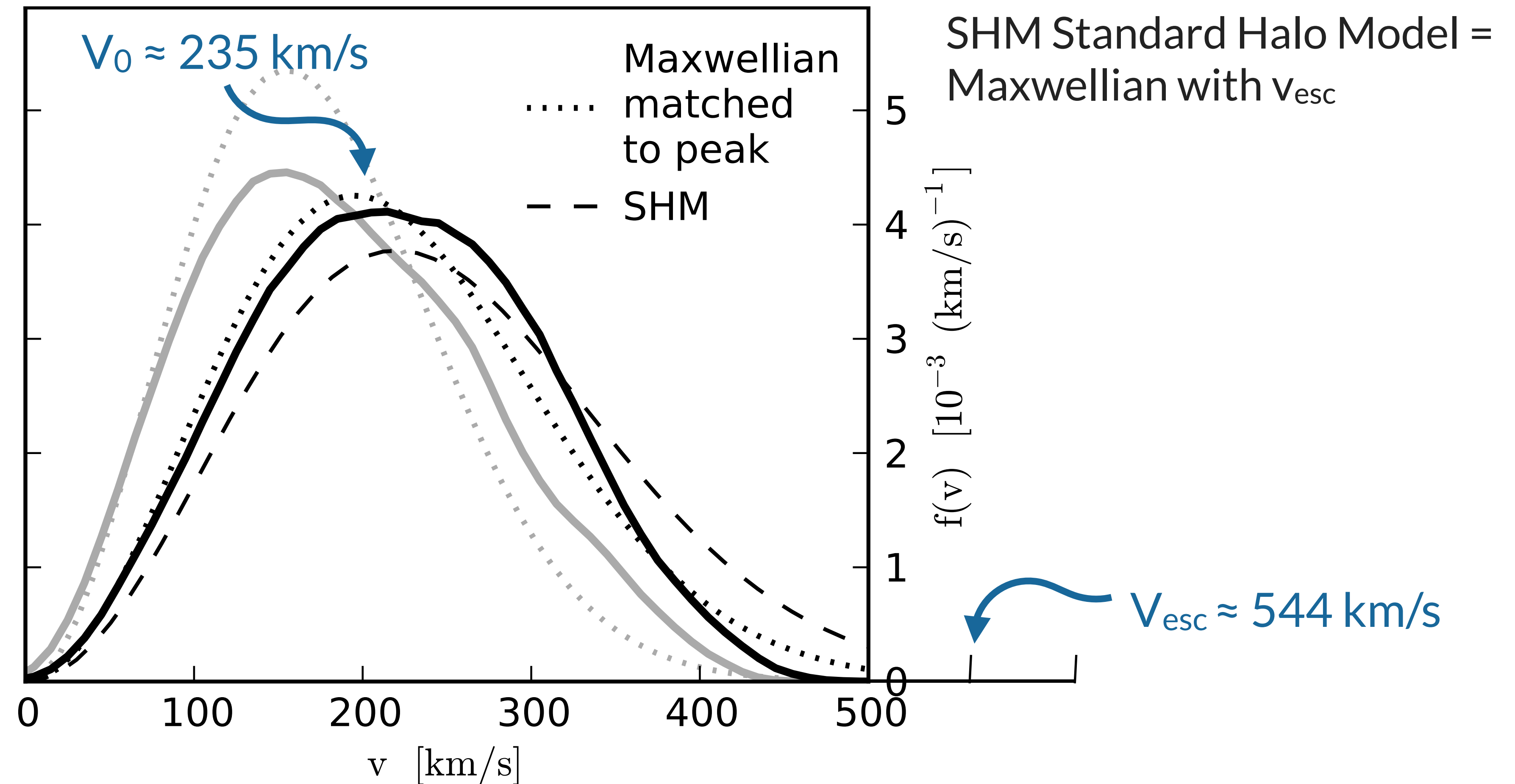
- Signal from the Sun or Earth in neutrinos cannot be mis-interpreted as an astrophysical source.
- Halo models agree in the Solar System.
- **Equilibrium** ($dN/dt = 0$) assumed for the Sun, not for Earth!

$$\frac{dN}{dt} = \underbrace{C_c}_{\text{capture } \sigma_{\chi-N}} - \underbrace{C_A}_{\text{annihilation } \sigma_A} N^2 - \underbrace{C_E N}_{\text{escape}}$$



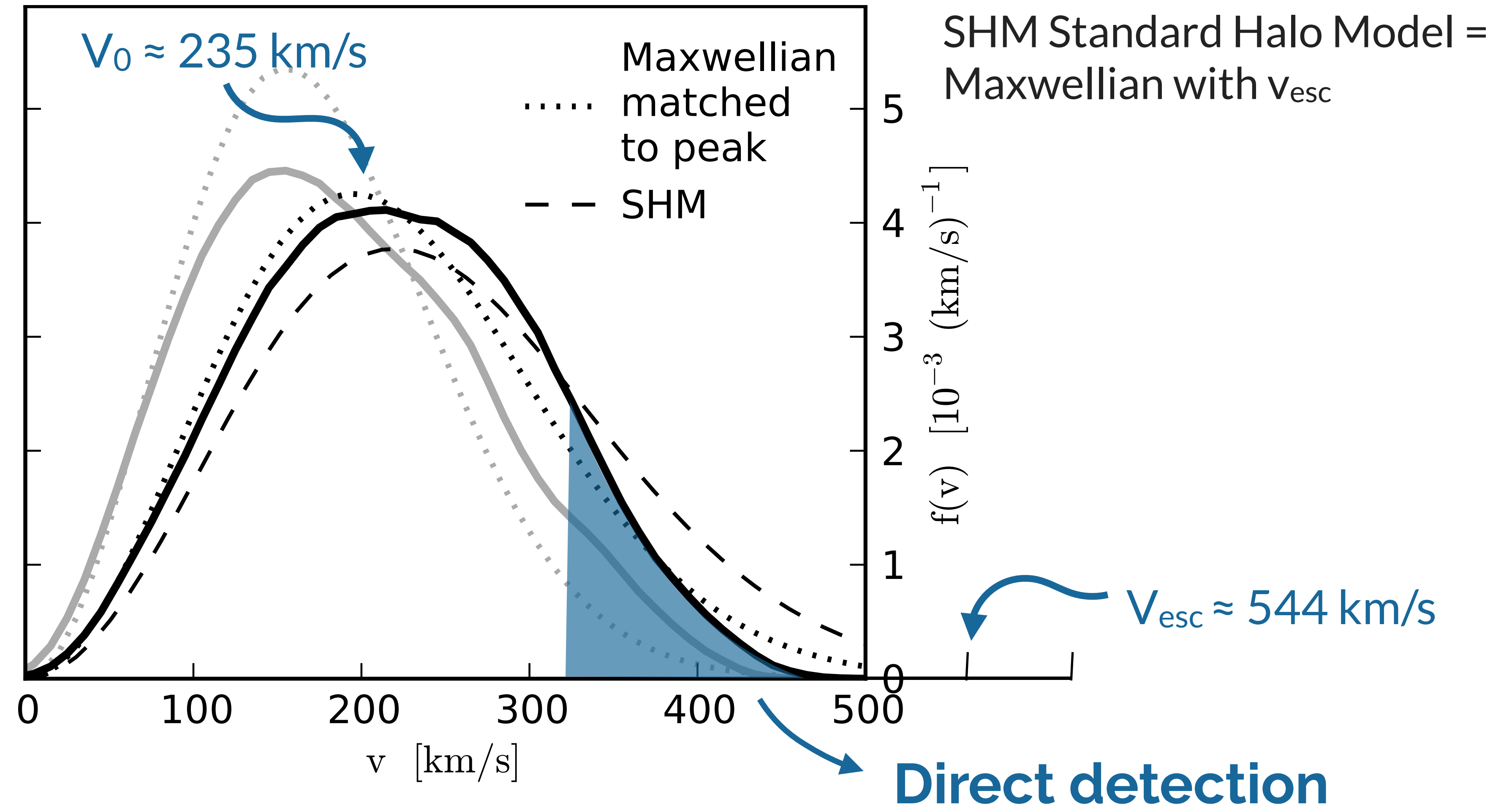
Heavy dark matter particles can only be captured at low velocities

Figure from <https://arxiv.org/pdf/1308.1703.pdf>



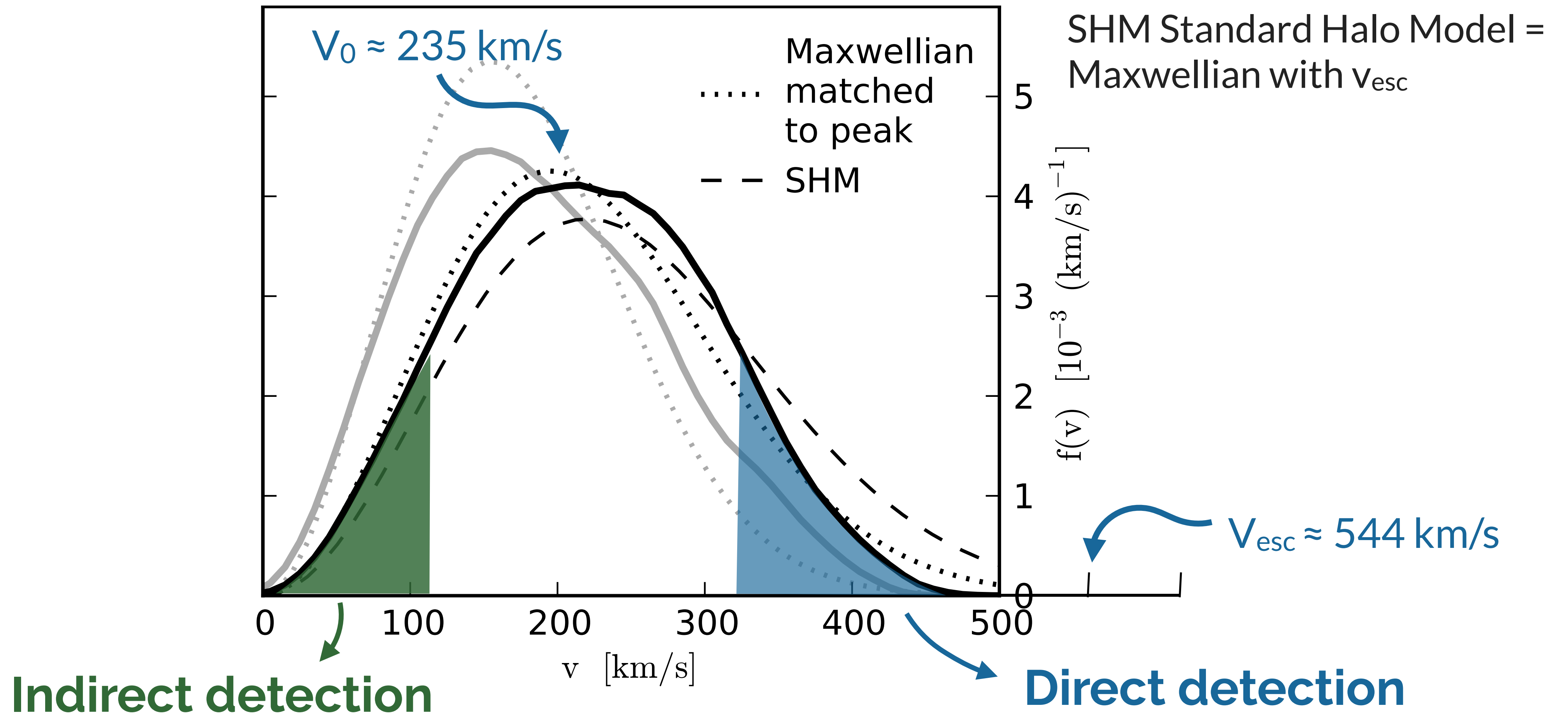
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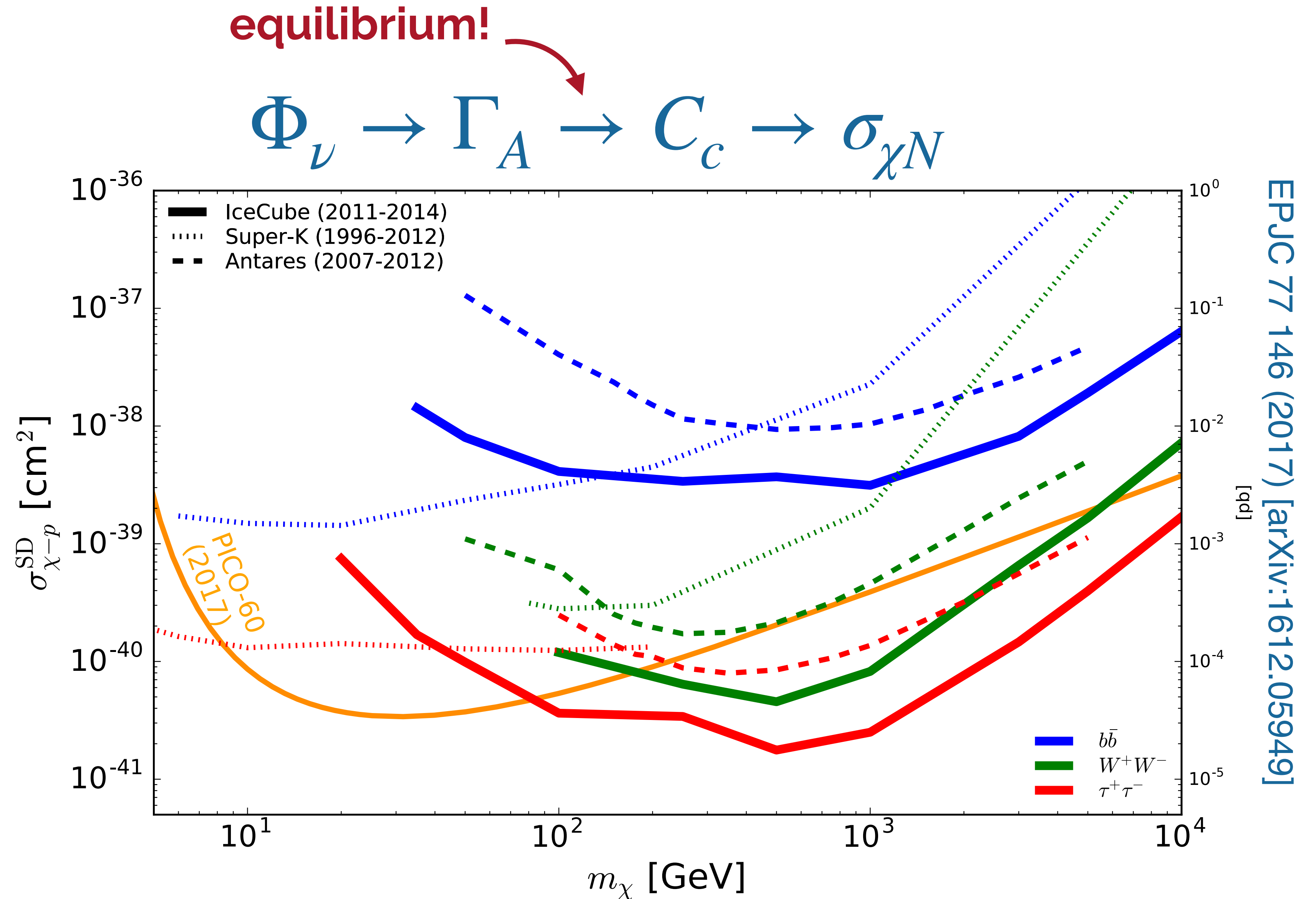


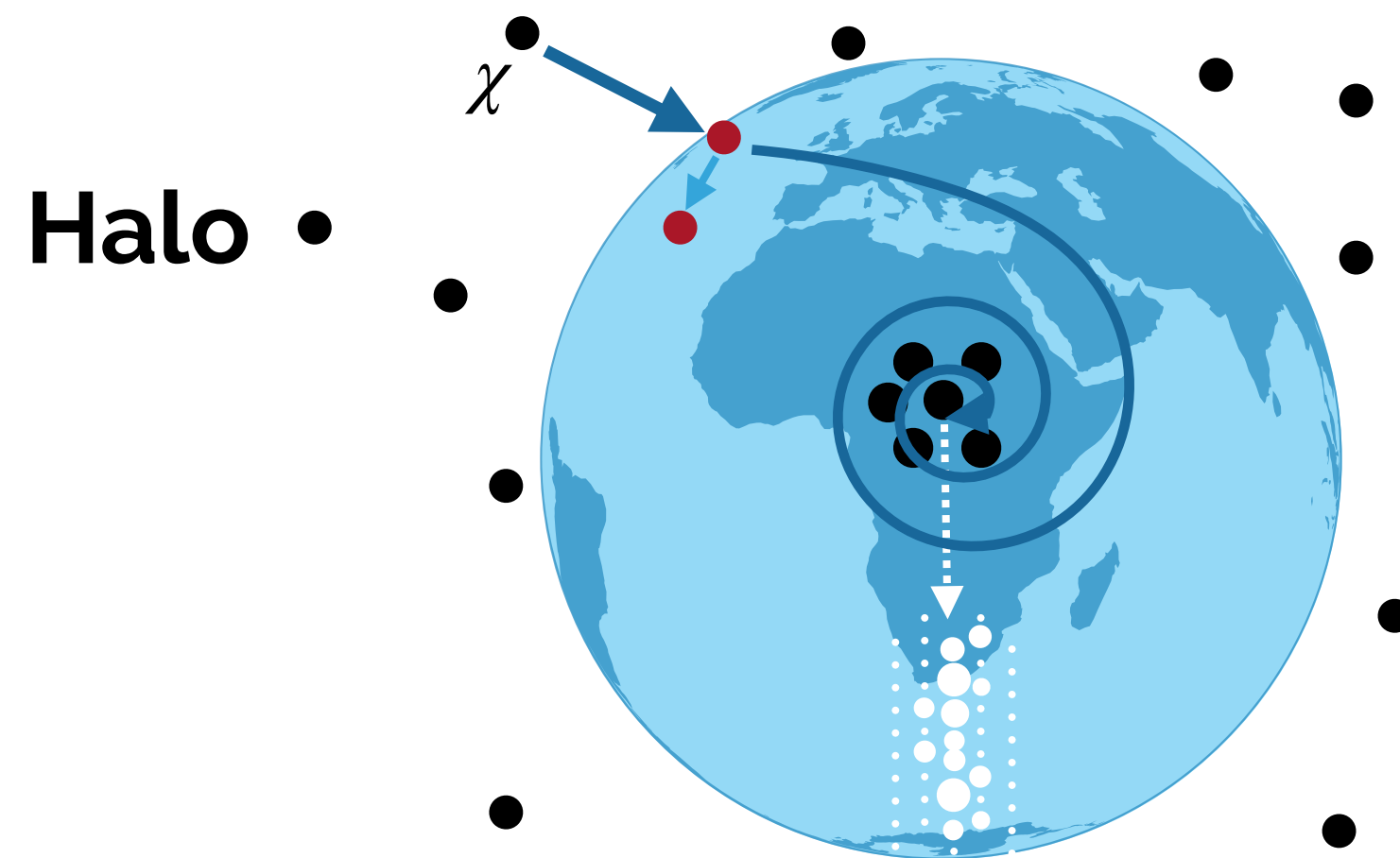
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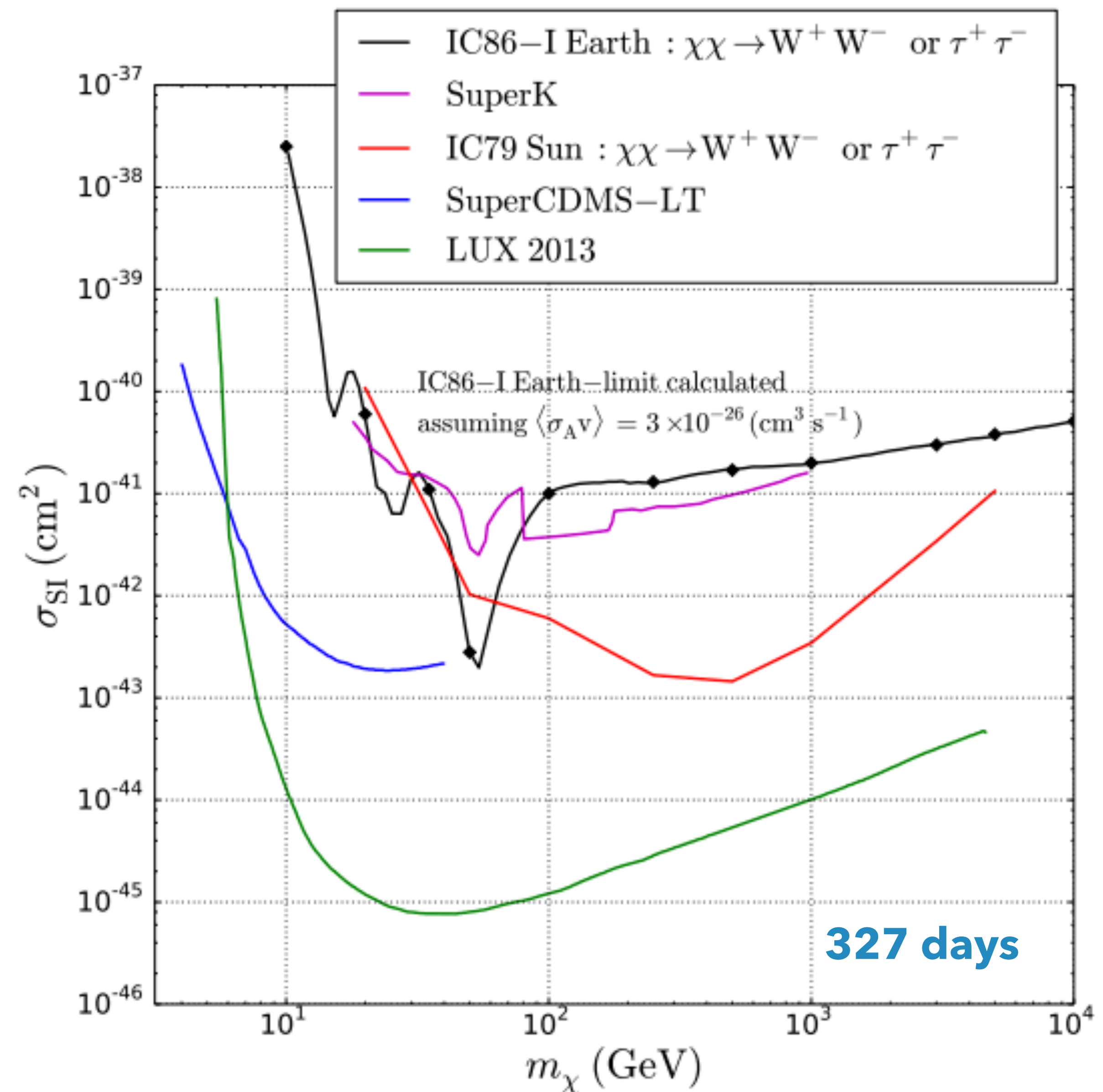


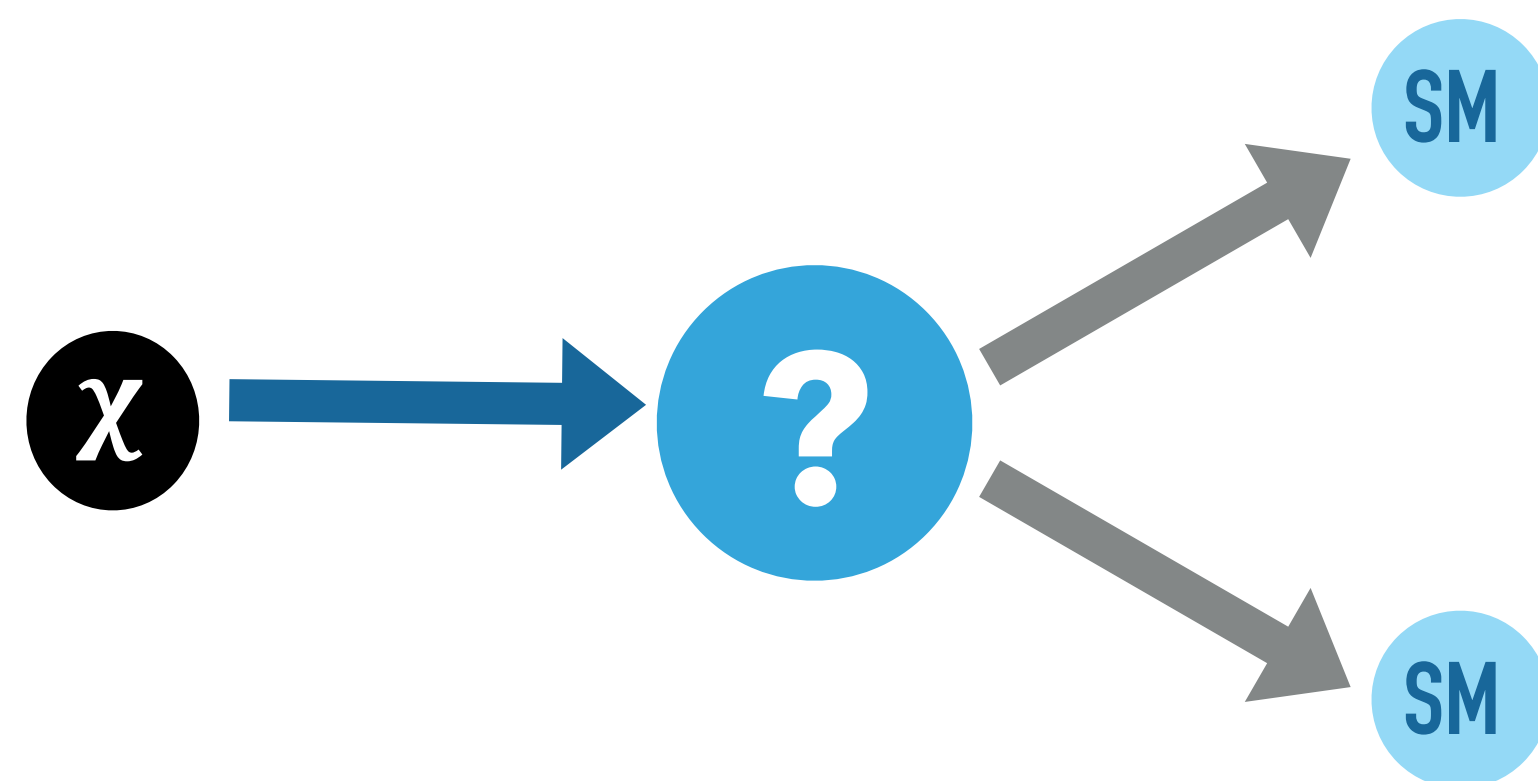
- Only events when Sun is below the horizon: **532 days of livetime**
- The mean free path of TeV neutrinos smaller than the Sun radius: **Low energy analysis**
- Best limits **$m_\chi > 80$ GeV**





- **No thermal equilibrium**
- Analysis very sensitive to astrophysical uncertainties (dark disc, velocity distribution)
- More years of data currently being analyzed





If dark matter is unstable its lifetime must exceed the age of the Universe

Galactic

$$\frac{d\Phi_\nu}{dE_\nu} = \frac{1}{4\pi} \frac{1}{m_\chi \tau_\chi} \frac{dN_\nu}{dE_\nu} \int_0^{\Delta\Omega} d\Omega \int_{l.o.s} \rho_\chi(r(s, \Psi, \theta)) ds$$

Not factor ², no so pronounced in the direction of the GC

Extra-Galactic

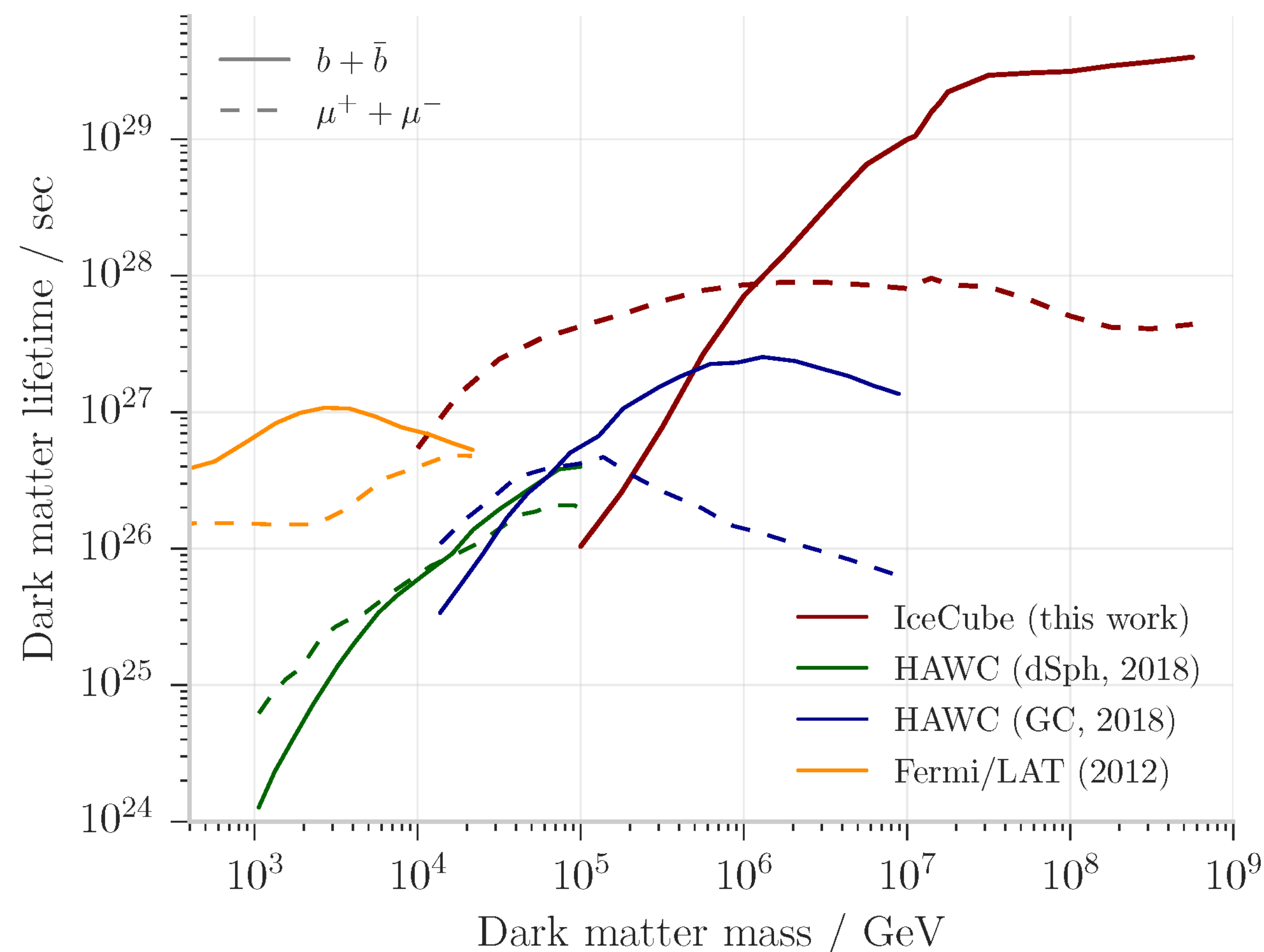
$$\frac{d\Phi_\nu}{dE_\nu} = \frac{1}{4\pi} \frac{\Omega_\chi \rho_c}{m_\chi \tau_\chi} \int_0^\infty dz \frac{c}{H(z)} \frac{dN_\nu}{dE_\nu} \Big|_{E=E_\nu(1+z)}$$

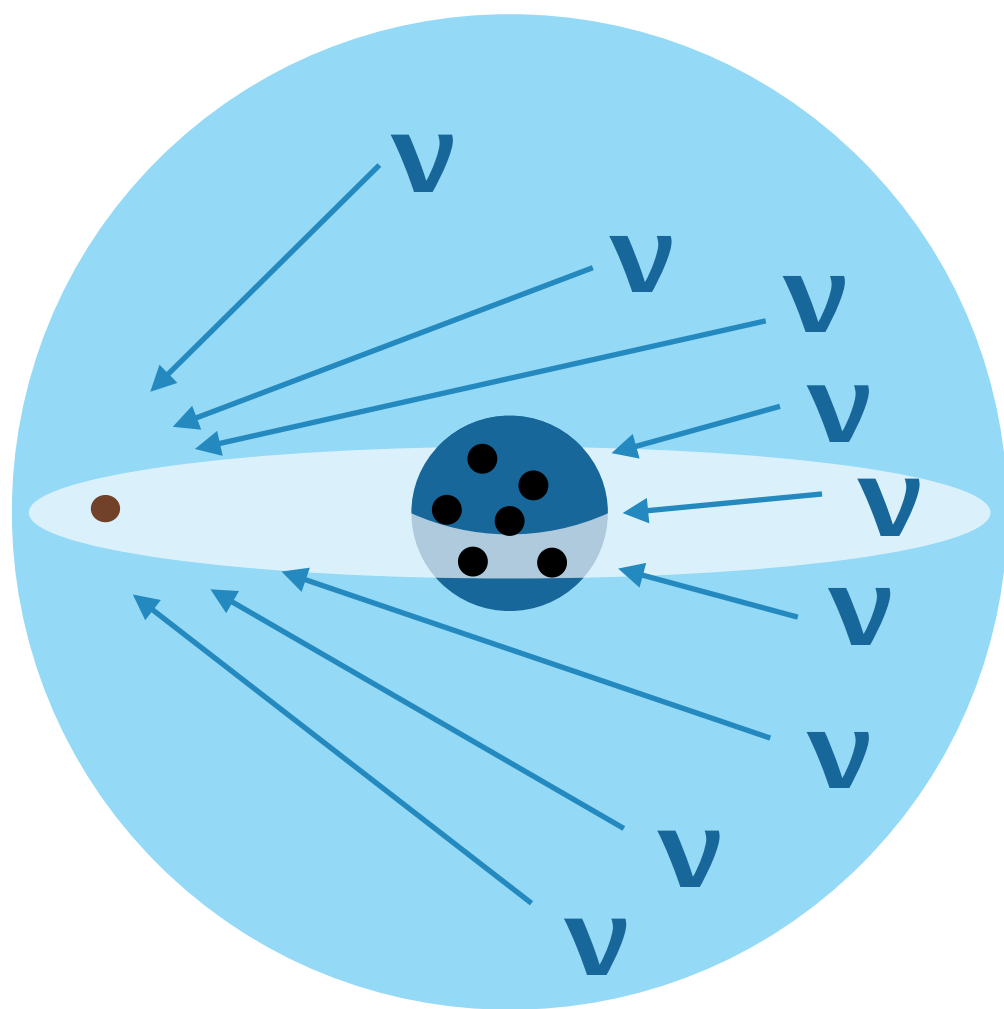
Extra-galactic component can be ignored for $\Psi < 10\text{deg}$.

constrain

- Two IceCube independent data samples:
 - Track-like with six years of data
 - Cascade-like with two years of data
- Dark Matter alone cannot explain IceCube neutrino flux.
- Best limits > 10 TeV

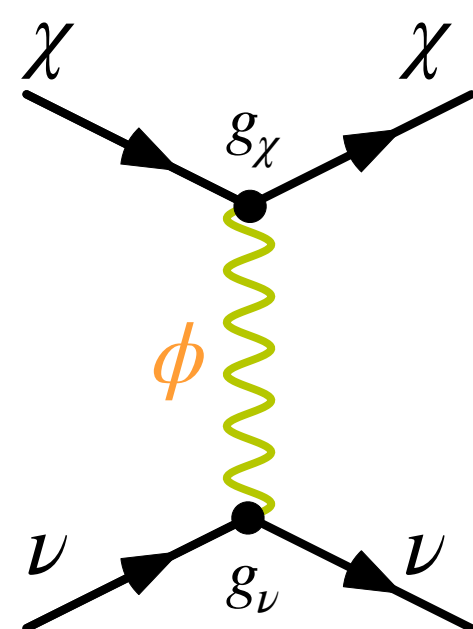
Eur. Phys. J. C (2018) 78: 831



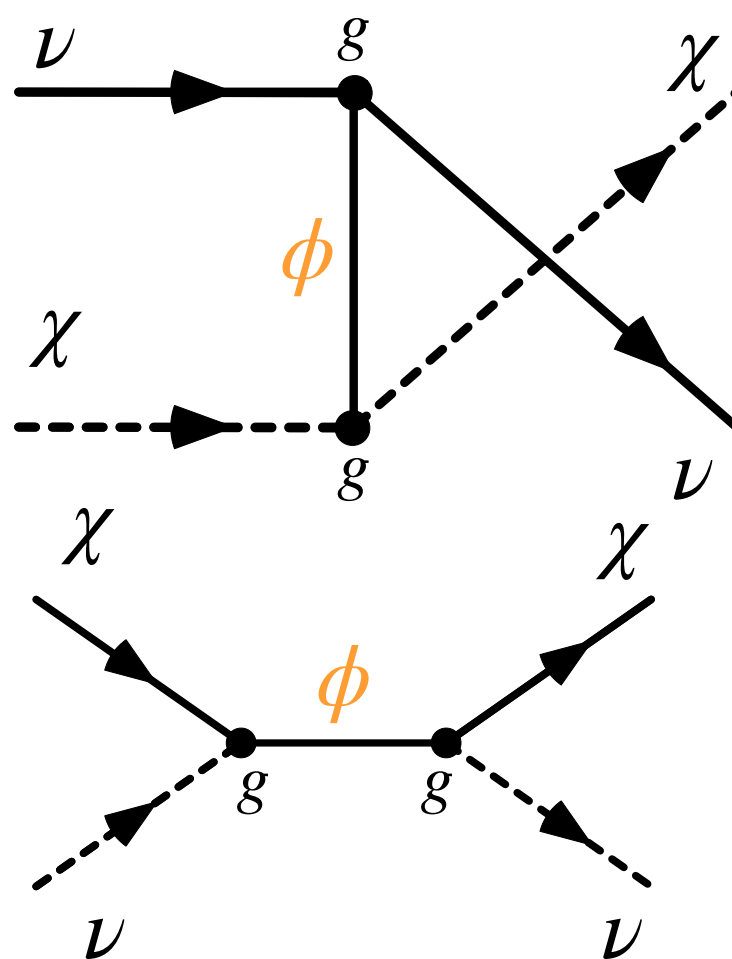


- Scattering of high energy cosmic neutrinos on DM in the halo can lead to a **deficit of high energy neutrinos** from the GC
- Focusing on HE neutrinos (cross-section increases with energy)

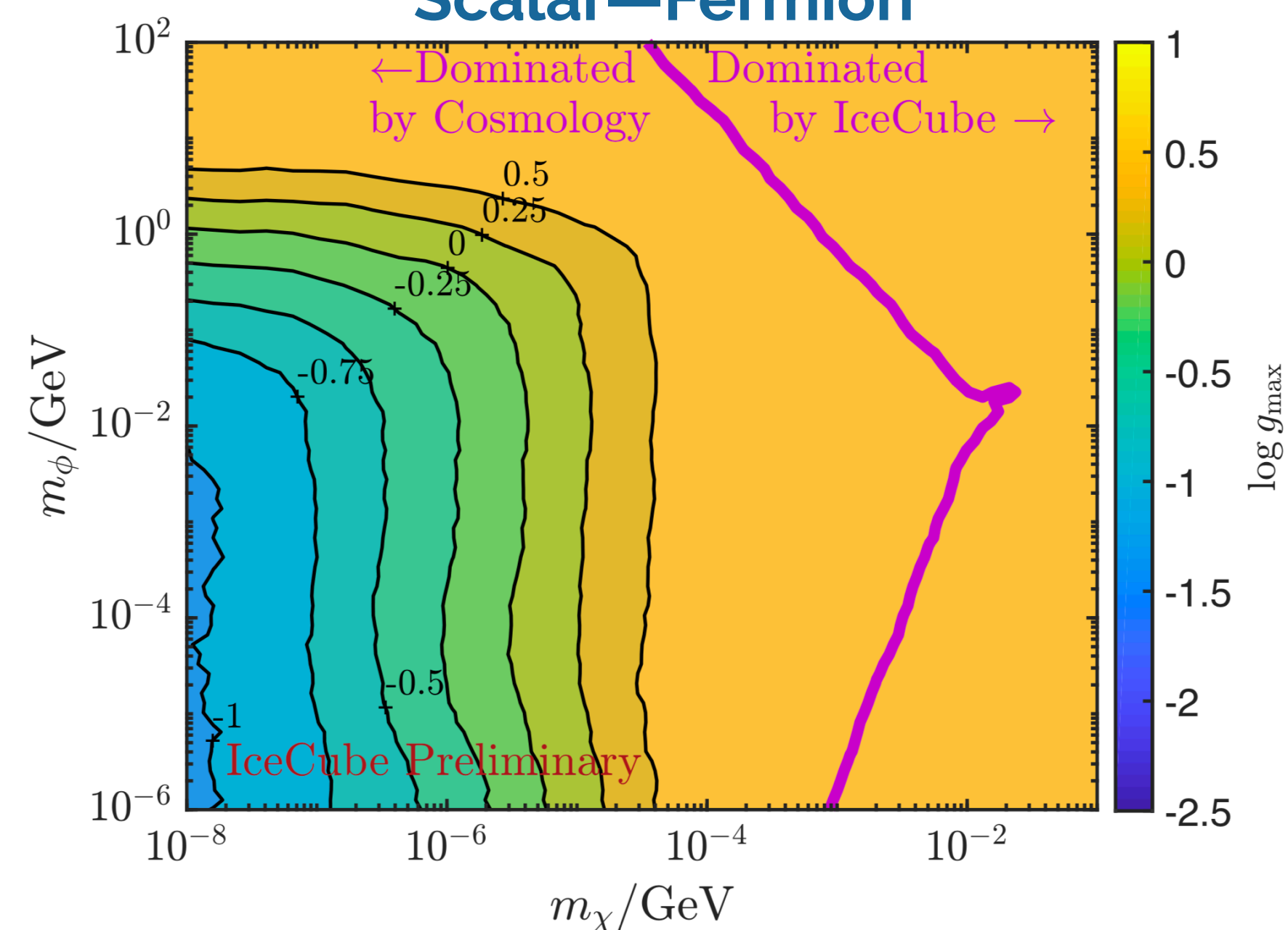
Fermion—vector



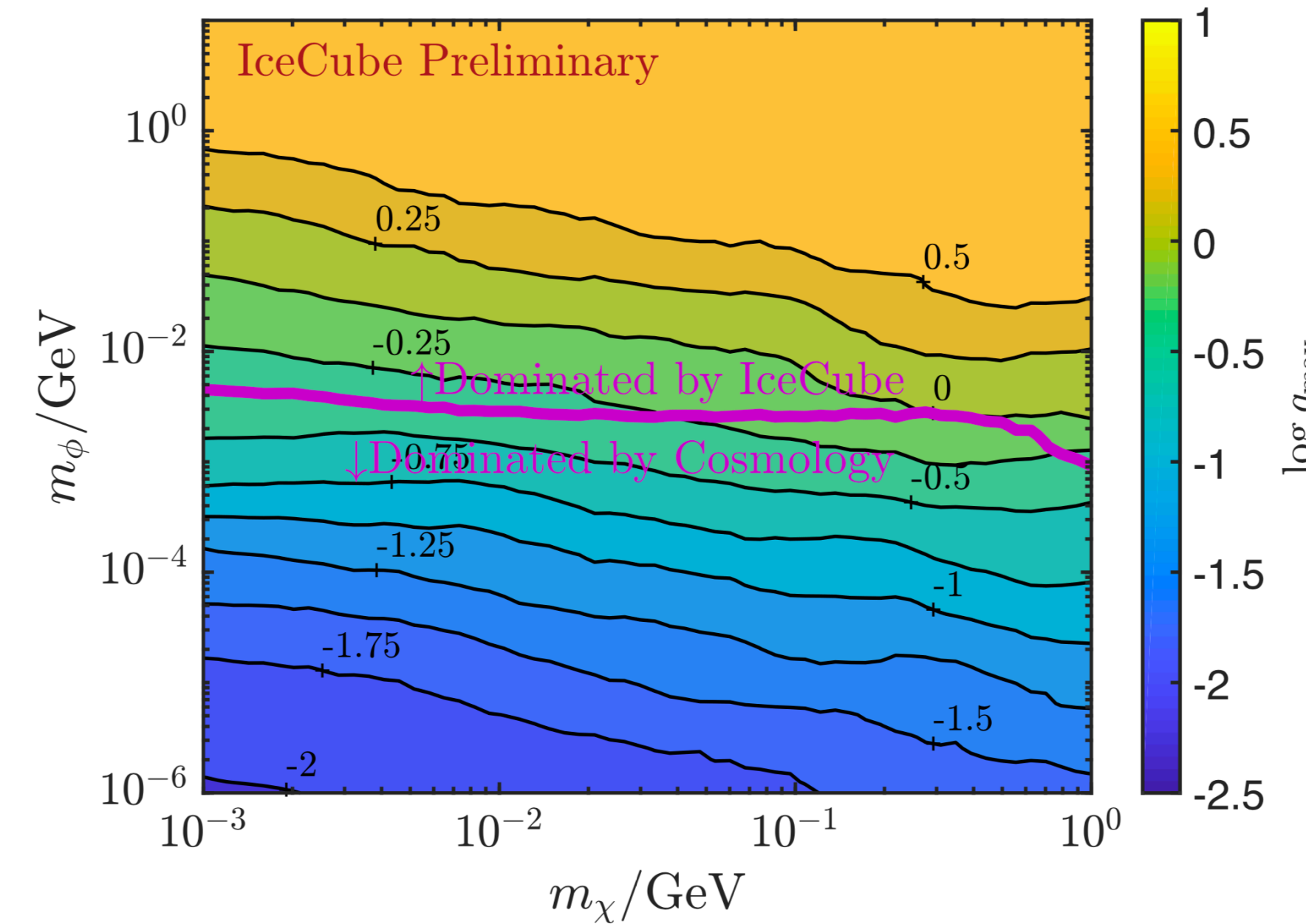
Scalar—Fermion



Scalar—Fermion



Fermion—vector

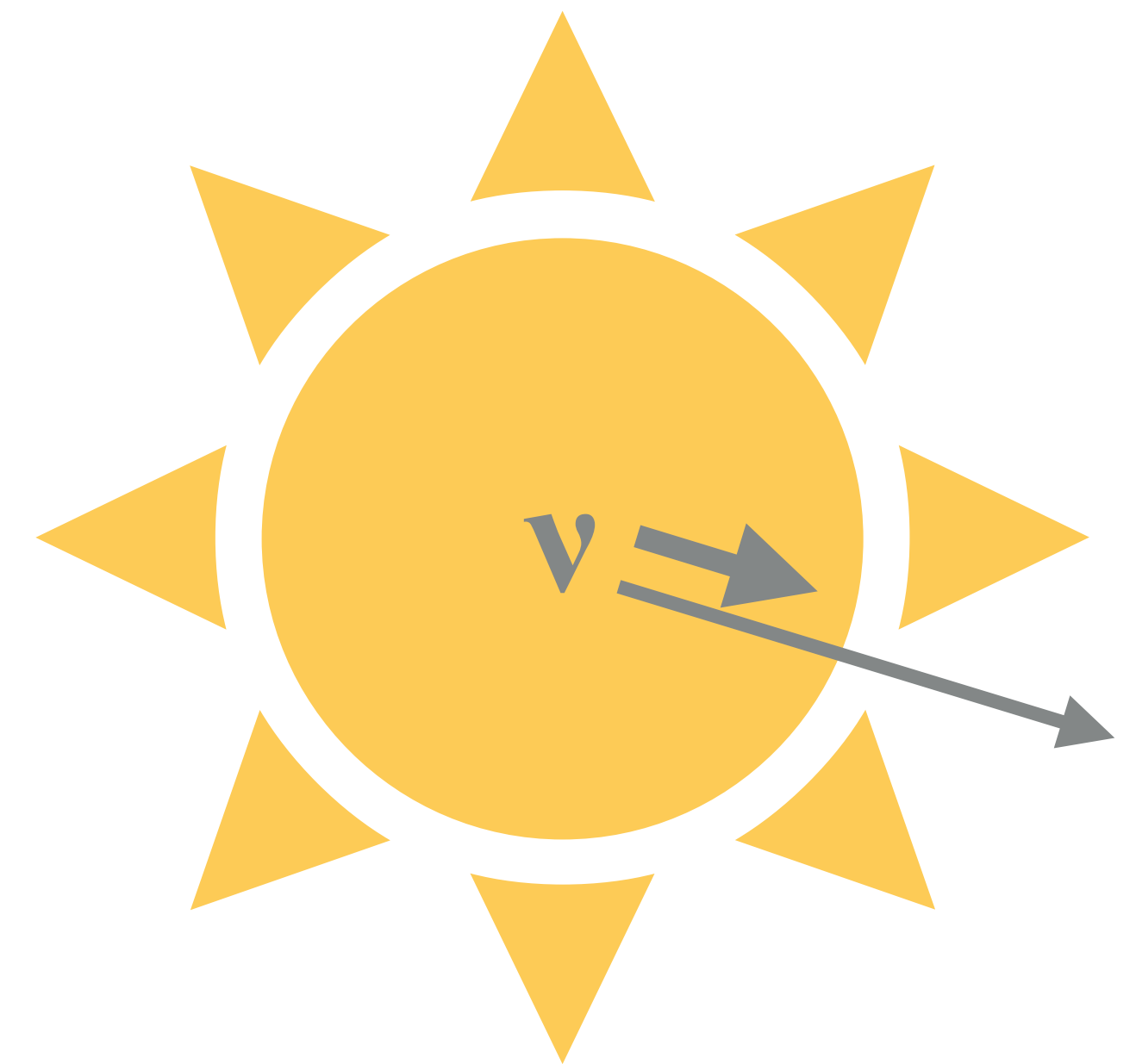
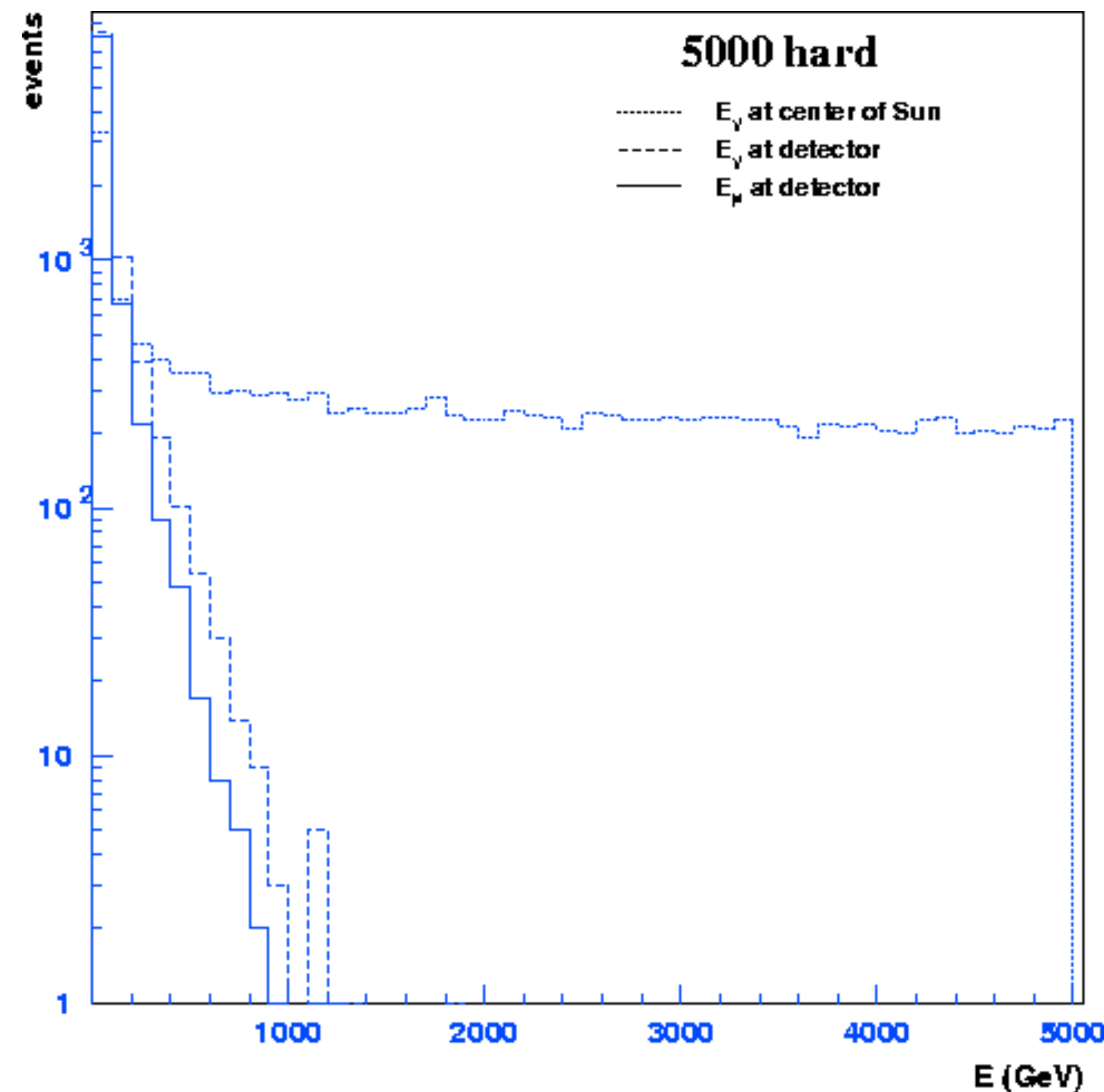


[Neutrino 2012, doi:10.5281/zenodo.1300506]

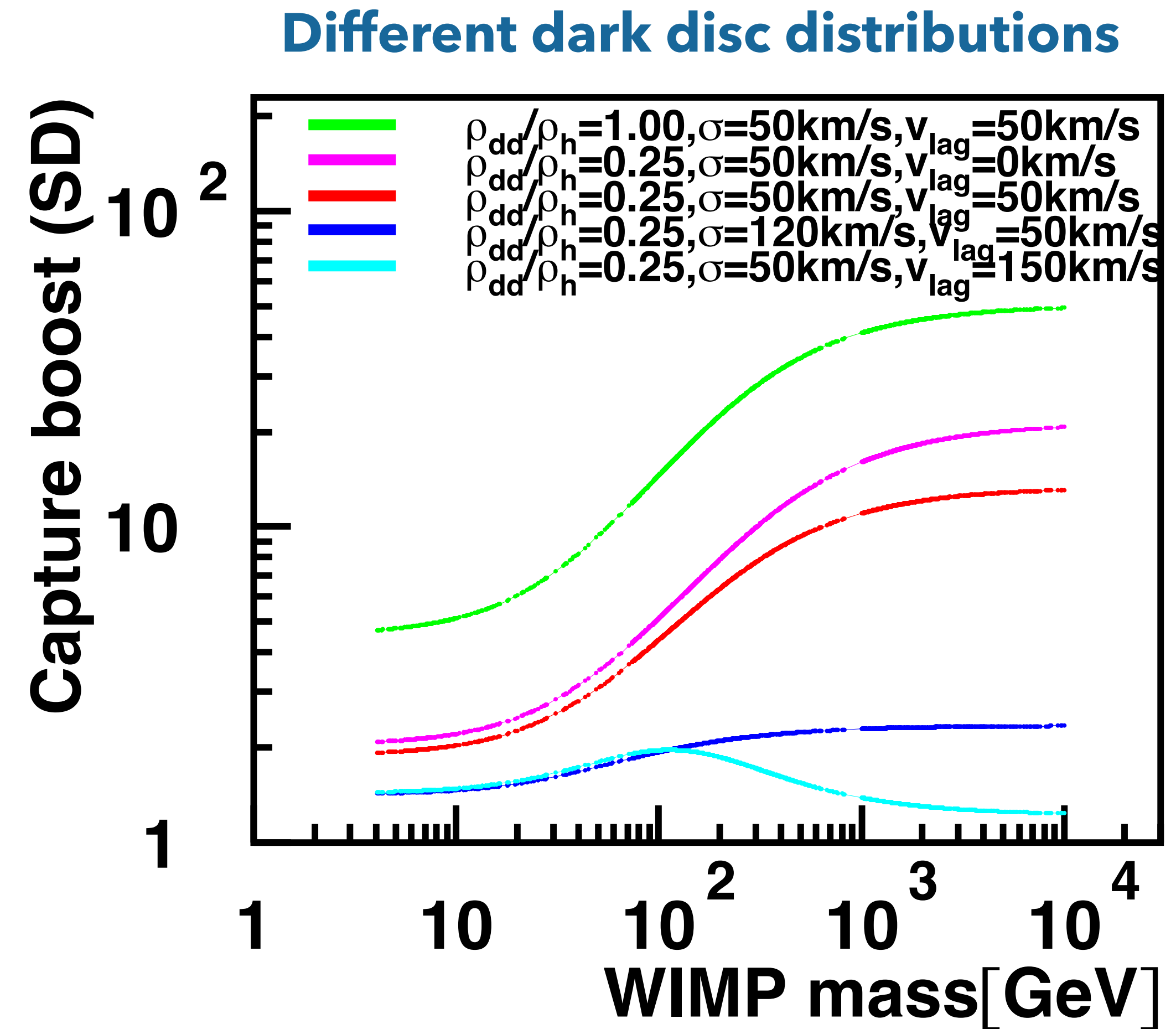
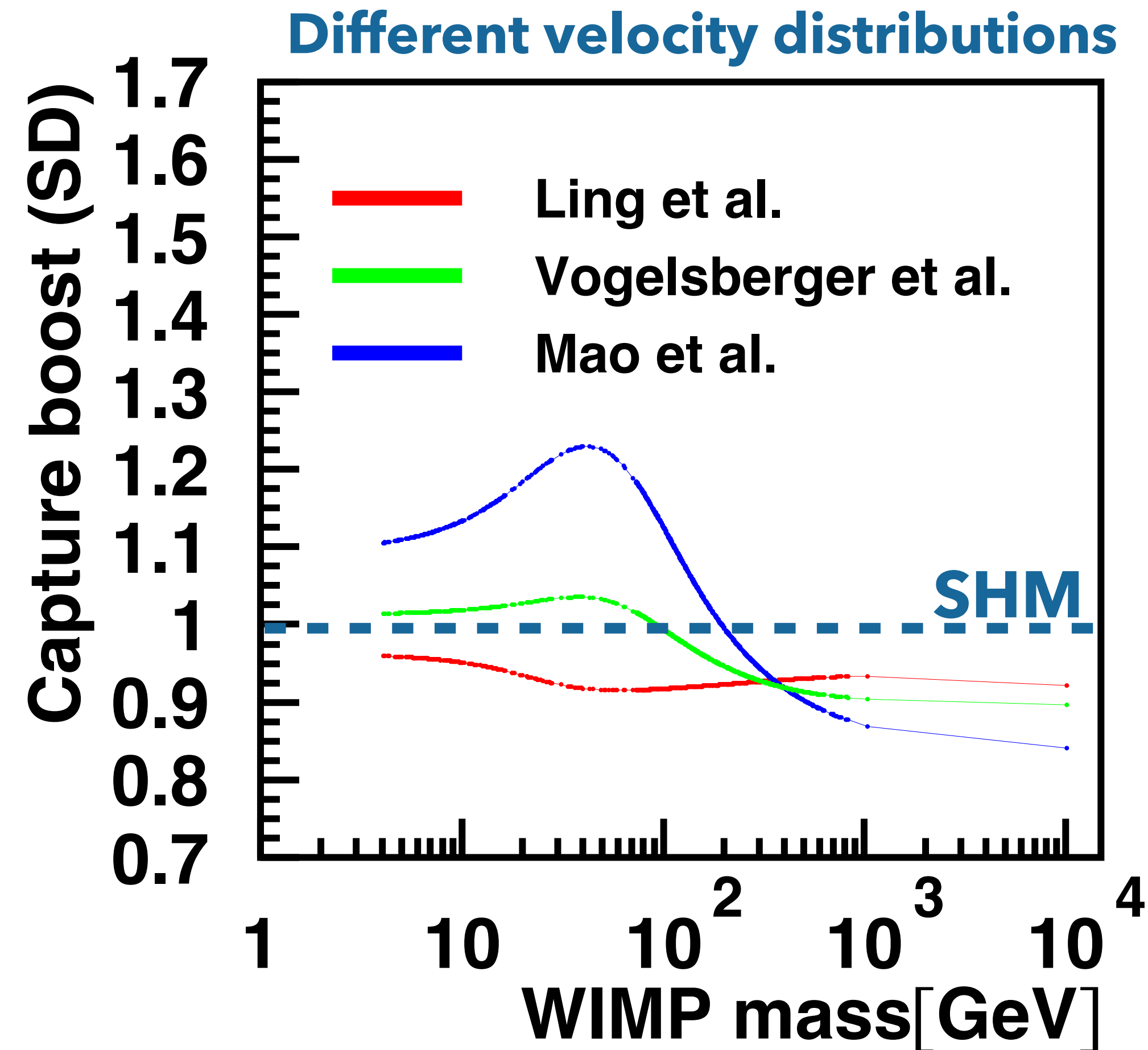
- **Indirect detection of Dark Matter** with neutrino telescopes provides complementarity to other techniques due to different backgrounds and systematics.
- Many astrophysical signals can be interpreted as Dark Matter. We need **strong corroboration** from all searching strategies.
- IceCube has a **lively program of Dark Matter searches**, with very competitive results.

backups

- ▶ The mean free path of neutrinos of 5000 GeV is smaller than the Sun radius
- ▶ Indirect searches from the Sun are **low-energy analysis** even for the highest dark matter masses.

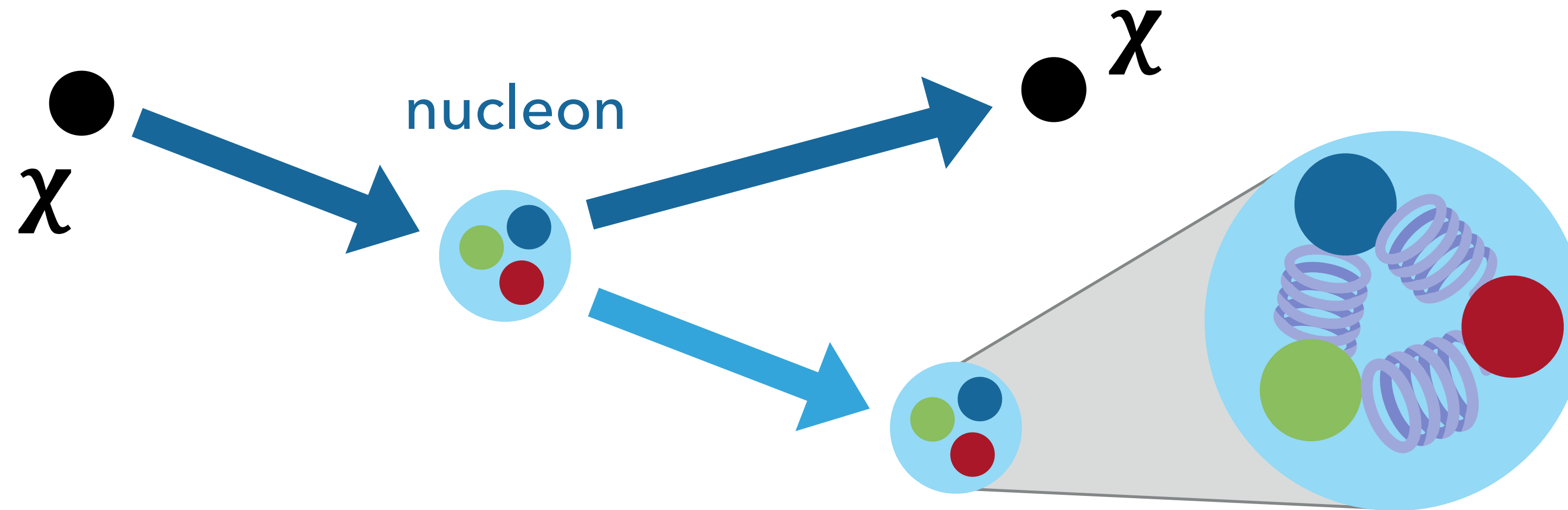


Effect of uncertainties in velocity distributions for Sun results:



Choi, Rott, Itow arXiv:1312.0273

A **dark matter disc** will have a significant (good) impact on the capture rate for the Sun/Earth



Both direct detection and indirect detection (gravitational capture) depend on the WIMP-nucleon cross-section.

$$\sigma_{SI} \propto A^2$$

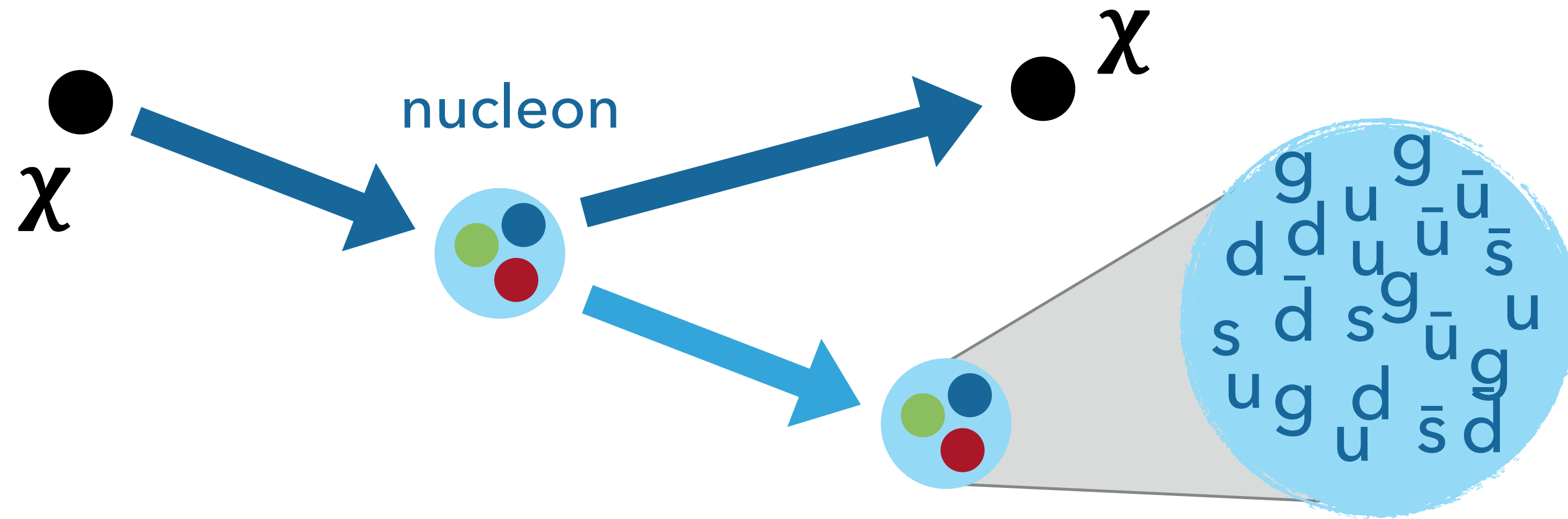
Spin independent

Use heavy nuclei as target: **Direct detection**

$$\sigma_{SD} \propto (a_p \langle S_p \rangle + a_n \langle S_n \rangle) \frac{J+1}{J} \frac{S(|\vec{q}|)}{S(0)}$$

Spin dependent

Sun is full of protons: **Indirect detection**



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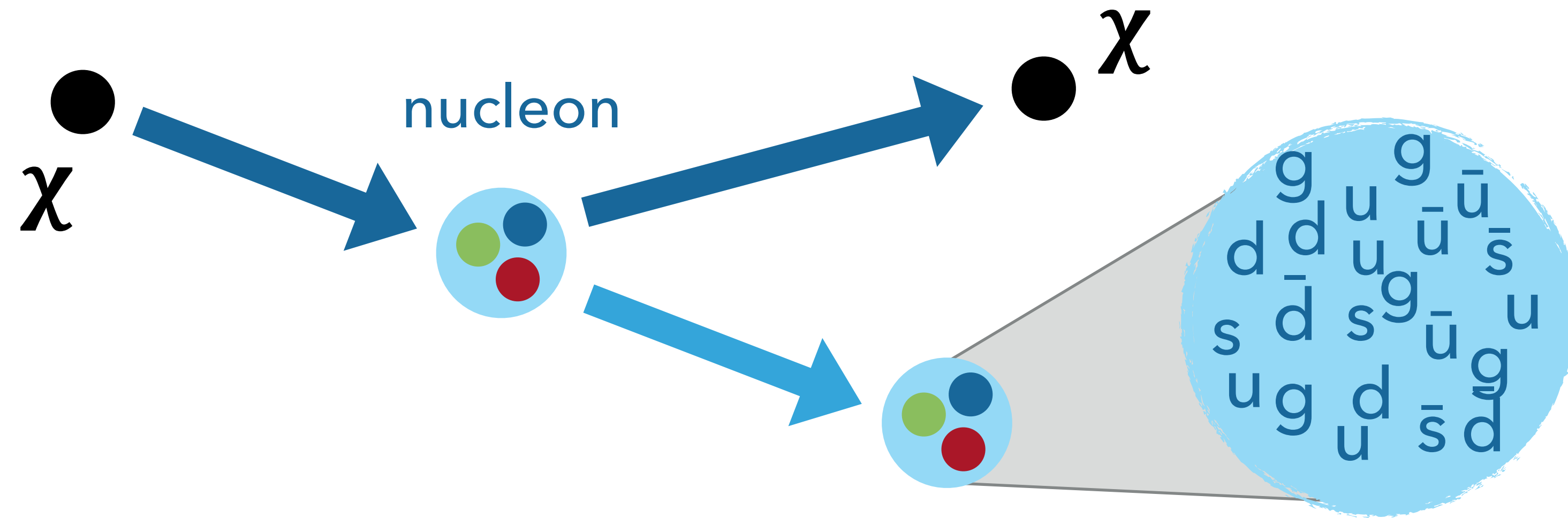
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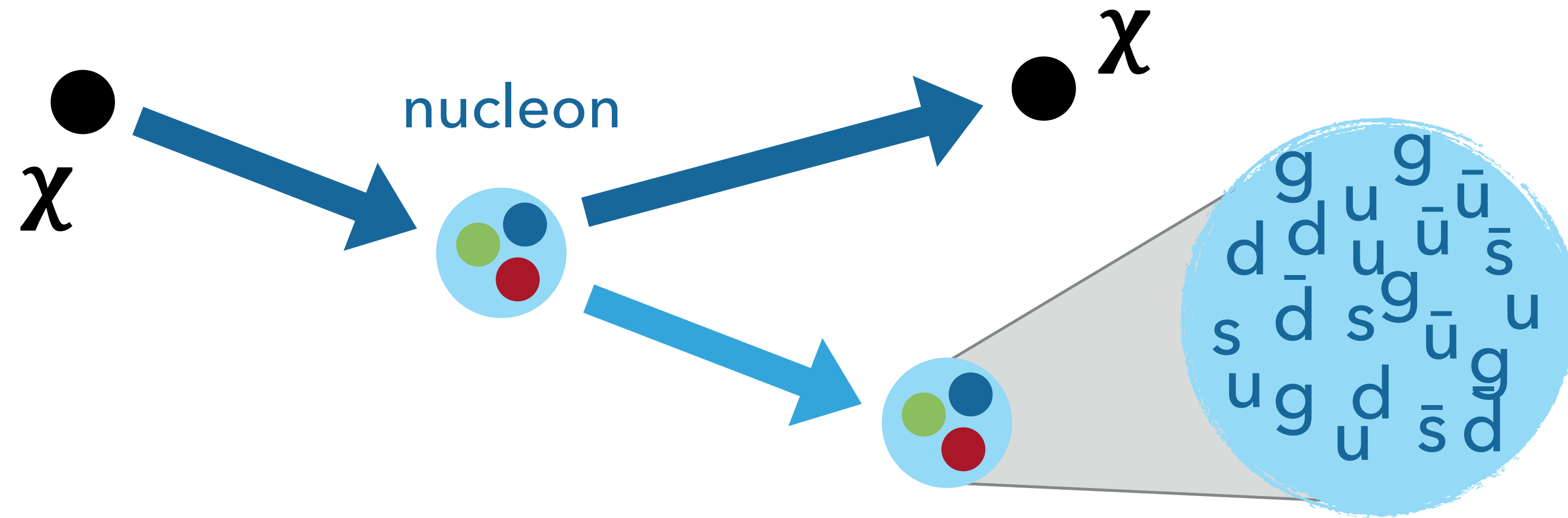
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Sun is full of protons: **Indirect detection**

The nucleon structure plays an essential role in calculating observables



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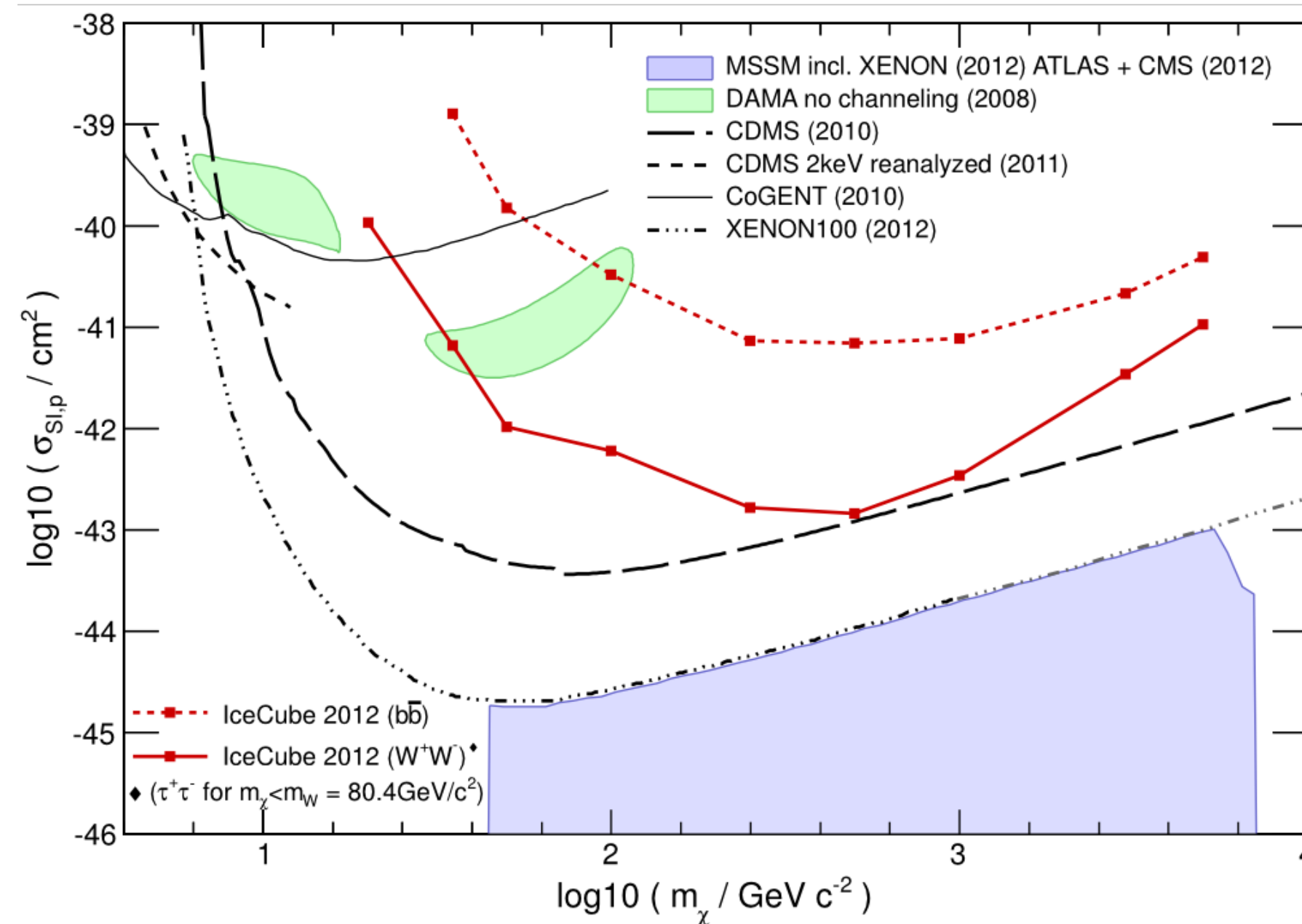
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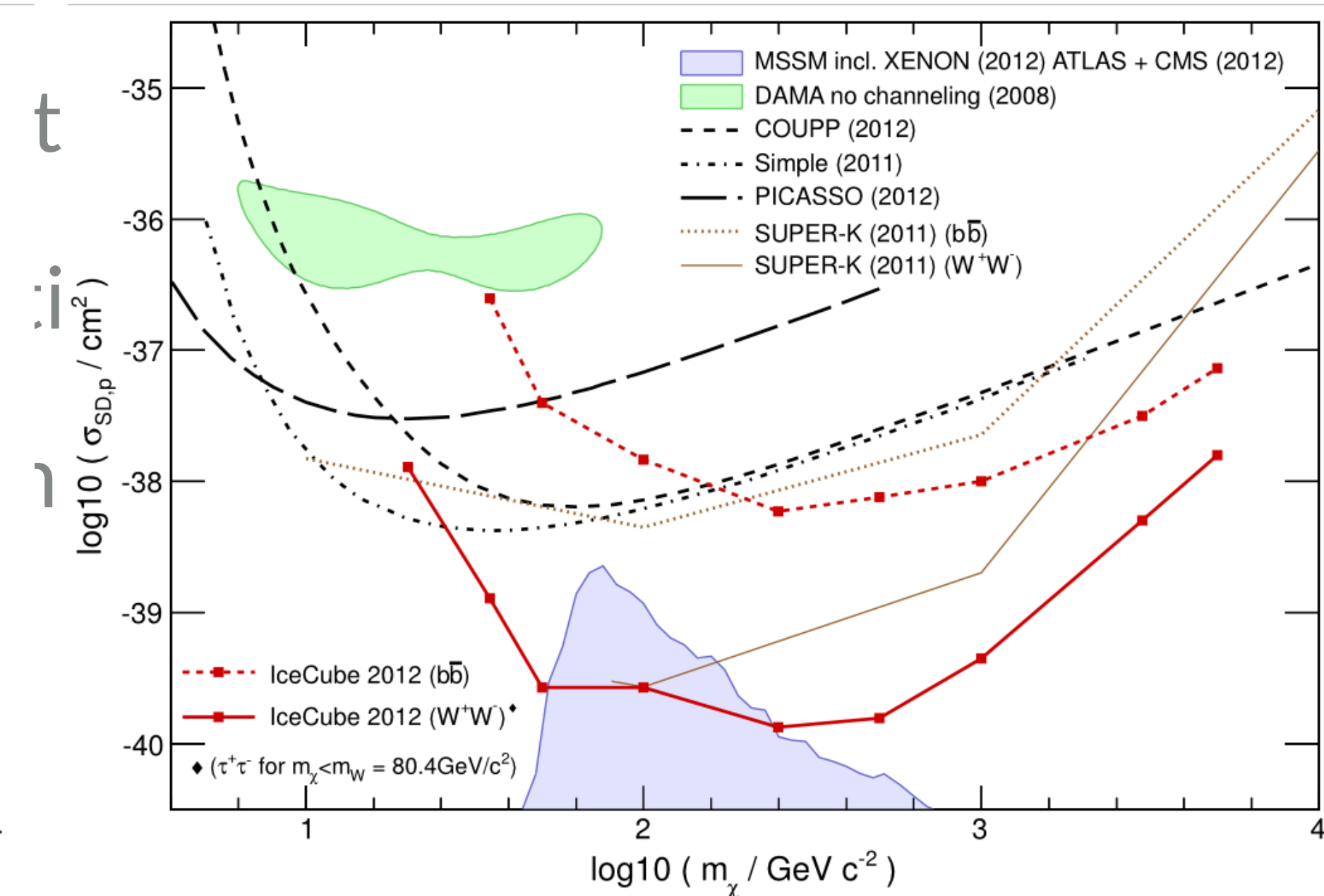
The nucleon structure plays an essential role in calculating observables

But it seems to affect more σ^{SI} than σ^{SD}

R. Ruiz, C. de los Heros arXiv:1307.6668



90% CL χ -p cross-section (spin-independent)



90% CL χ -p cross-section (spin-dependent)

Complementary to direct detection search efforts
fills out WIMP picture by testing other properties
Most stringent SD cross-section limit for most models

